

DPG-Frühjahrstagung Bonn, March 15–19, 2010

# Electroweak contributions to $e^+e^- \rightarrow W^+W^-b\bar{b}$ in the $t\bar{t}$ resonance region

# **Bernd Jantzen**

RWTH Aachen University

In collaboration with Martin Beneke and Pedro Ruiz-Femenía

- I Top-pair production at linear colliders near threshold
- II Evaluation of electroweak NLO hard contributions
- III Results & comparisons
- IV Conclusions & outlook

# **Top-pair production at linear colliders near threshold**

# **Future linear colliders** (ILC/CLIC)

with  $\sqrt{s} \gtrsim 2m_t \approx 350 \text{ GeV} \rightsquigarrow \text{ produce } t\bar{t} \text{ pairs:}$ clean initial state of  $e^+e^- \rightarrow t\bar{t}$  allows threshold scans with  $\sqrt{s} \sim 2m_t$ 

 $\hookrightarrow$  Precise determination of the top mass  $m_t$ , the width  $\Gamma_t$ , the Yukawa coupling  $\lambda_t$ without the uncertainties/ambiguities of hadron colliders. Martinez, Miquel '02

## Need also precise theoretical prediction

 $\Rightarrow$   $\delta\sigma/\sigma\sim$  2–3% ( $\delta\sigma\sim$  5 fb below threshold)

QCD corrections are known (almost) up to NNNLO order, but electroweak (EW) NLO hard contributions are missing!

Also: decay  $t\bar{t} \to (bW^+)(\bar{b}W^-)$  is an EW effect  $\Rightarrow$  describe  $t\bar{t}$  production in terms of the more physical process  $e^+e^- \to W^+W^-b\bar{b}$ and allow for invariant-mass cuts on reconstructed  $t, \bar{t}$ .

## Perturbative expansion: NRQCD

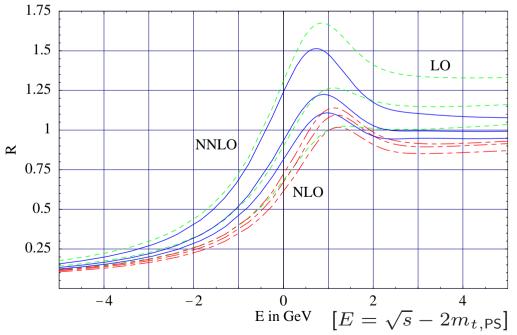
Top quarks move slowly at threshold:  $v = \sqrt{1 - \frac{4m_t^2}{s}} \sim \alpha_s \ll 1$   $\hookrightarrow \operatorname{sum} \left(\frac{\alpha_s}{v}\right)^n$  from "Coulomb gluons" to all orders:  $R = \frac{\sigma_{t\bar{t}}}{\sigma_{\mu^+\mu^-}} = v \sum_n \left(\frac{\alpha_s}{v}\right)^n \left(\{1\}_{\mathsf{LO}} + \{\alpha_s, v\}_{\mathsf{NLO}} + \{\alpha_s^2, \alpha_s v, v^2\}_{\mathsf{NNLO}} + \dots\right)$ 

Further RGE improvement by summing also  $(\alpha_s \ln v)^m$  to all orders: LL, NLL, ...

## Status of QCD corrections

- NNLO QCD corrections

   Hoang, Teubner '98–'99; Melnikov, Yelkhovsky '98;
   Yakovlev '98; Beneke, Signer, Smirnov '99 [see plot];
   Nagano, Ota, Sumino '99; Penin, Pivovarov '98–'99
- NNLO & (partial) NNLL
   Hoang, Manohar, Stewart, Teubner '00–'01;
   Hoang '03; Pineda, Signer '06
- (partial) NNNLO
   Beneke, Kiyo, Schuller '05-'08 [+ contributions from
   Kiyo, Seidel, Steinhauser '08; Anzai, Kiyo, Sumino '09; Smirnov, Smirnov, Steinhauser '09-'10]



3/12

**Effective field theory (EFT)** for pair production of unstable particles near threshold Beneke, Chapovsky, Khoze, Signer, Stirling, Zanderighi '01–'04; Actis, Beneke, Falgari, Schwinn, Signer, Zanderighi '07–'08

Non-relativistic power counting:

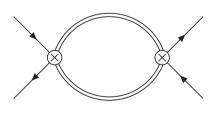
$$\label{eq:alphase} \boxed{\alpha_s^2 \sim \alpha_{\rm ew} \sim \frac{\Gamma_t}{m_t} \sim v^2 \approx \delta = \frac{s}{4m_t^2} - 1}$$

- Integrate out hard modes  $\sim m_t \rightsquigarrow$  EFT with potential (nearly on-shell) top quarks.
- Extract cross section  $e^+e^- \rightarrow W^+W^-b\bar{b}$  from appropriate cuts of the

 $e^+e^- \rightarrow e^+e^-$  forward-scattering amplitude:

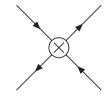
$$i\mathcal{A} = \underbrace{\sum_{k,l} \int \mathrm{d}^4 x \, \langle e^+ e^- | T \left[ i\mathcal{O}_p^{(k)\dagger}(0) \, i\mathcal{O}_p^{(l)}(x) \right] | e^+ e^- \rangle}_{k} + \underbrace{\sum_{k} \langle e^+ e^- | i\mathcal{O}_{4e}^{(k)}(0) | e^+ e^- \rangle}_{k}$$

resonant contributions with  $t\bar{t}$  production operators  $\mathcal{O}_p^{(k)}$ 



non-resonant contributions

4/12



correspond to full-theory  $e^+e^- \rightarrow e^+e^-$  with  $\Gamma_t=0$ 

⇒ Potential (+ soft . . .) corrections to resonant diagrams within EFT ⇒ Hard corrections to matching coefficients of operators  $\mathcal{O}_{p}^{(k)}$  and  $\mathcal{O}_{4e}^{(k)}$ .

## **Electroweak effects at LO**

• Replacement rule  $E = \sqrt{s} - 2m_t \rightarrow E + i\Gamma_t$  Fadin, Khoze '87

#### **Electroweak effects at NLO**

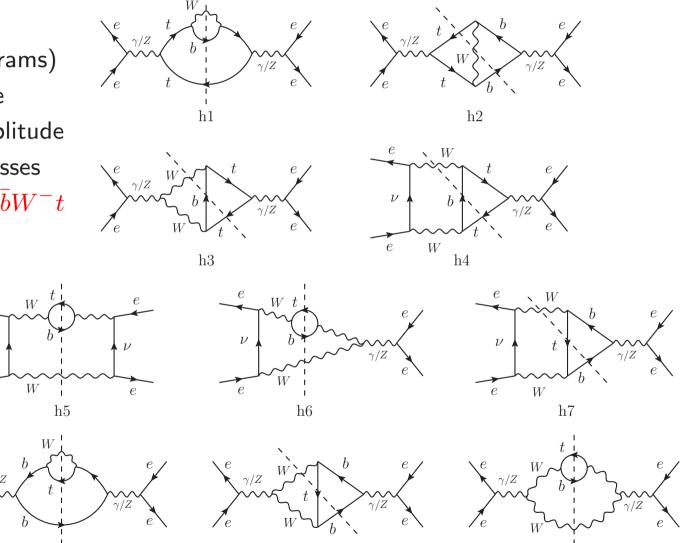
- Exchange of a "Coulomb photon": trivial extension of QCD corrections, available
- Gluon exchange between  $t, \bar{t}$  and their decay products: these contributions cancel at NLO in the total cross section, Fadin, Khoze, Martin '94; Melnikov, Yakovlev '94 they are negligible if the top invariant-mass cuts are loose enough.
- Hard corrections to the matching coefficient of the non-resonant operator  $\mathcal{O}_{4e}^{(k)}$  $\hookrightarrow$  topic of this talk!

The resonant NNLO corrections involve "finite-width divergences"  $\propto \alpha_s \frac{\Gamma_t}{\epsilon}$  (in dim. reg.). These must be cancelled by EW NNLO hard contributions.  $\hookrightarrow$  Motivation for calculating EW hard corrections (starting at NLO ...).

# **Evaluation of electroweak NLO hard contributions**

## Hard corrections at NLO:

- cuts through  $bW^+\bar{t}$  (see diagrams) and  $\bar{b}W^-t$  (not shown) in the 2-loop forward-scattering amplitude
- correspond to tree-level processes  $e^+e^- \rightarrow bW^+\bar{t}$  and  $e^+e^- \rightarrow \bar{b}W^-t$
- expansion in  $\delta = \frac{s}{4m_t^2} 1$  $\hookrightarrow$  at NLO:  $s = 4m_t^2$



h9

h8

ν

e

[symmetric diagrams not shown]

h10

6/12

## Form of hard contributions

With the reconstructed top momentum  $p_t = p_b + p_{W^+}$ , the contributions of diagrams h1-h10 (for  $s = 4m_t^2$ ) are of the form:

7/12

$$H_{i} = \int_{\Delta^{2}}^{m_{t}^{2}} \mathrm{d}p_{t}^{2} h_{i} \left(\frac{p_{t}^{2}}{m_{t}^{2}}, \frac{M_{W}^{2}}{m_{t}^{2}}\right)$$

with  $\Delta^2 = M_W^2$  for the total cross section.

#### Invariant-mass cuts

Restrict invariant masses of the reconstructed  $t, \bar{t}: |\sqrt{p_{t,\bar{t}}^2 - m_t}| \leq \Delta M$  $\hookrightarrow$  hard contributions with  $\Delta^2 = m_t^2 - \Lambda^2$  where  $\Lambda^2 = (2m_t - \Delta M)\Delta M$ .

We focus on loose cuts with  $\Lambda^2 \gg m_t \Gamma_t$  (typical offshellness of potential top quarks), corresponding to  $\Delta M \gg \Gamma_t \rightsquigarrow$  no cut needed for potential contributions.

In contrast: for tight cuts with  $\Lambda^2 \lesssim m_t \Gamma_t$  or  $\Delta M \lesssim \Gamma_t \rightsquigarrow$  hard contributions are absent and potential contributions need to be cut.

# III Results & comparisons

**EW NLO hard contributions:** numeric integration over  $p_t^2$  (and over one angle for some diagrams), the integrand is an analytic function of  $p_t^2/m_t^2$  and  $M_W^2/m_t^2$ , the cut-dependence enters in the integration limit.

**Parameters:**  $m_t = 172 \text{ GeV}, \ \Gamma_t = \Gamma_t^{\text{tree}} = 1.46550 \text{ GeV}, \text{ on-shell (pole) masses}, \alpha \text{ and } \sin^2 \theta_w \text{ from } G_{\text{F}}, M_W, M_Z$ 

**Comparison to recent alternative approach (HRR)** 

Hoang, Reißer, Ruiz-Femenía '10

• Here QCD & EW contributions expanded for moderate invariant-mass cuts  $15 \text{ GeV} \le \Delta M \le 35 \text{ GeV}$ 

 $\hookrightarrow$  our result is also valid for larger cuts up to the total cross section.

- EW contributions match expansion of our result in  $(\Lambda/m_t)^n$  up to n=1
- HRR had to neglect the non-t diagrams h5-h10, these "single-resonant" contributions are systematically included in our calculation.

#### **Comparison to MadGraph/MadEvent/MadAnalysis (MG)**

Alwall et al. '07

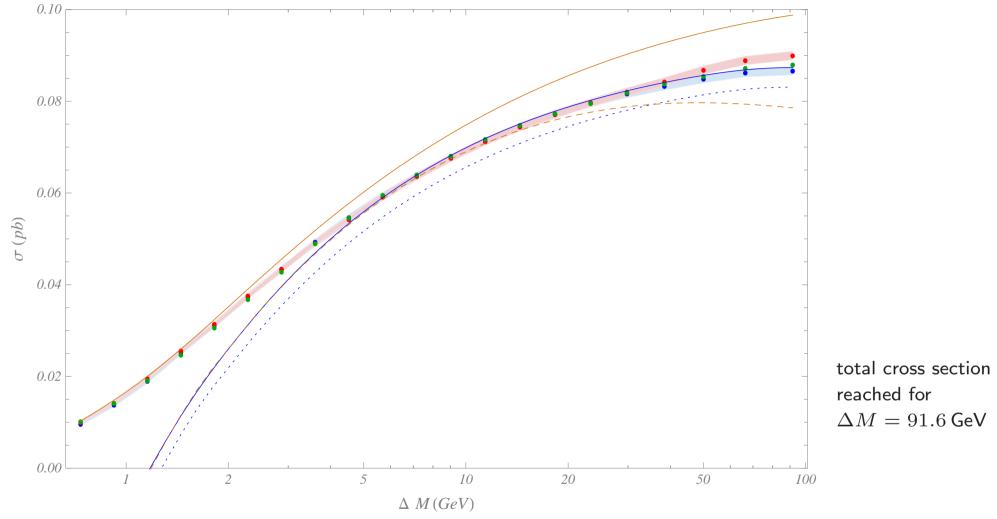
8/12

 $\hookrightarrow$  generated 10<sup>4</sup> events for  $e^+e^- \to W^+W^-b\bar{b}$ , analyzed cut-dependence

#### EW contributions: cut-dependence at threshold

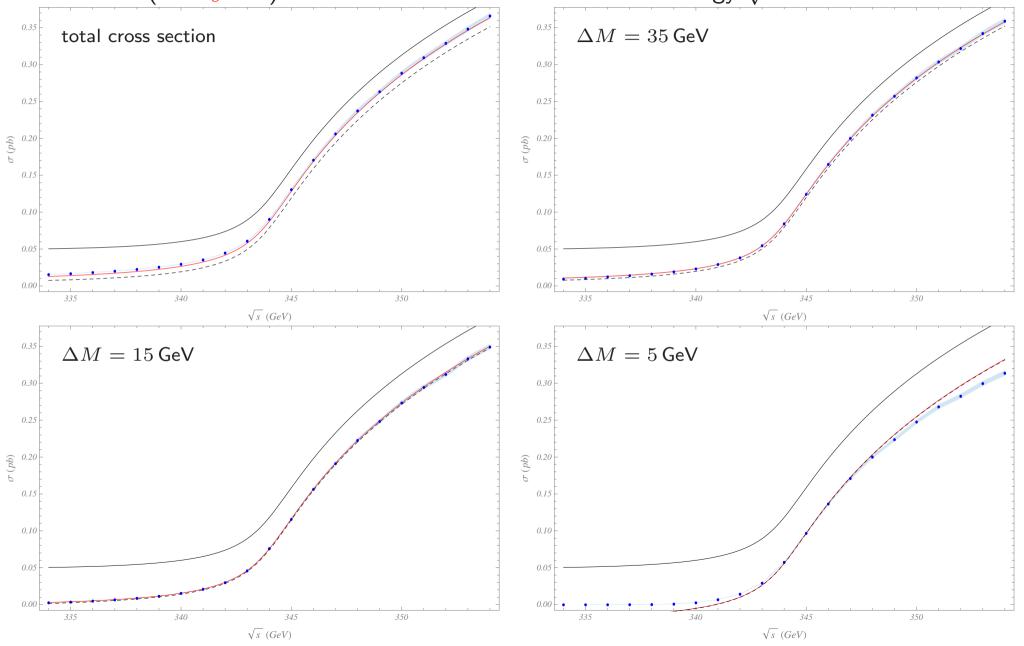
cross section (for  $\alpha_s = 0$ ) at threshold ( $s = 4m_t^2$ ) as a function of the invariant-mass cut  $\Delta M$ 

9/12



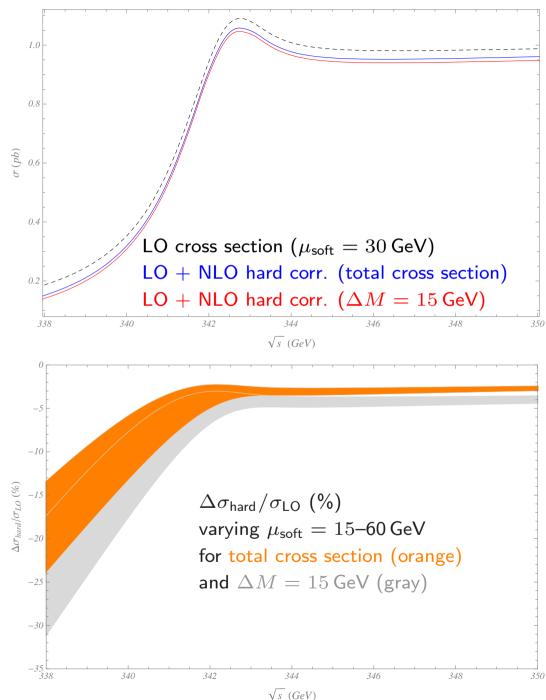
MG points (with statistical error bands): full, without Higgs, only *t*- or  $\bar{t}$ -diagrams Our result: EW NLO hard + LO (dashed-blue) / NNLO (solid-blue) potential contributions  $\hookrightarrow$  good agreement with MG for loose cuts  $\Delta M \gtrsim 5 \text{ GeV}$ Cut potential region (LO): solid-brown  $\Rightarrow$  good agreement with MG for tight cuts  $\Delta M \lesssim 1 \text{ GeV}$ HRR result: dashed-brown  $\Rightarrow$  agrees with our result for small  $\Delta M$ 

# **EW contributions: energy-dependence for different cuts** cross section (for $\alpha_s = 0$ ) as a function of the center-of-mass energy $\sqrt{s}$



MG (full) points & error band, our result, HRR result (dashed), only potential contributions (solid)

#### Full cross section with QCD & EW contributions



QCD contributions with  $\alpha_s^{\overline{\text{MS}}}(\mu_{\text{soft}})$  (obtained from  $\alpha_s^{\overline{\text{MS}}}(M_Z) = 0.118$ )

LO cross section (potential QCD & EW contributions, dashed-black) and including our EW NLO hard contributions (solid-colored)

[NLO QCD contributions not shown]

**Relative correction** (in %) of EW NLO hard contribution w.r.t. LO cross section

# **IV Conclusions & outlook**

Electroweak contributions to  $e^+e^- \rightarrow W^+W^-b\bar{b}$  in the  $t\bar{t}$  resonance region

- NLO contribution completed by EW non-resonant (hard) contributions for total cross section and with top invariant-mass cuts
- correction of  $\sim$  -30 fb (-3% above and much more below threshold) for total cross section, even more with invariant-mass cuts
- good agreement with MadGraph for loose cuts
- good agreement with Hoang–Reißer–Ruiz-Femenía result for small cuts
- $\hookrightarrow$  can be added to existing QCD results to improve accuracy of theoretical prediction

## **Future improvements**

- add initial-state radiation and convolution with electron distribution functions
- evaluate leading EW NNLO contributions  $\Rightarrow$  cancel finite-width divergences