

# **Quantum-mechanical tunneling time of an electron wave packet**

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- I Tunneling delay time measurements in Helium  
P. Eckle et al., Science 322 (2008) 1525
- II Free wave packet
- III Stationary solution for rectangular tunnel barrier
- IV Numerical solution for time-dependent tunneling in three cases:
  - A: “Slow” wave packet, thin barrier → “tunneling of the low momenta”
  - B: “Slow” wave packet, thick barrier → “transmission of the high momenta”
  - C: “Fast” wave packet, thick barrier → time-resolved transmission

# I Tunneling delay time measurements in Helium

P. Eckle<sup>1</sup>, A.N. Pfeiffer<sup>1</sup>, C. Cirelli<sup>1</sup>, A. Staudte<sup>2</sup>, R. Dörner<sup>3</sup>, H.G. Muller<sup>4</sup>,

M. Büttiker<sup>5</sup>, U. Keller<sup>1</sup>, *Attosecond Ionization and Tunneling Delay Time*

*Measurements in Helium, Science 322 (2008) 1525*

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**Aim:** Measure the tunneling time of a bound electron through the Coulomb potential of a helium atom during ionization.

## Measurement:

- Intense laser pulse hits helium atoms → lowers Coulomb potential → enables tunneling of bound electron to the outside → ionization
- At the exit of the tunnel the electron is accelerated by electric laser field → laser polarization determines direction of electron drift momentum
- Use circularly/elliptically polarized laser pulse → electron drift direction corresponds to laser field polarization at the time of tunnel exit  
→ time measurement with attosecond accuracy (“angular streaking”)
- Measure drift of helium instead of electrons → correspondence by simple semi-classical simulation assuming instantaneous tunneling (with Coulomb correction)

## Calibration of absolute time zero:

- ellipticity of laser polarization induces oscillations in the angular ion momentum distributions
- calibration by linear polarized laser pulse

**Result:** upper limit of 34 as for tunneling time

(intensity-averaged upper limit: 12 as)

## Comparison to theoretical predictions:

- Instantaneous tunneling agrees with numerical simulations
- **Wigner–Eisenbud–Smith** delay time: follow peak of a wave packet;  
delay time can be much shorter than time for light propagation (“superluminal”)
- **Buttiker-Landauer** traversal time: modulating tunnel barrier height in time;  
predicts 450–560 as for present conditions

## II Free wave packet

**Schrödinger equation** in 1 dimension ( $\hbar = 1$ ):  $i\partial_t \psi(t, x) = \left( -\frac{\partial_x^2}{2m} + V(x) \right) \psi(t, x)$

**Free electron wave** (without barrier potential,  $V \equiv 0$ ):

single-momentum contribution (Fourier component):  $\tilde{\psi}_{\text{free}}(k; t, x) = \exp\left(i k x - i \frac{k^2}{2m} t\right)$

**Gaussian wave packet:** amplitude  $A(k) = \frac{1}{\sqrt[4]{2\pi}} \frac{1}{\sqrt{\sigma_k}} \exp\left(-\frac{(k - k_0)^2}{4\sigma_k^2} - ikx_0\right)$

↪ normalized momentum distribution  $|A(k)|^2 = \frac{1}{\sqrt{2\pi}} \frac{1}{\sigma_k} \exp\left(-\frac{(k - k_0)^2}{2\sigma_k^2}\right)$

with  $\langle k \rangle = \int_{-\infty}^{\infty} dk k |A(k)|^2 = k_0$  and  $\sigma_k = \sqrt{\langle k^2 \rangle - \langle k \rangle^2}$

**Time-dependent free wave packet:** integration of  $\tilde{\psi}_{\text{free}}$  over  $k$  with amplitudes  $A$

$$\psi_{\text{free}}(t, x) = \int_{-\infty}^{\infty} \frac{dk}{\sqrt{2\pi}} A(k) \tilde{\psi}_{\text{free}}(k; t, x)$$

↪ Fourier transformation with factor  $e^{ikx}$  included in  $\tilde{\psi}_{\text{free}}(k; t, x)$

## Time-dependent free wave packet: analytical solution

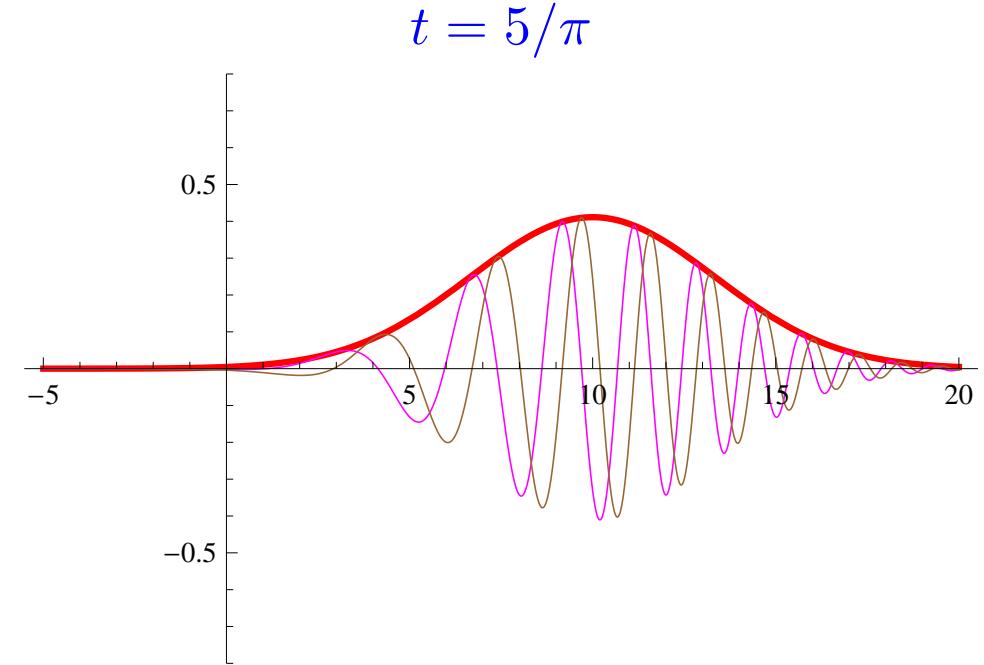
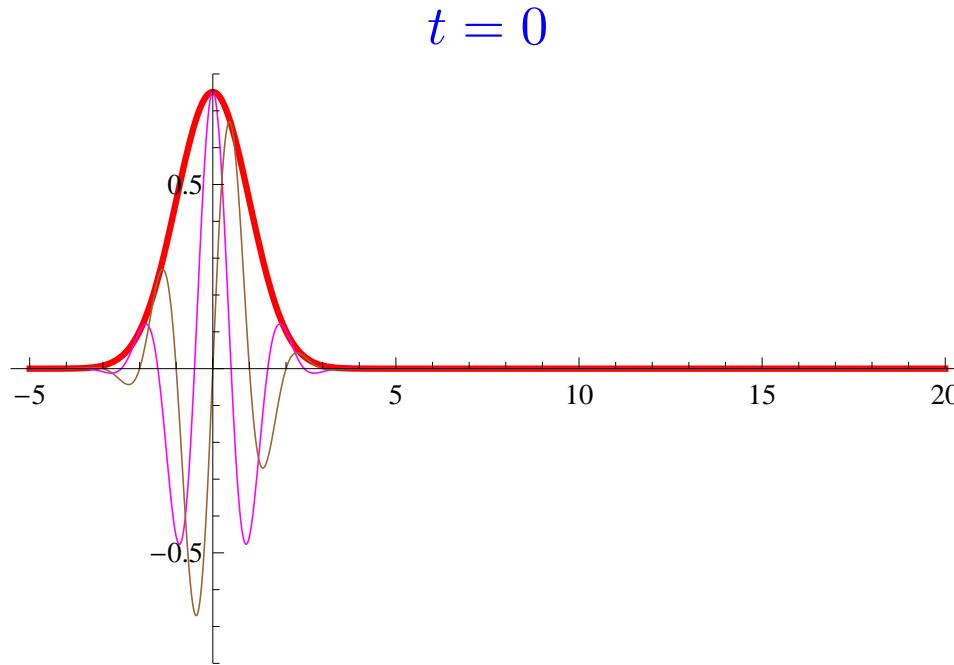
$$\psi_{\text{free}}(t, x) = \frac{1}{\sqrt[4]{2\pi}} \sqrt{\frac{2\sigma_k}{1 + i\frac{2\sigma_k^2}{m}t}} \exp\left(\frac{-\sigma_k^2(x - x_0)^2 + ik_0(x - x_0) - i\frac{k_0^2}{2m}t}{1 + i\frac{2\sigma_k^2}{m}t}\right)$$

↪ normalized probability distribution

$$|\psi_{\text{free}}(t, x)|^2 = \frac{1}{\sqrt{2\pi}} \frac{2\sigma_k}{\sqrt{1 + \frac{4\sigma_k^4}{m^2}t^2}} \exp\left(-\frac{1}{2} \frac{(2\sigma_k)^2}{1 + \frac{4\sigma_k^4}{m^2}t^2} \left[x - \left(x_0 + \frac{k_0}{m}t\right)\right]^2\right)$$

$$\text{with } \langle x \rangle = \int_{-\infty}^{\infty} dx x |\psi_{\text{free}}(t, x)|^2 = x_0 + \frac{k_0}{m}t \text{ and } \sigma_x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \frac{\sqrt{1 + \frac{4\sigma_k^4}{m^2}t^2}}{2\sigma_k}$$

**Free wave packet** with parameters  $k_0 = \pi$ ,  $\sigma_k = 1/\sqrt{2} \approx 0.7$ ,  $x_0 = 0$ ,  $2m = 1$ ,  
plot  $|\psi_{\text{free}}(t, x)|$ ,  $\text{Re } \psi_{\text{free}}(t, x)$  and  $\text{Im } \psi_{\text{free}}(t, x)$ :



Velocity of the packet maximum:  $v = \frac{k_0}{m} = 2k_0 = 2\pi$ ,  
but the front of the wave packet is faster!

Numerical integration of Fourier components  $A(k) \tilde{\psi}_{\text{free}}(k; t, x) \Rightarrow$  same result.

### III Stationary solution for rectangular tunnel barrier

**Rectangular tunnel barrier** with height  $V$  and length  $L$ :  $V(x) = V \theta(x)\theta(L - x)$

**Stationary solution** for incoming momentum  $k$ :

$$\tilde{\psi}(k; t, x) = \exp\left(-i\frac{k^2}{2m}t\right) \begin{cases} e^{ikx} + \frac{(k^2 + \kappa^2) \sinh(\kappa L)}{\Delta(k, \kappa)} e^{-ikx}, & x \leq 0, \\ \frac{k(k+i\kappa)}{\Delta(k, \kappa)} e^{-\kappa(x-L)} - \frac{k(k-i\kappa)}{\Delta(k, \kappa)} e^{\kappa(x-L)}, & 0 \leq x \leq L, \\ \frac{2ik\kappa}{\Delta(k, \kappa)} e^{ik(x-L)}, & L \leq x \end{cases}$$

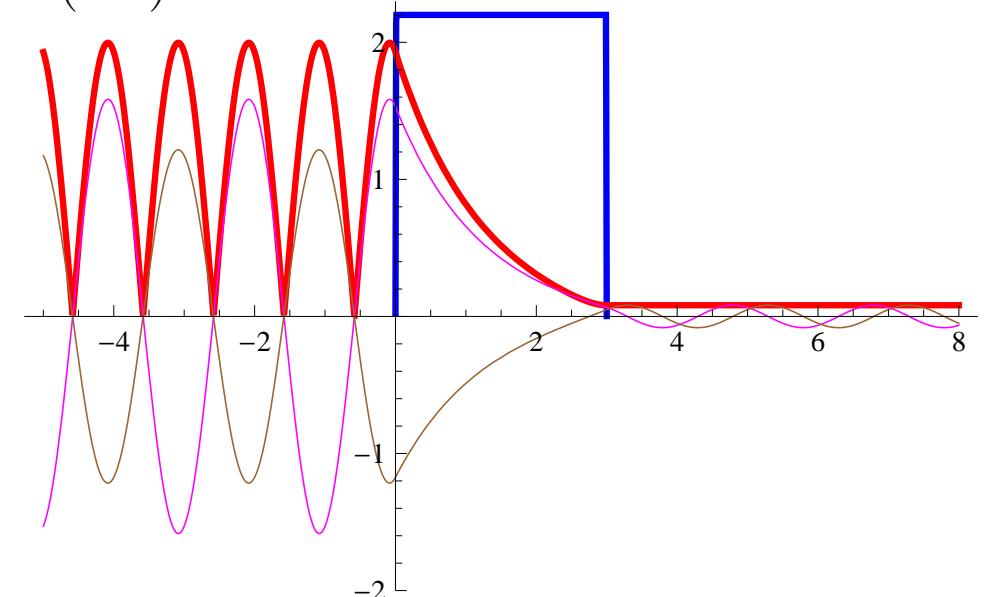
with  $\Delta(k, \kappa) = (k^2 - \kappa^2) \sinh(\kappa L) + 2ik\kappa \cosh(\kappa L)$

and  $\kappa = \sqrt{2mV - k^2}$

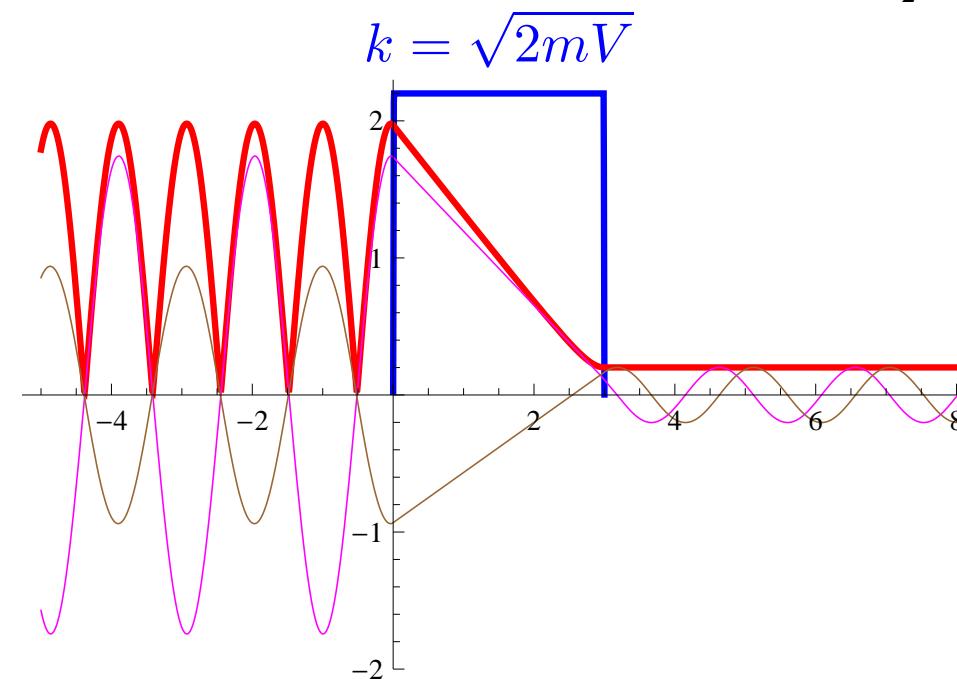
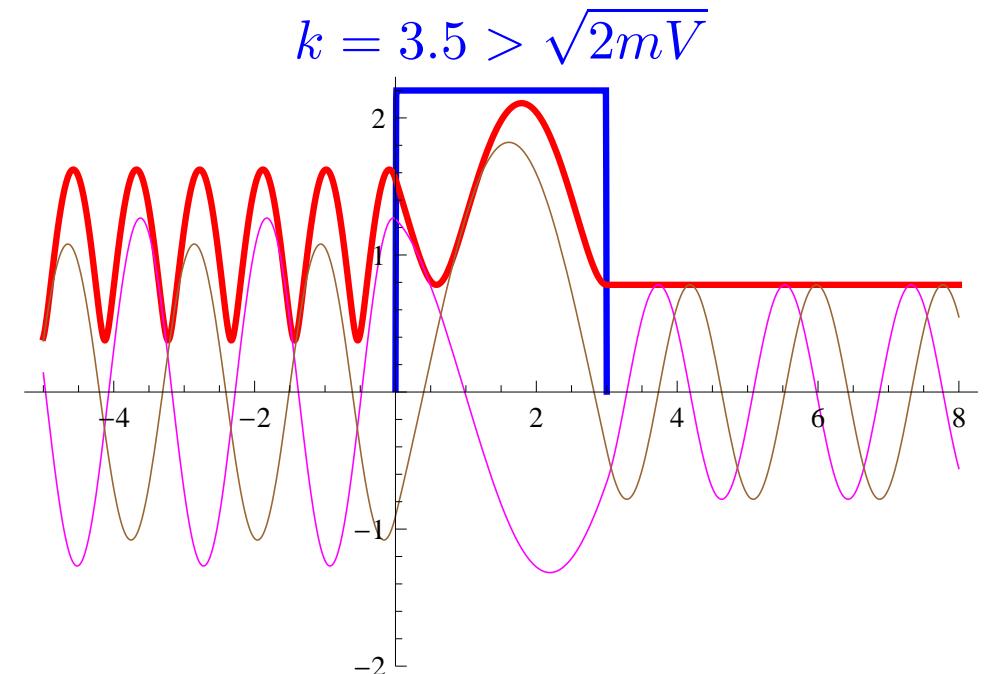
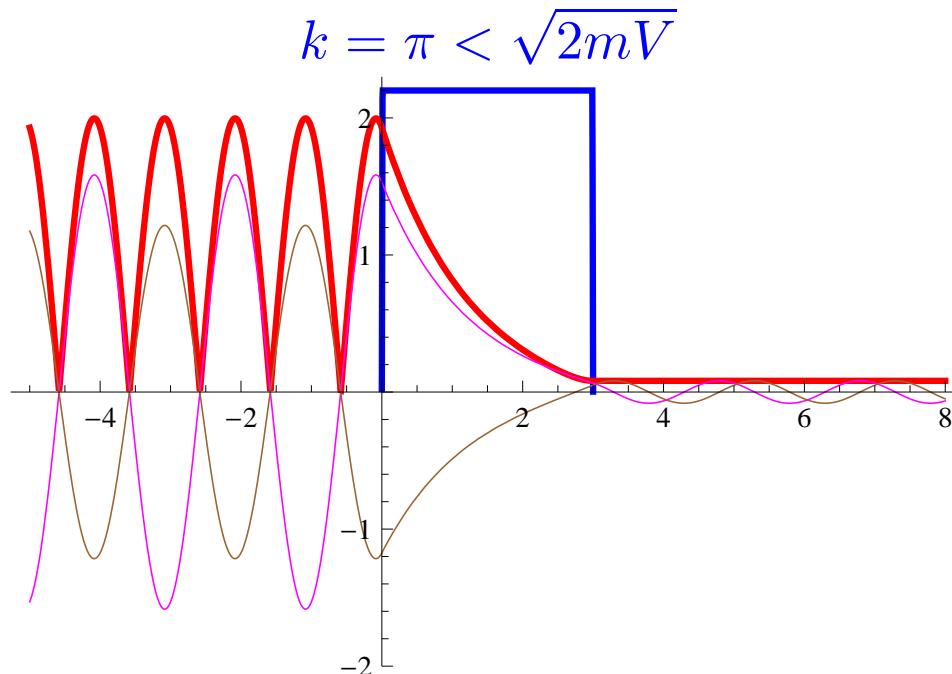
(or  $\kappa = -i\sqrt{k^2 - 2mV}$  for  $k^2 > 2mV$ )

**Example** for  $k^2 < 2mV$ :

plot barrier (height in arbitrary units),  
 $|\tilde{\psi}(k; t, x)|$ ,  $\text{Re } \tilde{\psi}(k; t, x)$  and  $\text{Im } \tilde{\psi}(k; t, x)$



**Stationary solution** for  $L = 3$  and  $\sqrt{2mV} = 3.25$ :



## IV Numerical solution for time-dependent tunneling

Multiply stationary solution  $\tilde{\psi}(k; t, x)$  with Gaussian wave packet  $A(k)$  and integrate numerically over  $k$ :

$$\psi(t, x) = \int_{-\infty}^{\infty} \frac{dk}{\sqrt{2\pi}} A(k) \tilde{\psi}(k; t, x)$$

**Numerical integration** performed by MATHEMATICA with a working precision of 20–35 digits, depending on the parameters.

### Three cases:

- A: “slow” wave packet ( $k_0 = \pi \gtrsim \sigma_k$ ), thin barrier ( $L = 1 = \pi/k_0 \sim \sigma_x$ )
- B: “slow” wave packet ( $k_0 = \pi \gtrsim \sigma_k$ ), thick barrier ( $L = 5 > \{\sigma_x, \pi/k_0\}$ )
- C: “fast” wave packet ( $k_0 = 20 \gg \sigma_k$ ), thick barrier ( $L = 5 > \sigma_x \gg \pi/k_0$ )

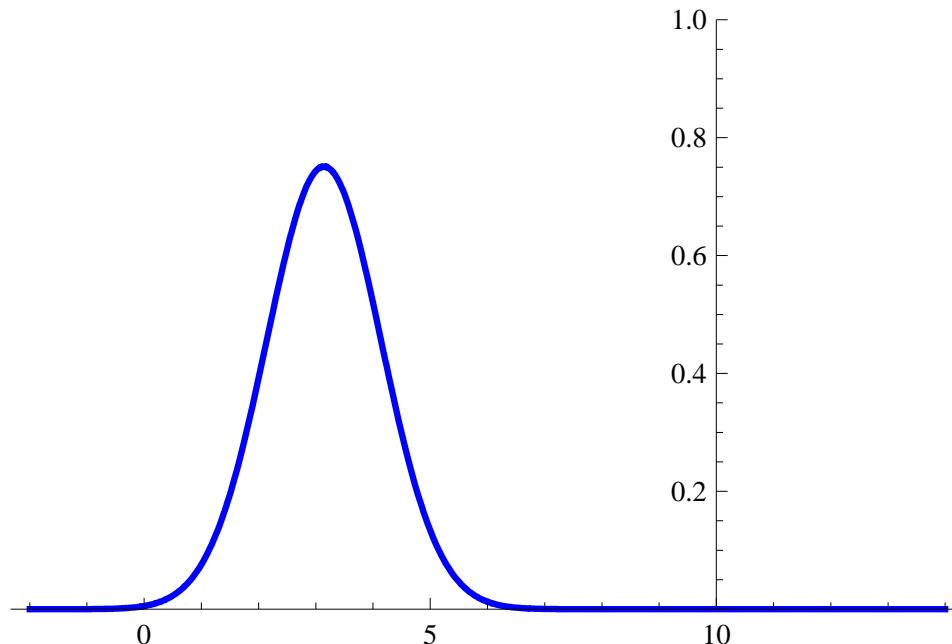
## A: “Slow” wave packet, thin barrier

**Parameters:**  $k_0 = \pi \gtrsim \sigma_k$ ,

$L = 1 = \pi/k_0 \sim \sigma_x$ ,  $\sigma_k = 1/\sqrt{2} \approx 0.7$ ,

$x_0 = -5$ ,  $2m = 1$ ,  $V = 10^2 \gg k_0^2$

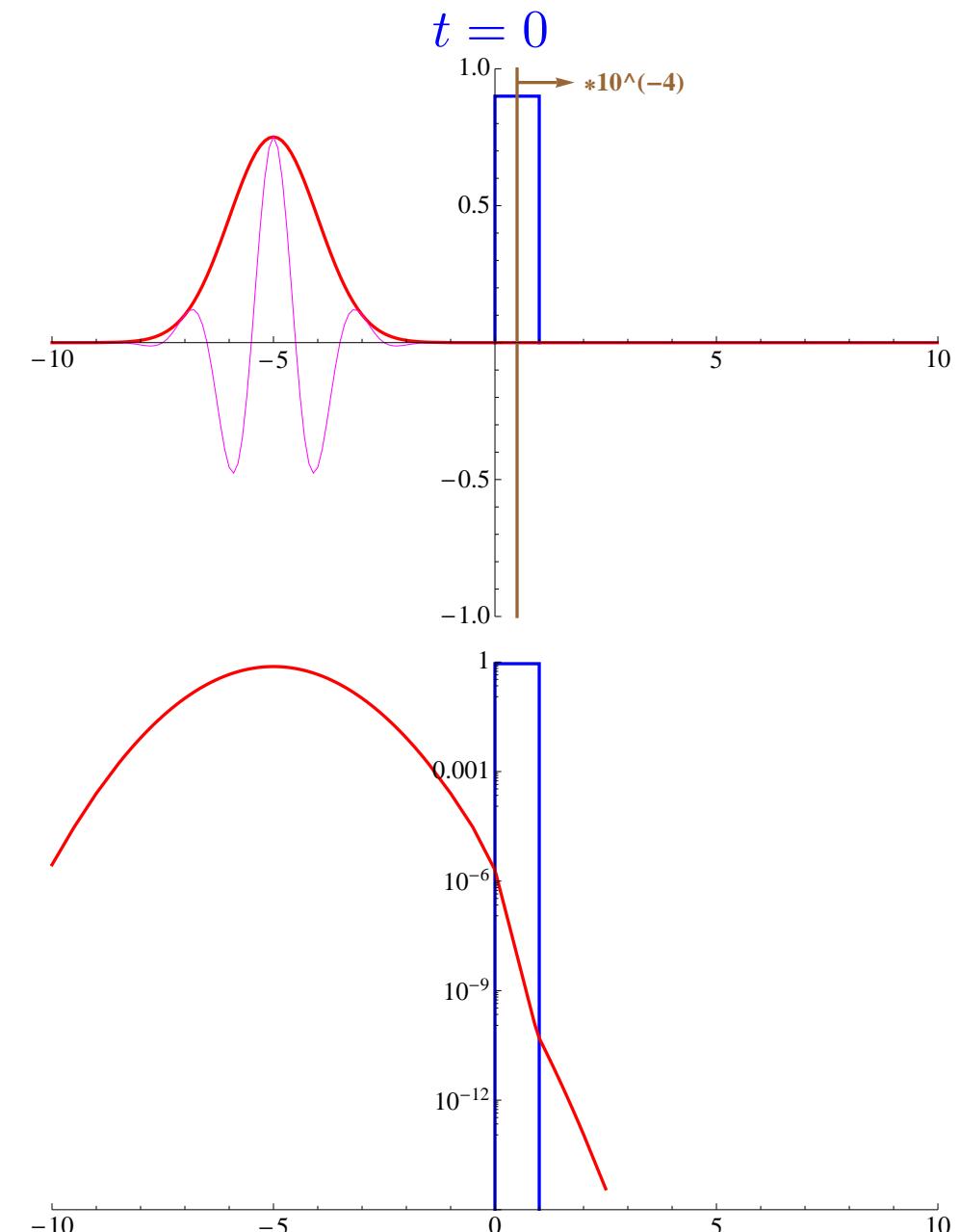
**Momentum distribution**  $|A(k)|$ :



**Right plots:** barrier,  $|\psi(t, x)|$ ,  $\text{Re } \psi(t, x)$ ;

top: to scale by  $10^{-4}$  for  $x > 0.5$ ;

bottom: logarithmic scale

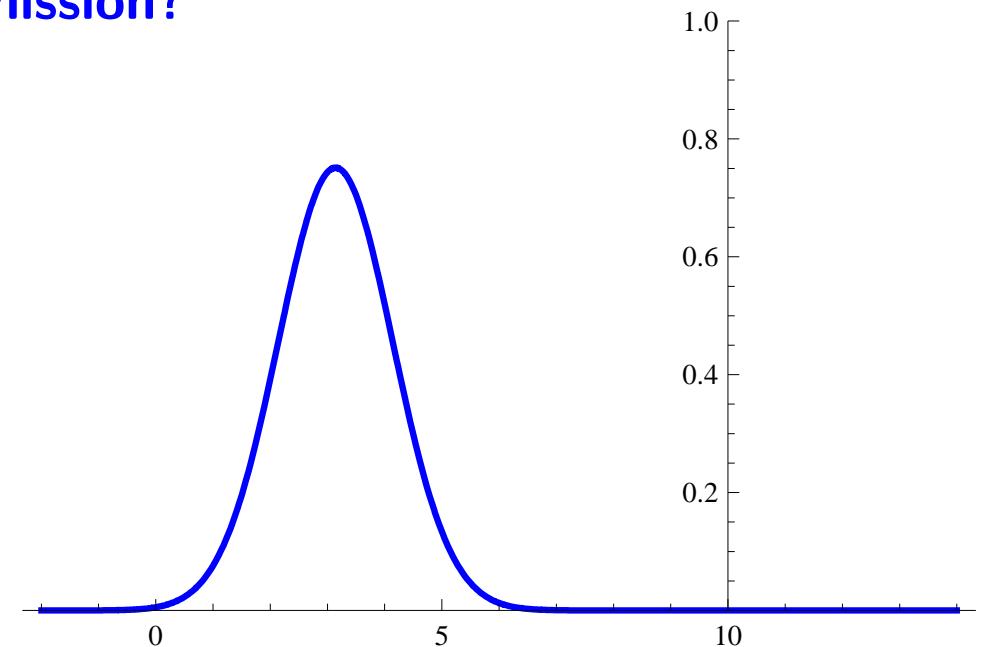


## A: “Slow” wave packet, thin barrier — Which momenta contribute to the transmission?

Momentum distribution of incoming wave

packet, i.e. part of  $\tilde{\psi}(k; t, x) \propto e^{ikx}$  at

$x < 0$ :  $|A(k)|$



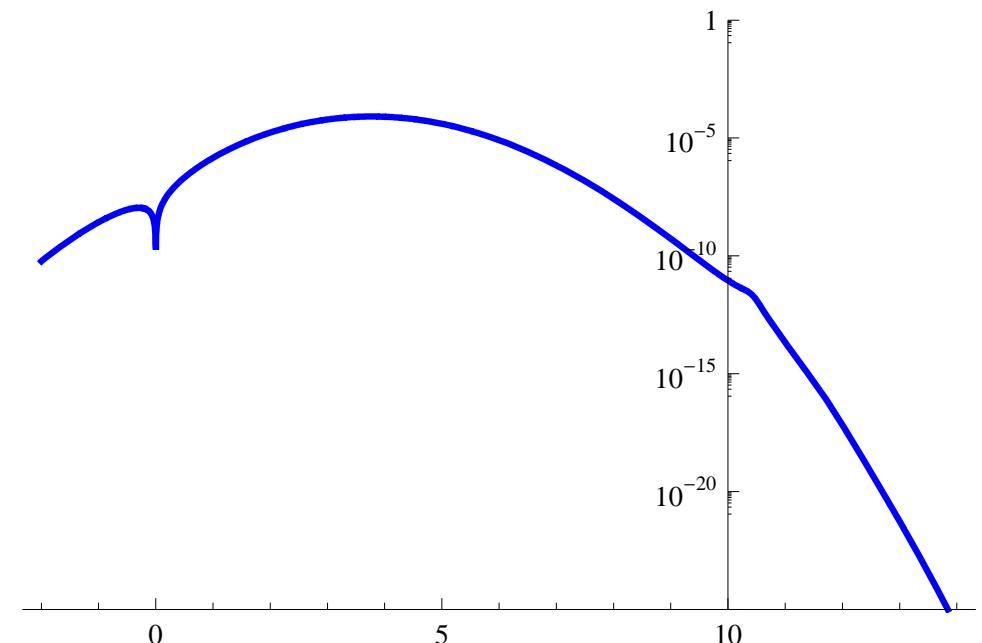
Momentum distribution of transmitted

wave packet, i.e.  $\tilde{\psi}(k; t, L)$  at  $x > L$ :

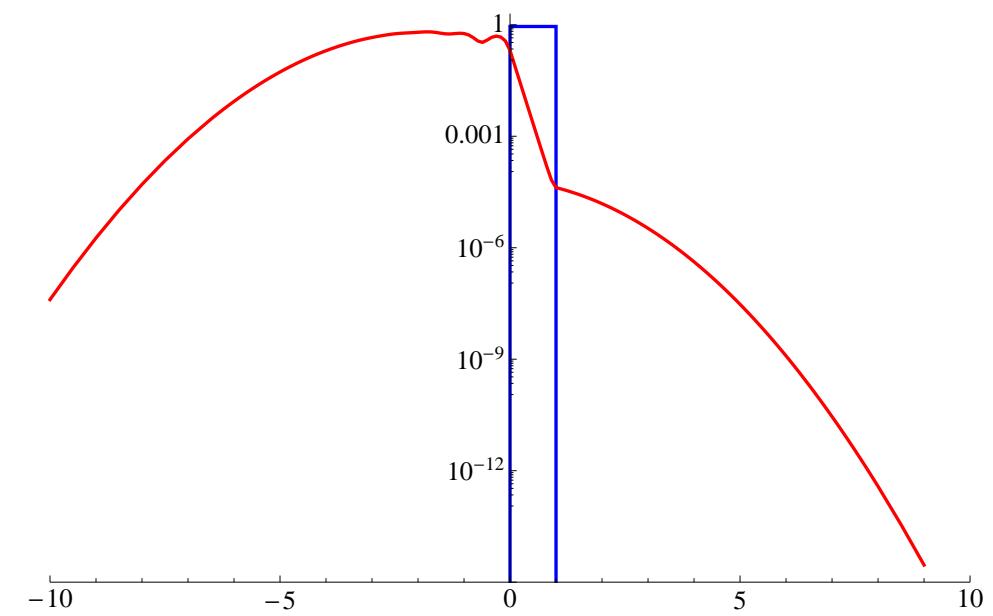
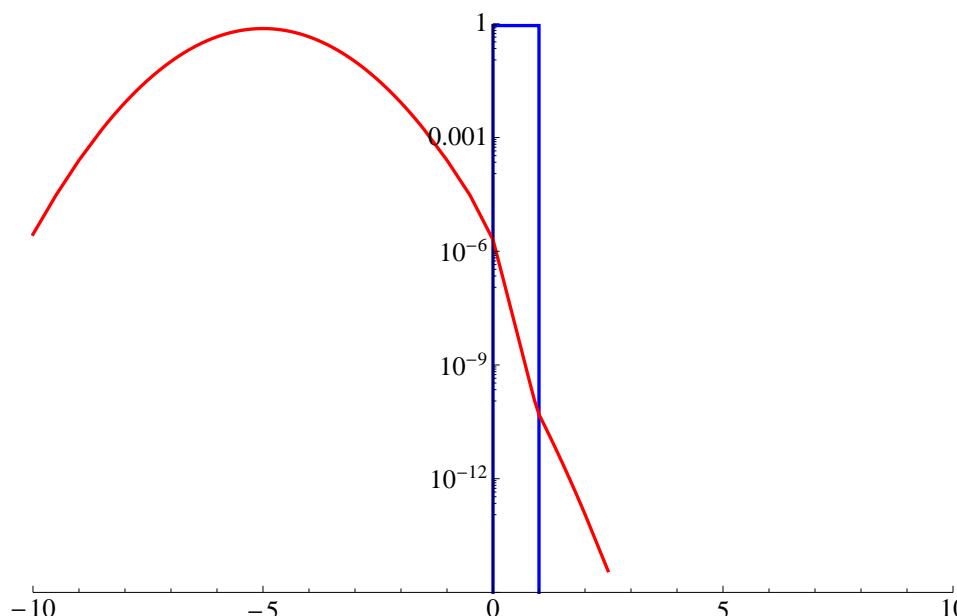
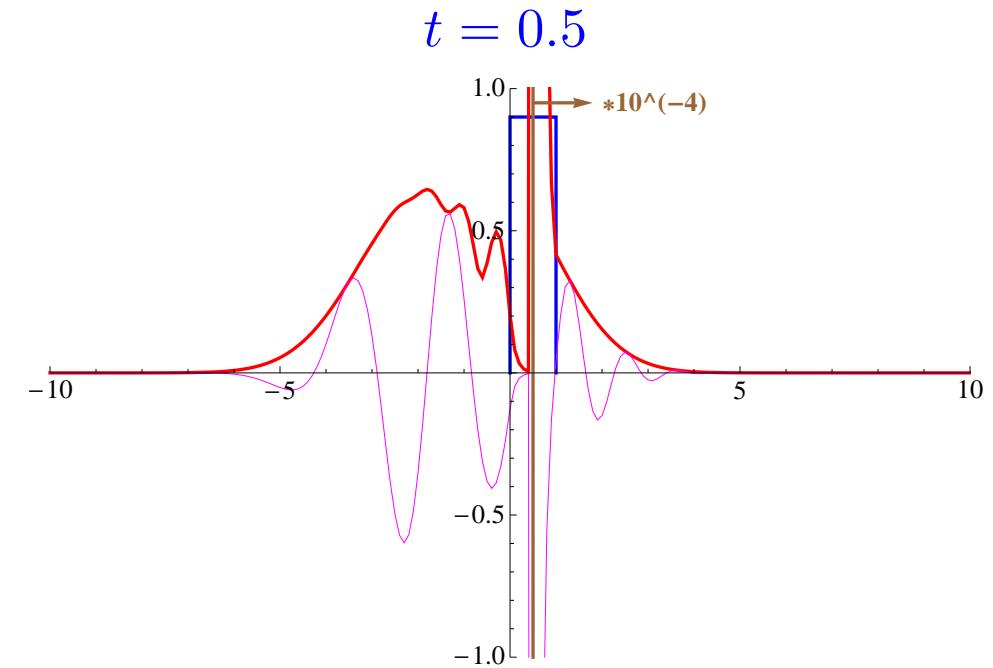
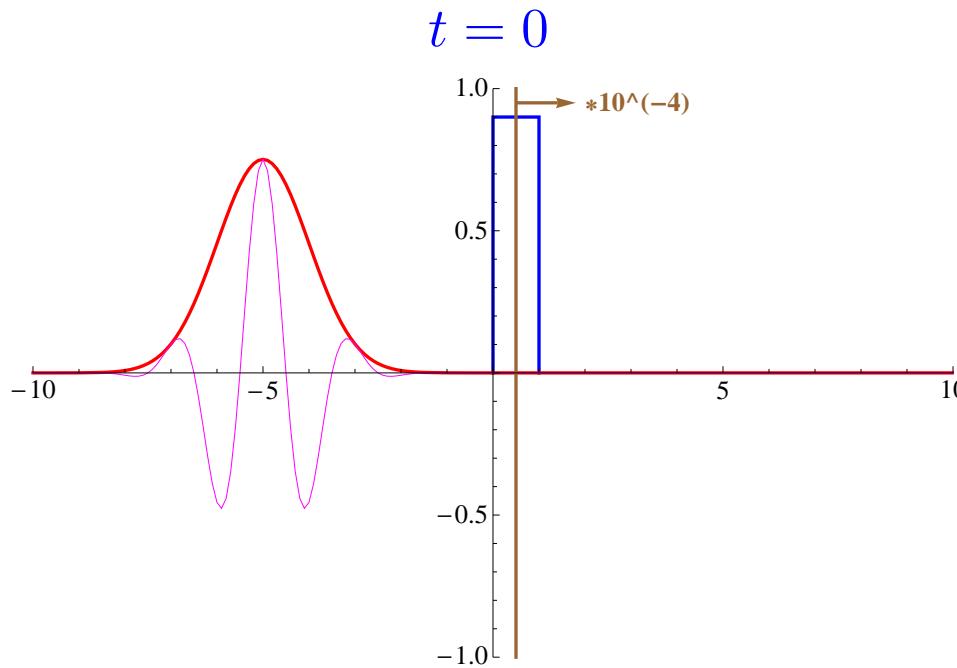
$|A(k) \tilde{\psi}(k; t, L)|$  (logarithmic scale)

⇒ transmission dominated by  
central momenta  $k \approx k_0 < \sqrt{V}$

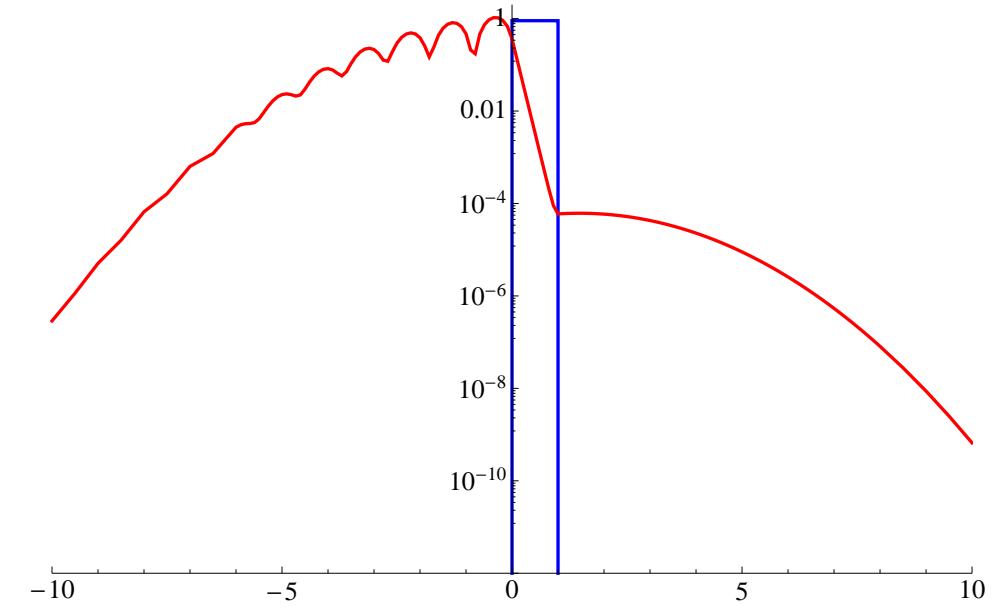
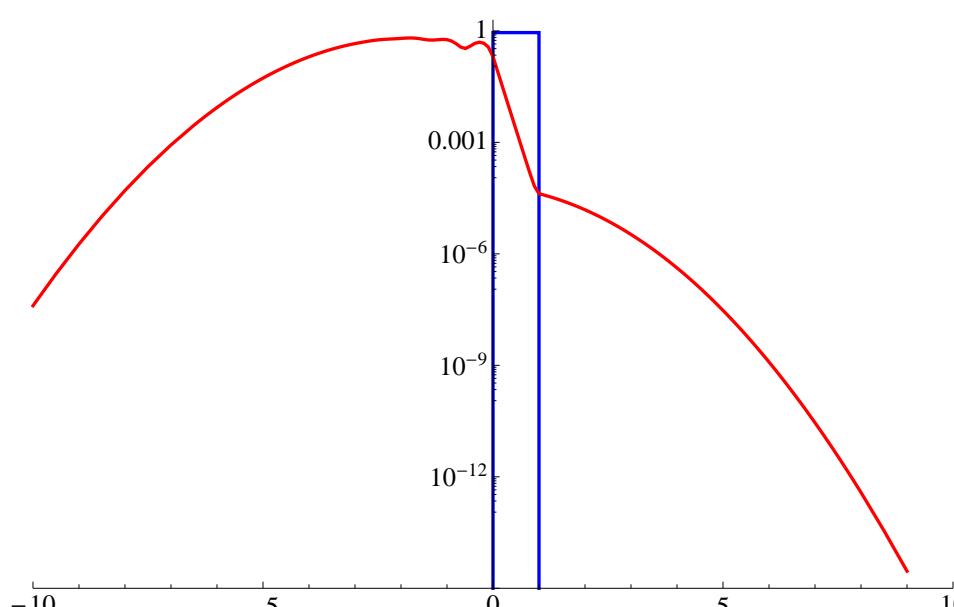
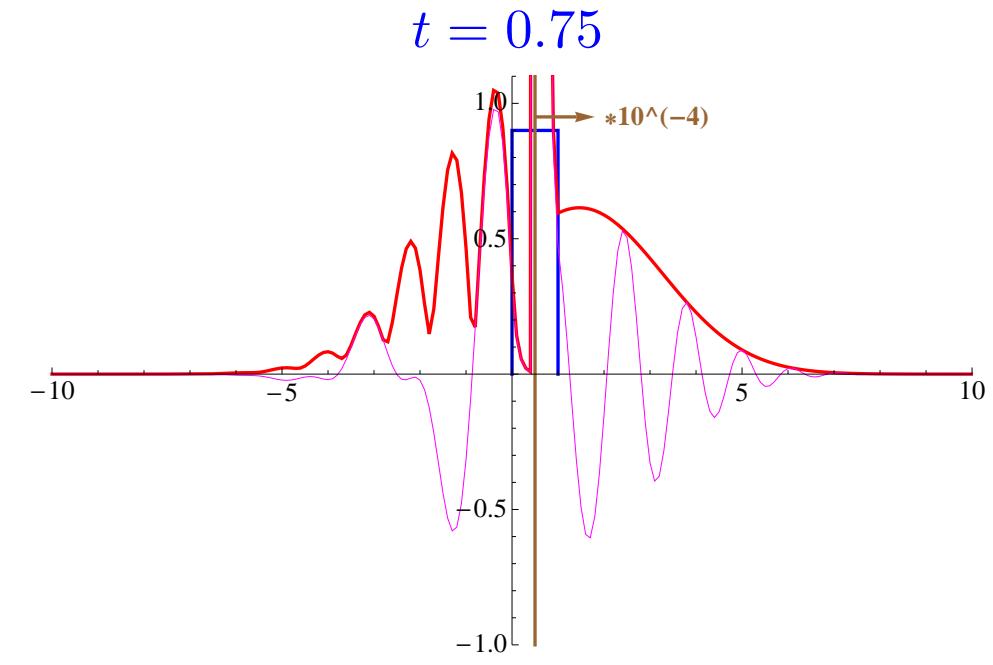
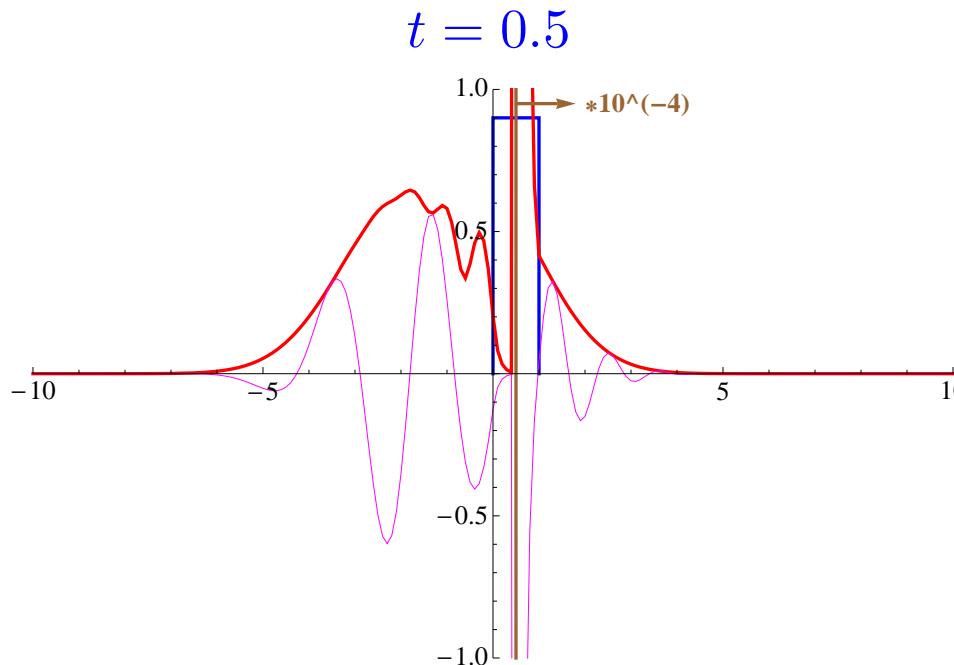
⇒ “real” tunneling



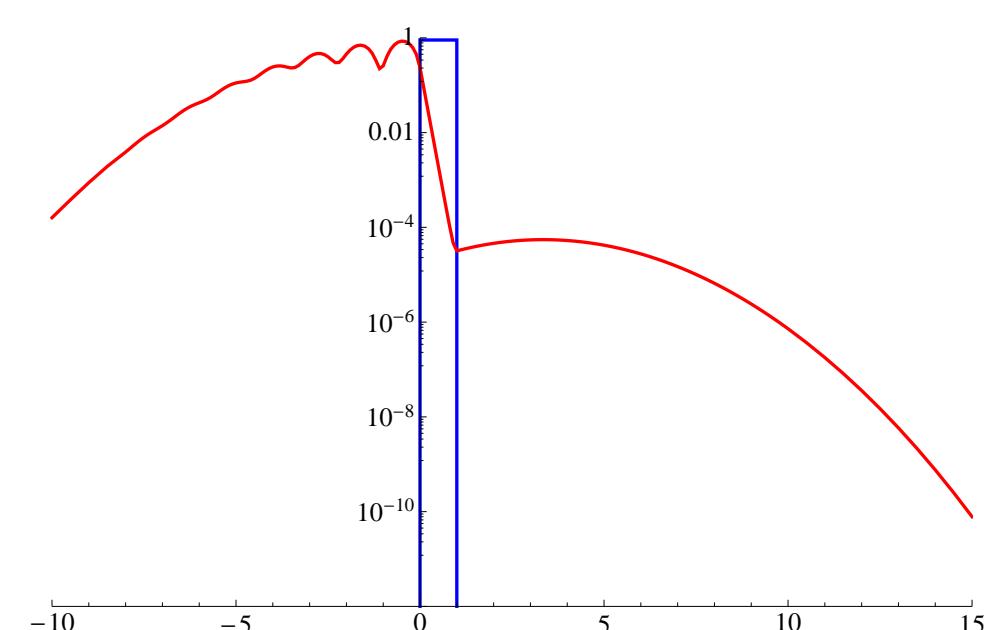
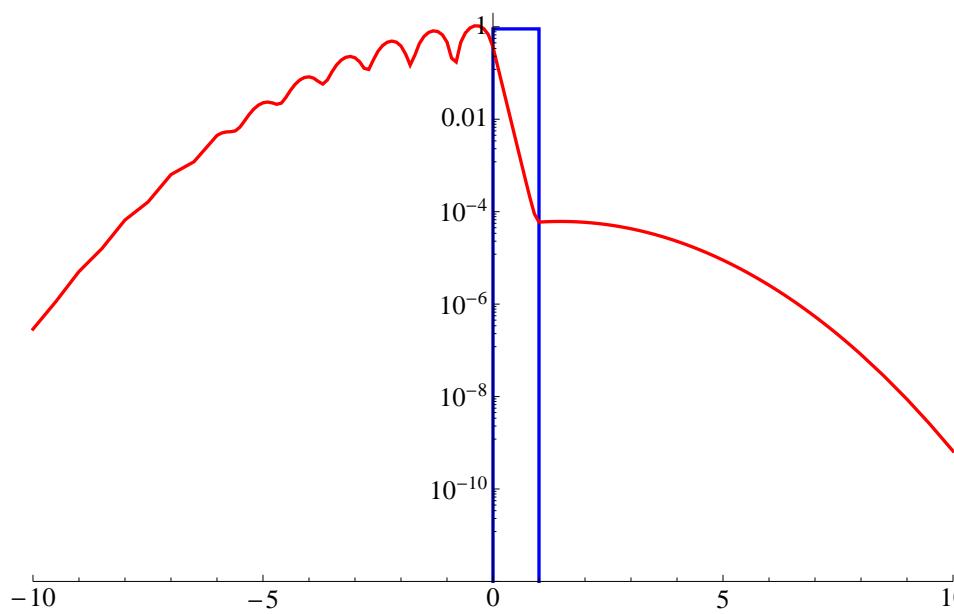
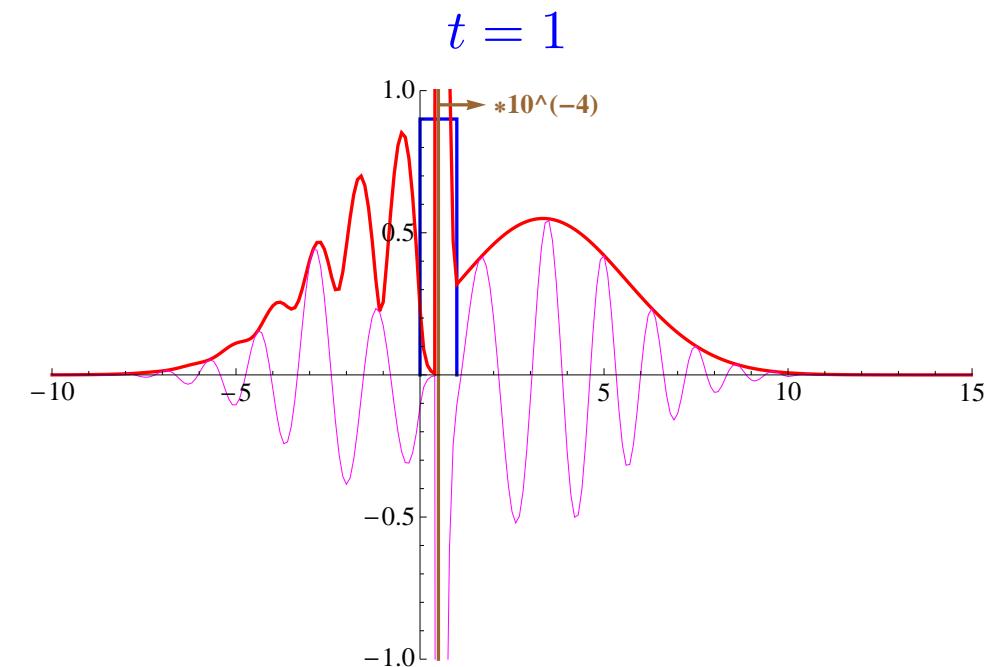
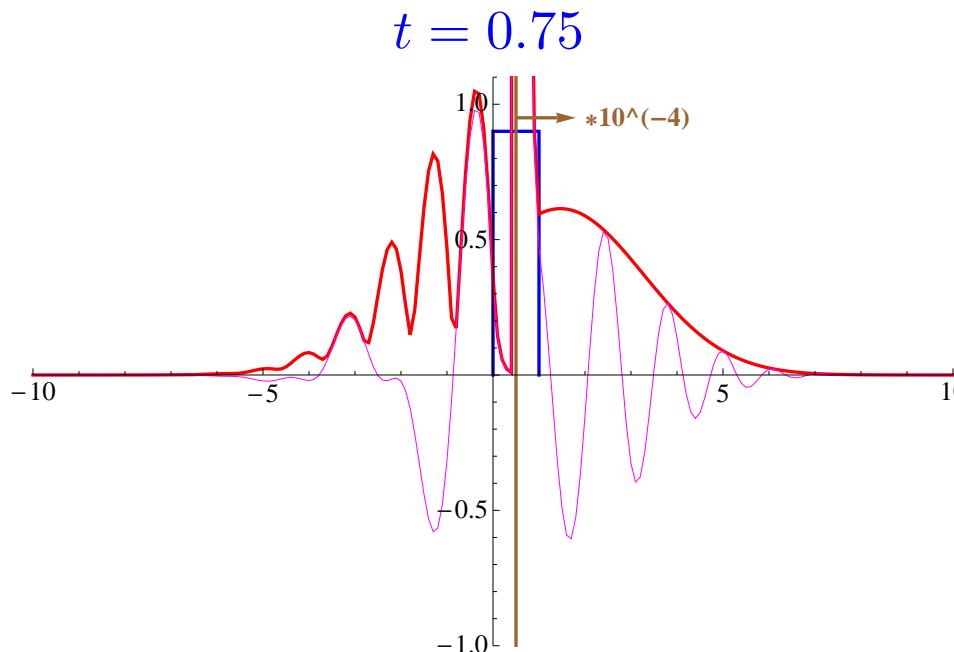
## A: “Slow” wave packet, thin barrier



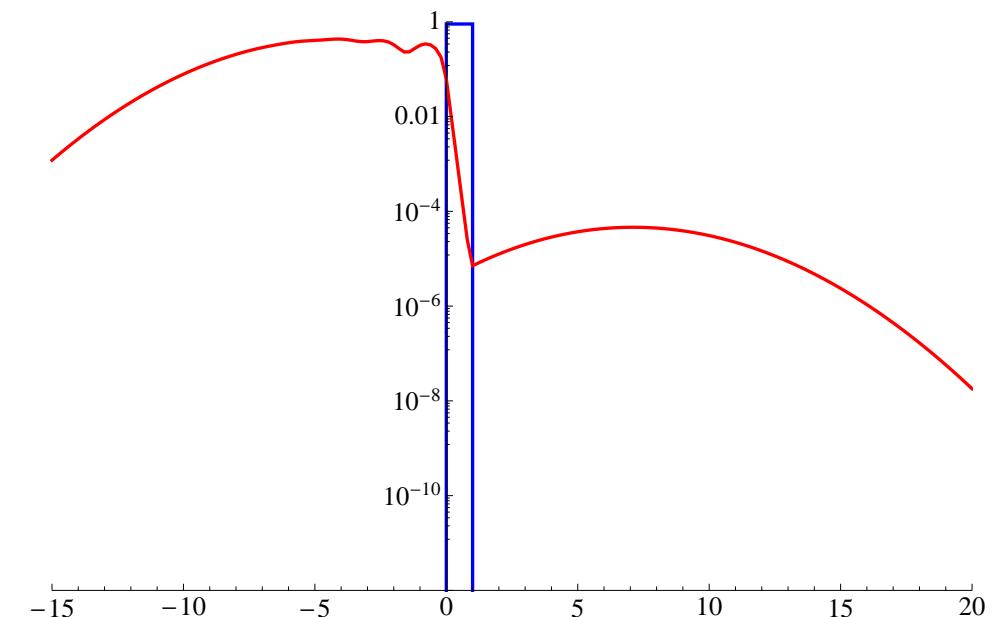
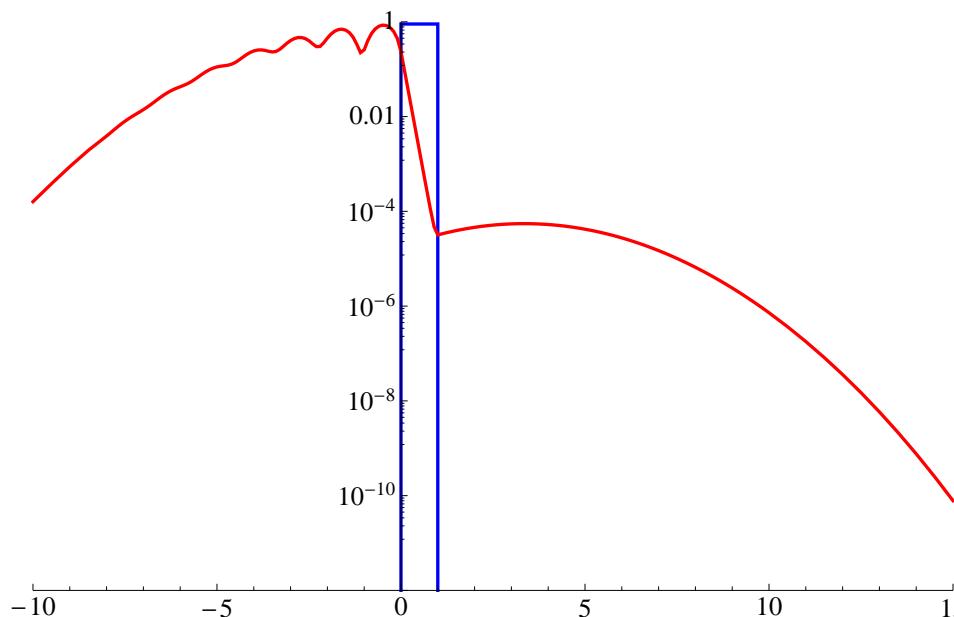
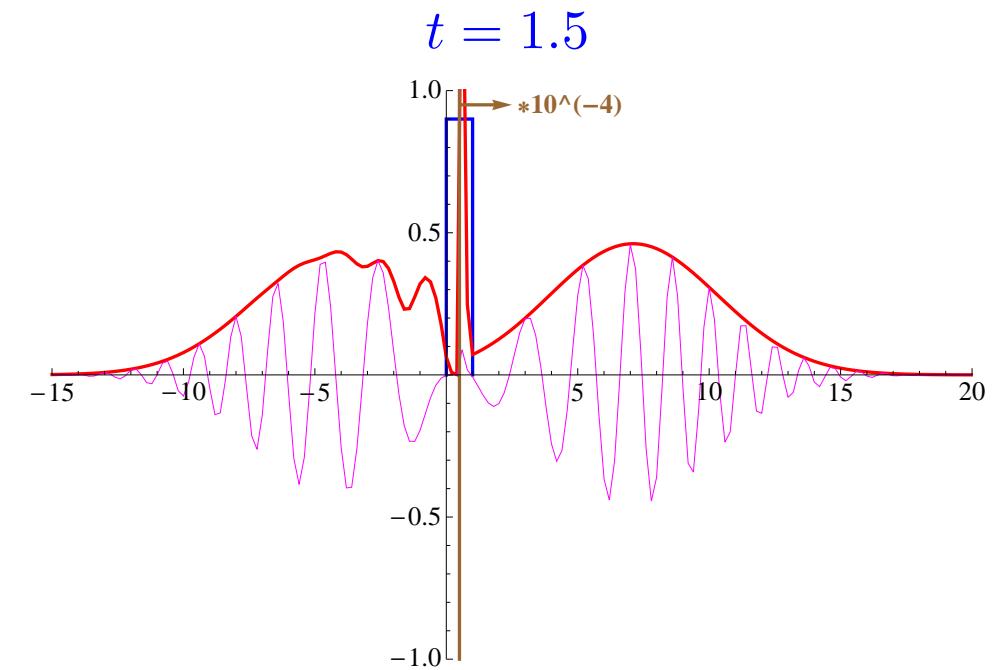
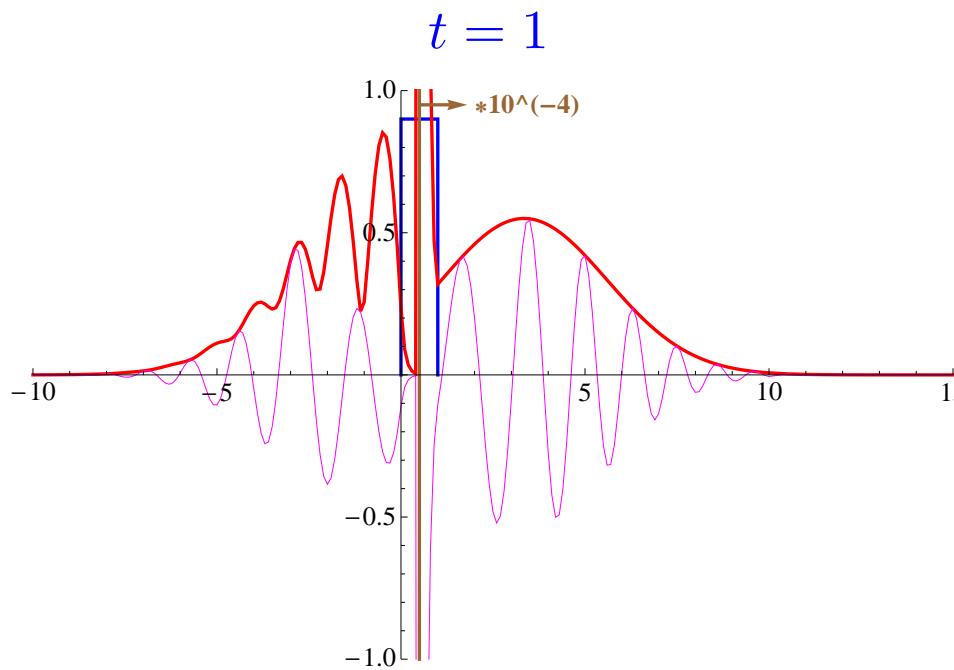
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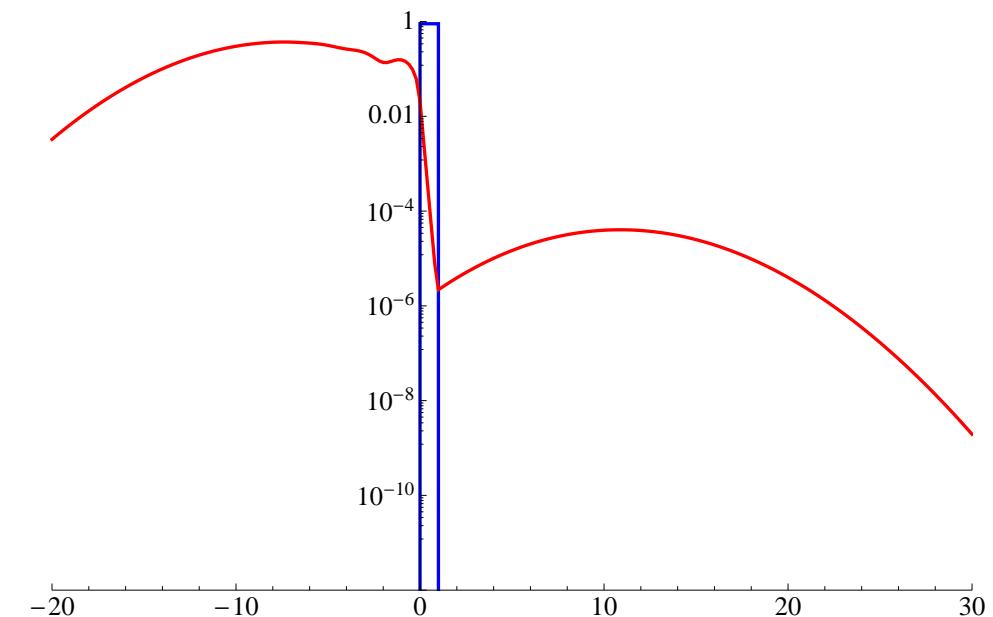
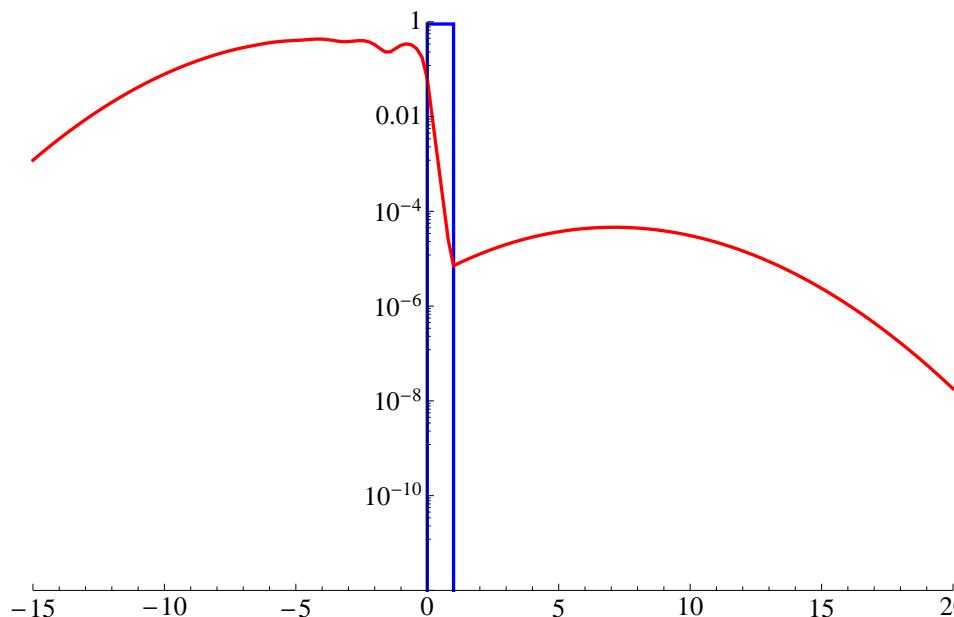
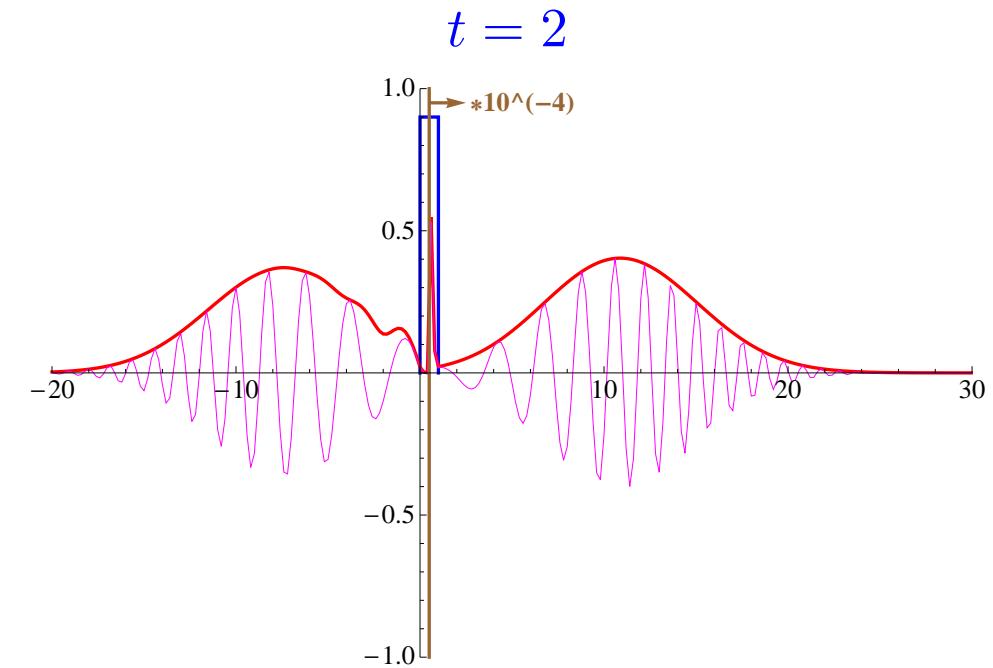
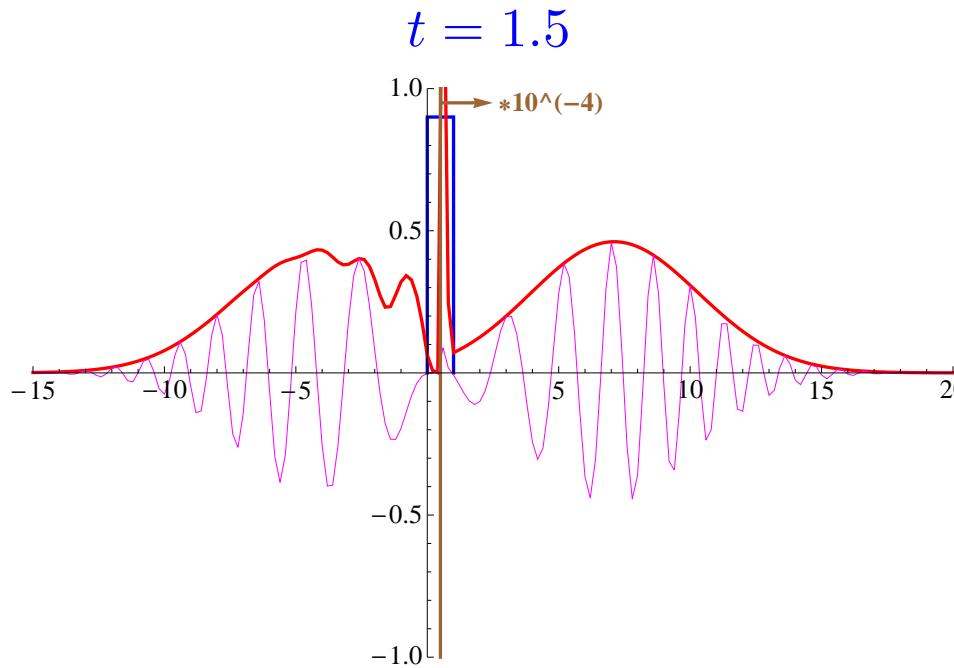
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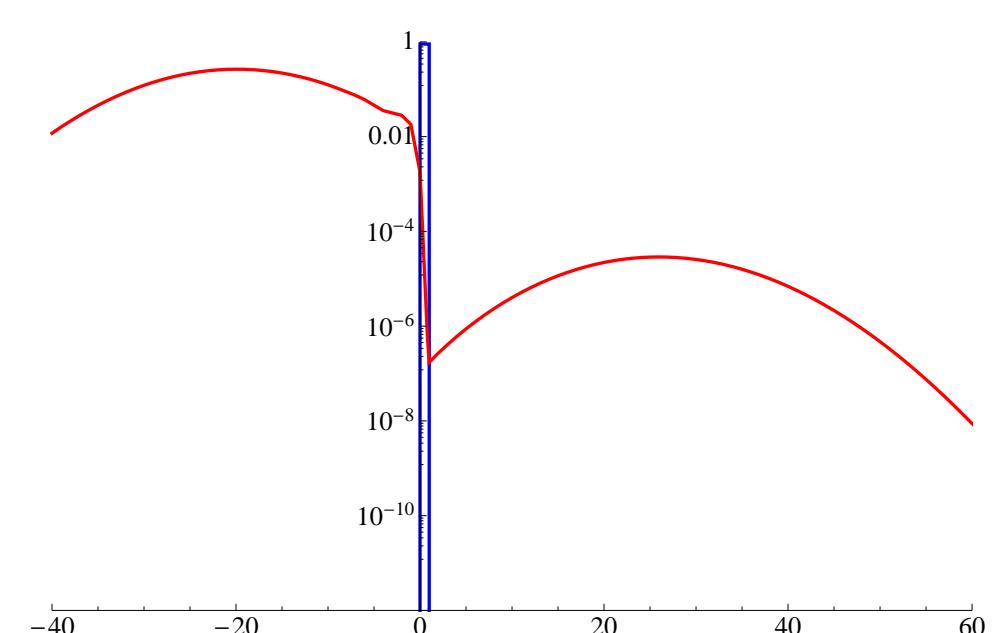
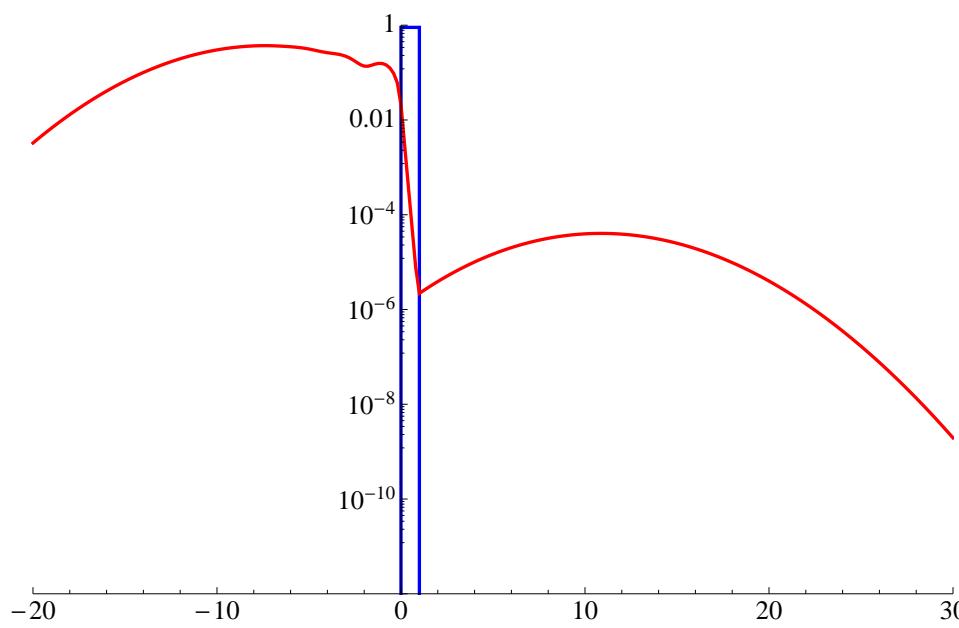
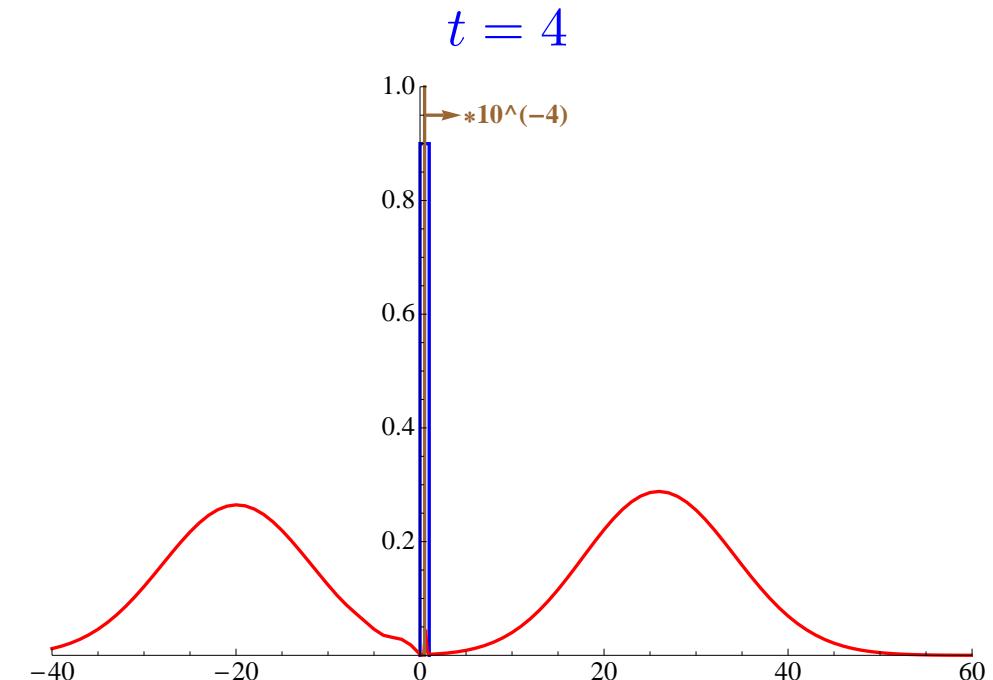
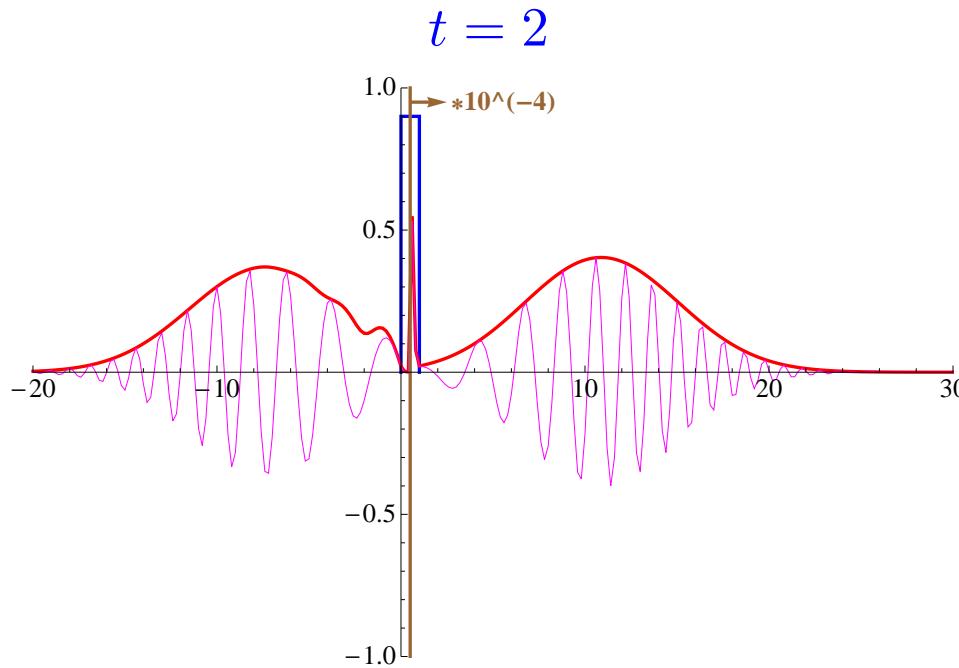


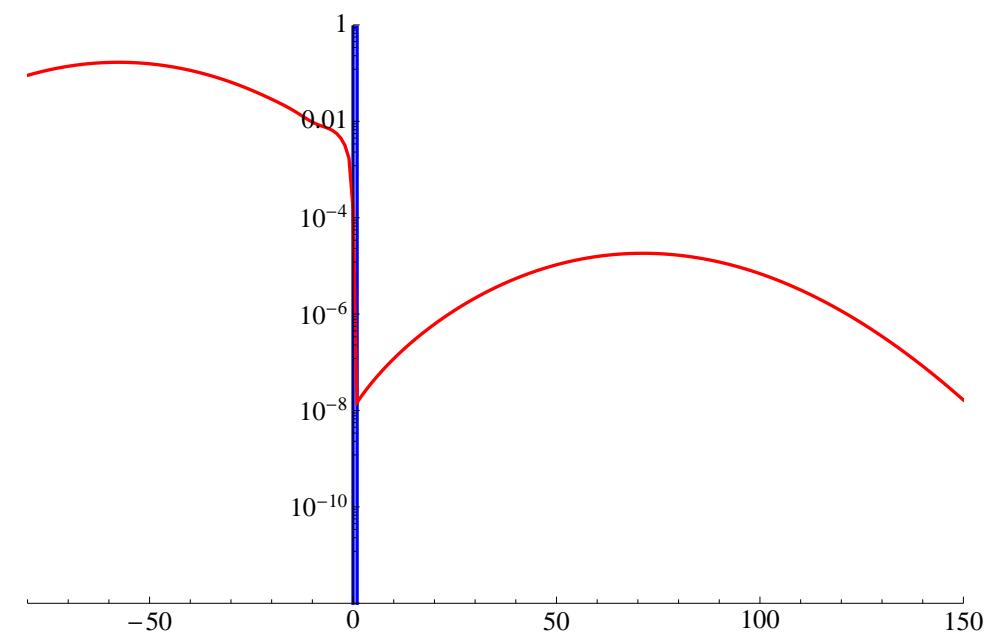
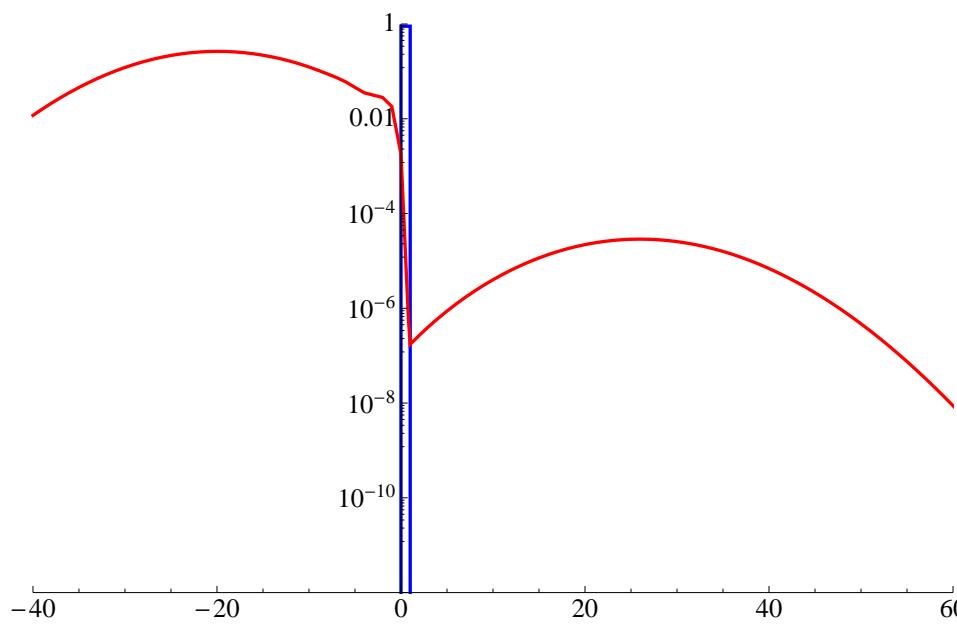
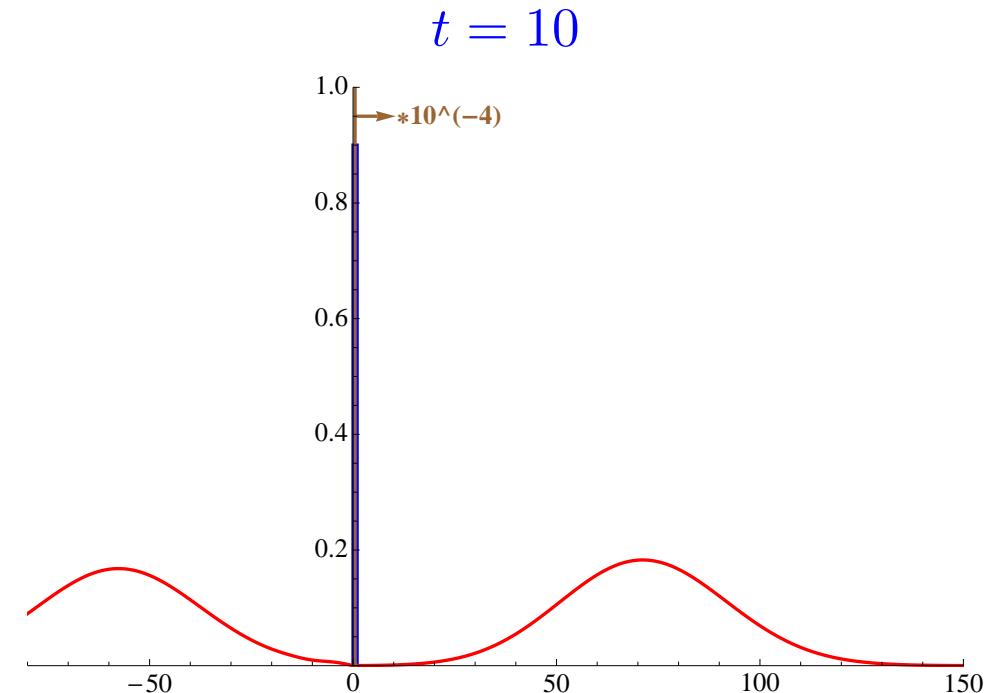
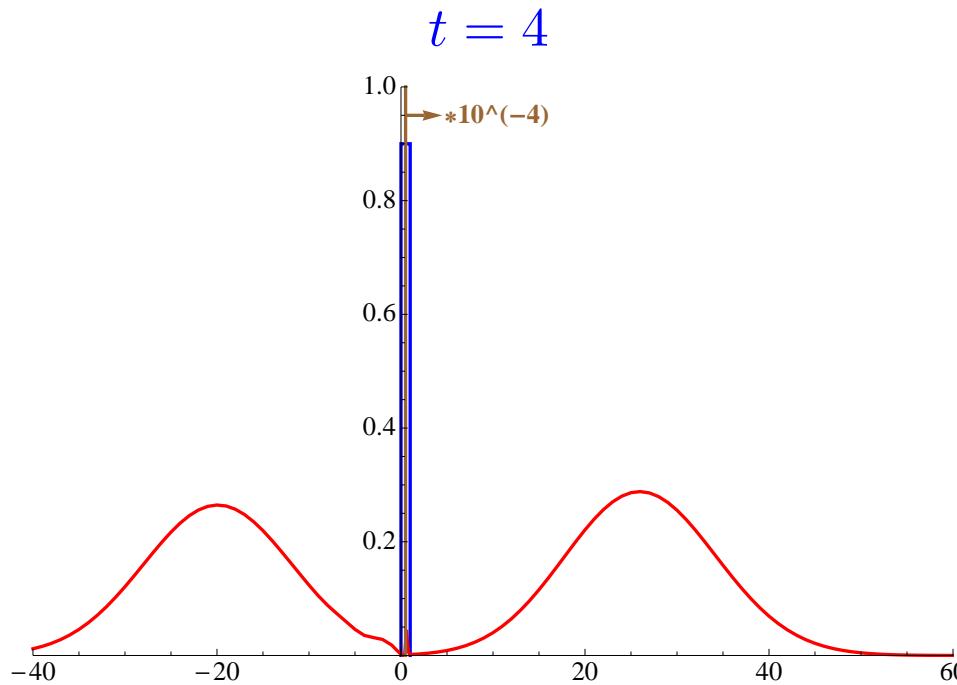
## A: “Slow” wave packet, thin barrier



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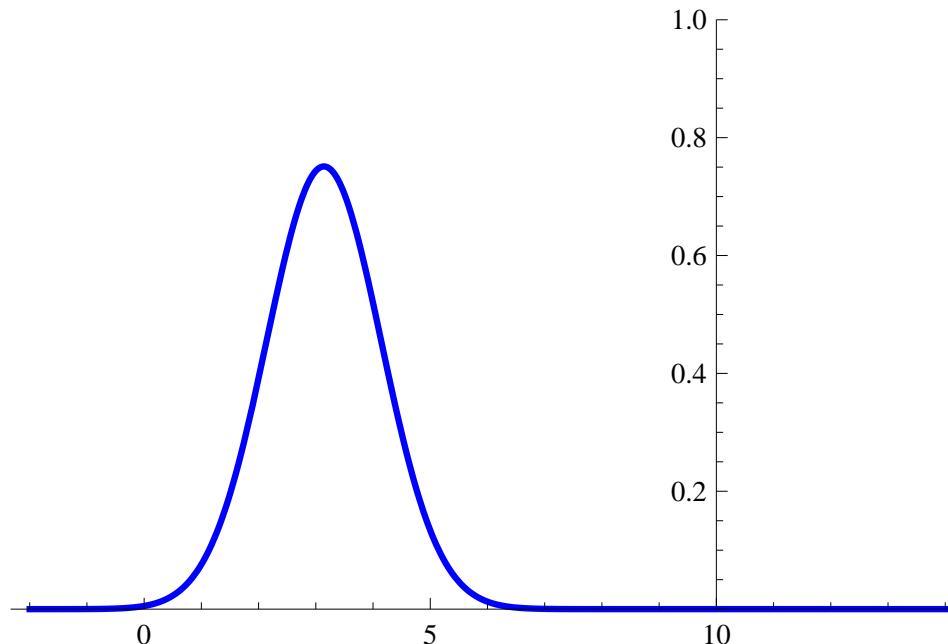
## A: “Slow” wave packet, thin barrier — Conclusion

- Thin barrier: only modest exponential damping through the tunneling ( $\sim 10^{-4}$ )  
↪ transmission dominated by **tunneling of low-momentum components**.
- Incoming and reflected wave packet: wave length ( $2\pi/k$ ) and velocity of maximum ( $2k$ ) correspond to  $k_0 = \pi$  and to a **reflection at  $x = 0$** .
- In the tunnel: exponential decrease, no structure
- Transmitted wave packet: wave length  $\rightsquigarrow k \approx 4$ , velocity of maximum  $\rightsquigarrow 3.5 \lesssim k \lesssim 4 \Rightarrow k \gtrsim k_0$ ,  $k^2 \ll V$
- Packet maximum enters tunnel at  $t = 0.8$  and exits at  $t = 0.7$   
 $\Rightarrow \Delta t_{\text{tunnel}} \approx -0.1 \lesssim 0$
- Reflected and transmitted wave packets are approximately Gaussian.  
Only while overlapping with the barrier, the incoming/reflected wave packet develops **fringes** (“standing wave”).

## B: “Slow” wave packet, thick barrier

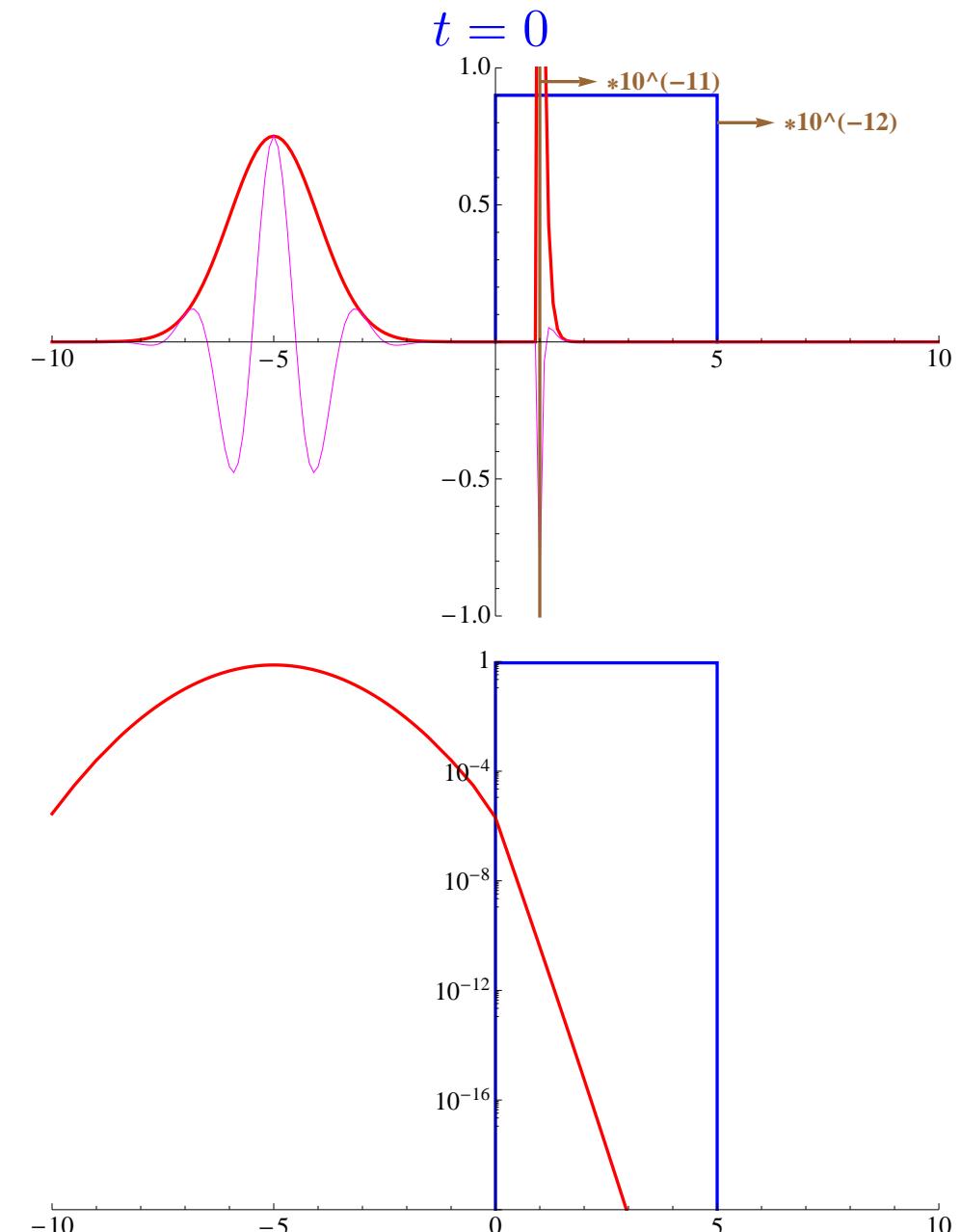
**Parameters:**  $k_0 = \pi \gtrsim \sigma_k$ ,  
 $L = 5 > \{\sigma_x, \pi/k_0\}$ ,  $\sigma_k = 1/\sqrt{2} \approx 0.7$ ,  
 $x_0 = -5$ ,  $2m = 1$ ,  $V = 10^2 \gg k_0^2$

**Momentum distribution**  $|A(k)|$ :



**Right plots:** barrier,  $|\psi(t, x)|$ ,  $\text{Re } \psi(t, x)$ ;

top: to scale by  $10^{-11}$  for  $1 < x < L$   
and by  $10^{-12}$  for  $x > L$ ;

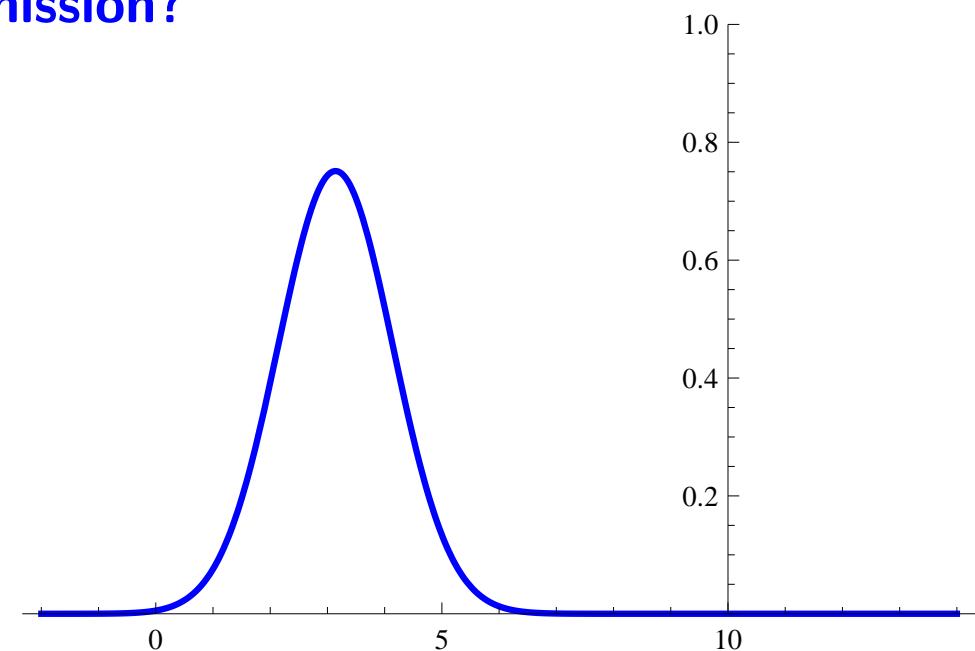


## B: “Slow” wave packet, thick barrier — Which momenta contribute to the transmission?

Momentum distribution of incoming wave

packet, i.e. part of  $\tilde{\psi}(k; t, x) \propto e^{ikx}$  at

$x < 0$ :  $|A(k)|$



Momentum distribution of transmitted

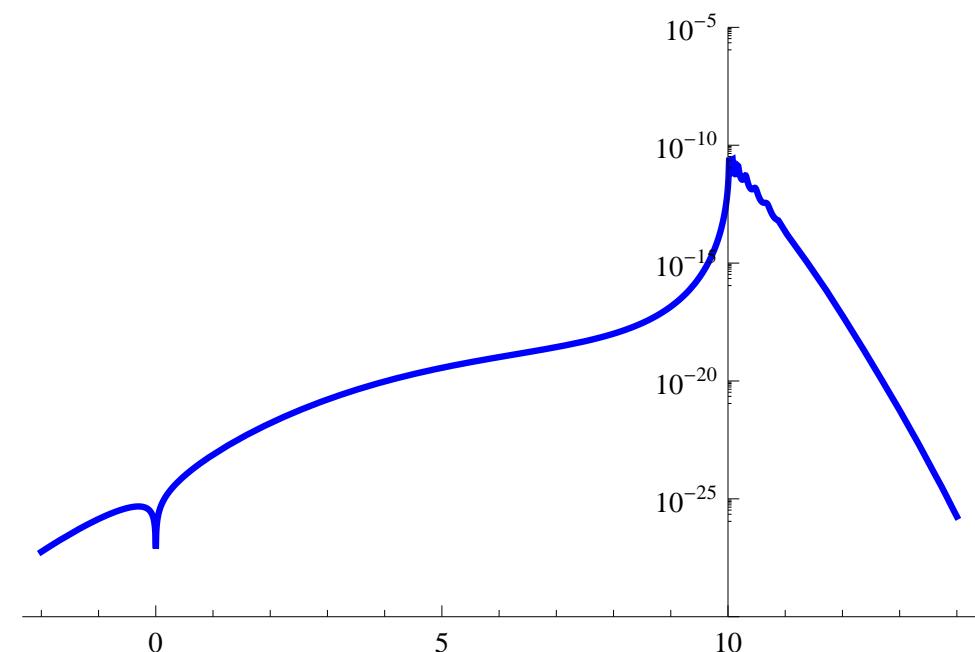
wave packet, i.e.  $\tilde{\psi}(k; t, L)$  at  $x > L$ :

$|A(k) \tilde{\psi}(k; t, L)|$  (logarithmic scale)

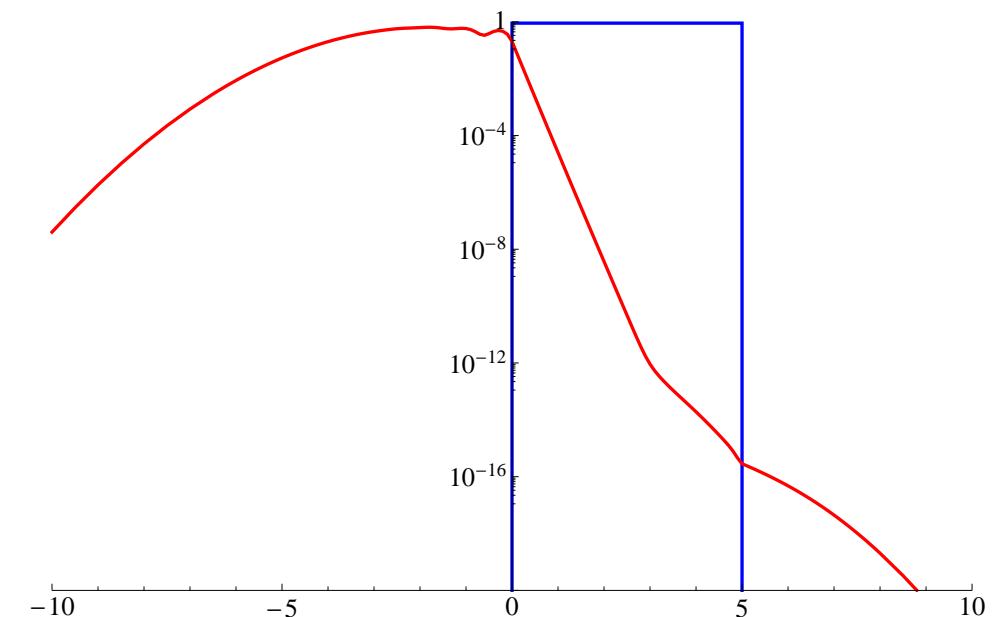
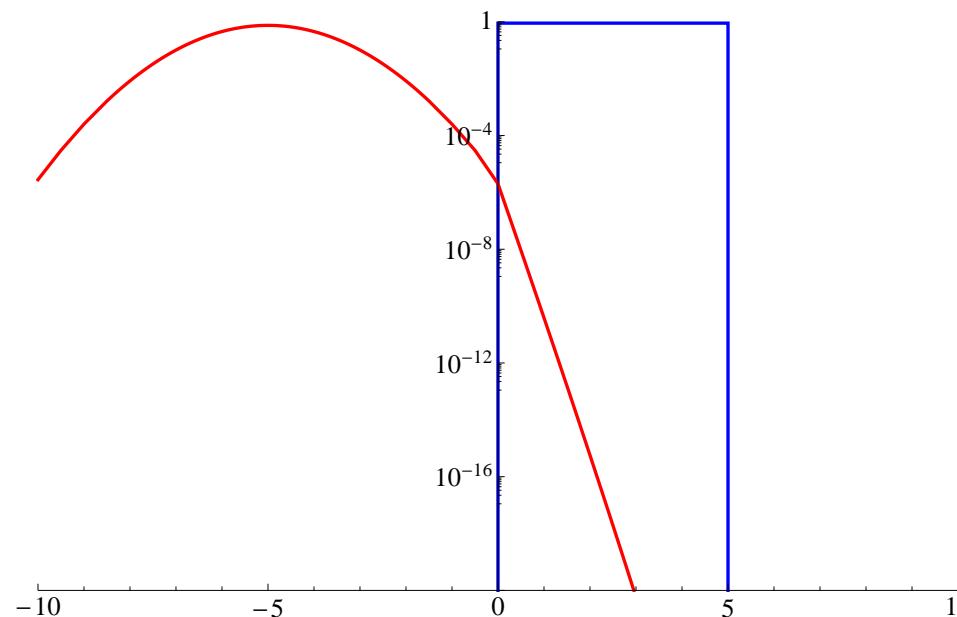
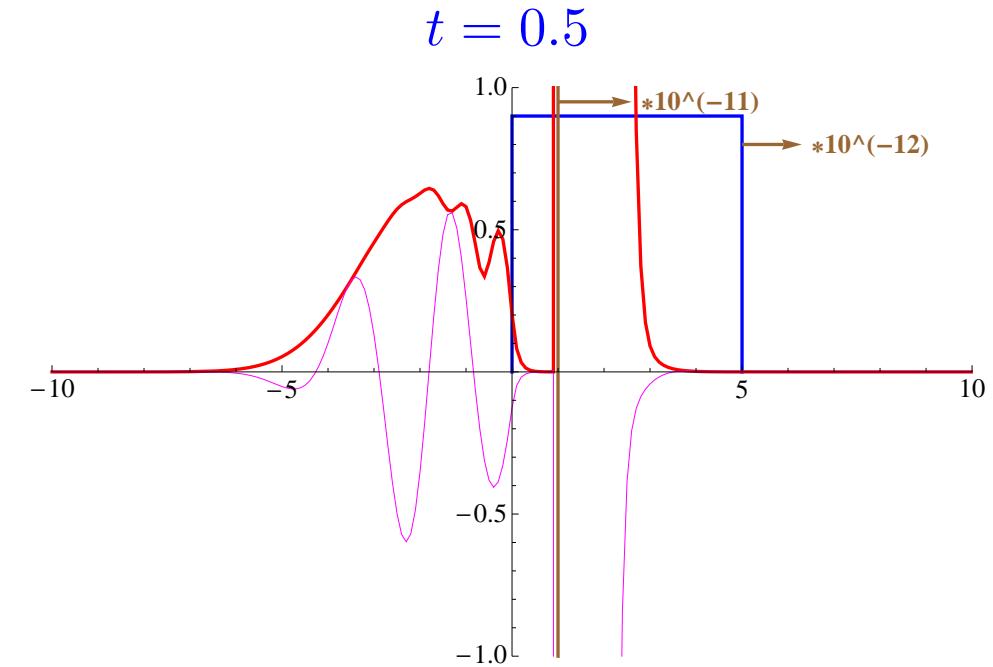
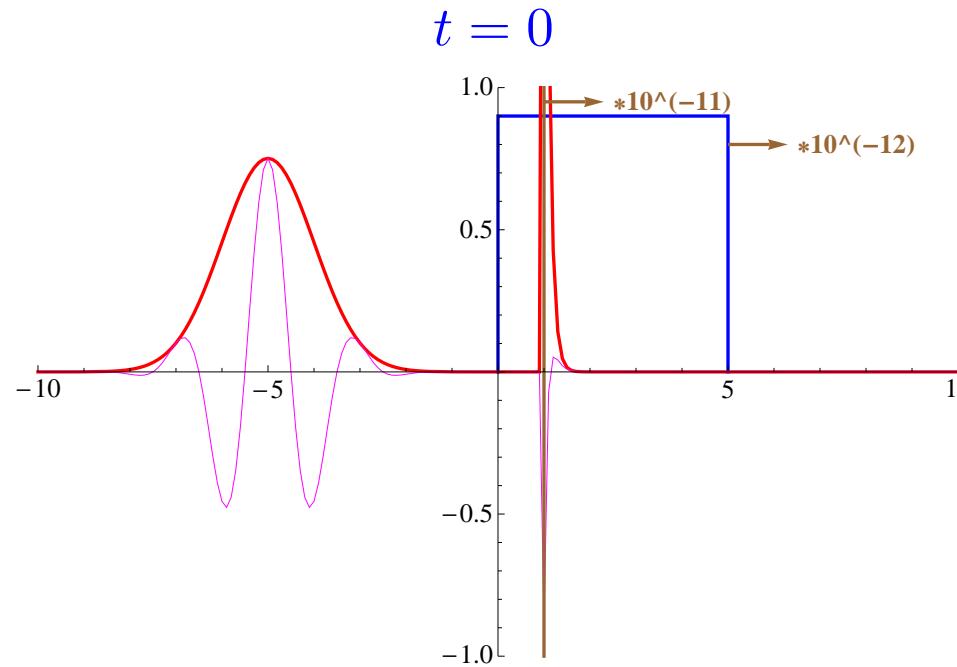
⇒ transmission dominated by

high momenta  $k \gtrsim \sqrt{V}$

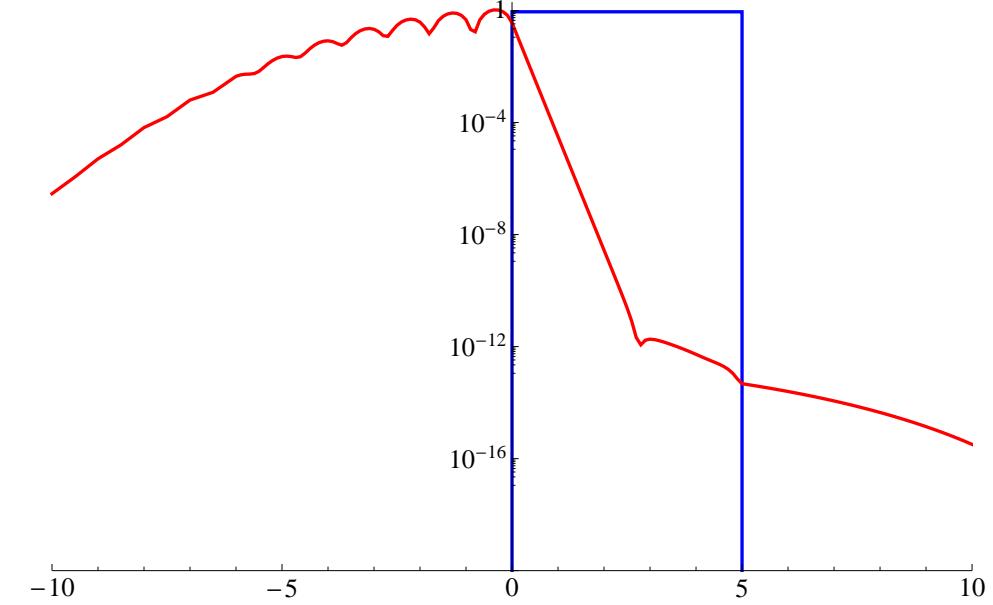
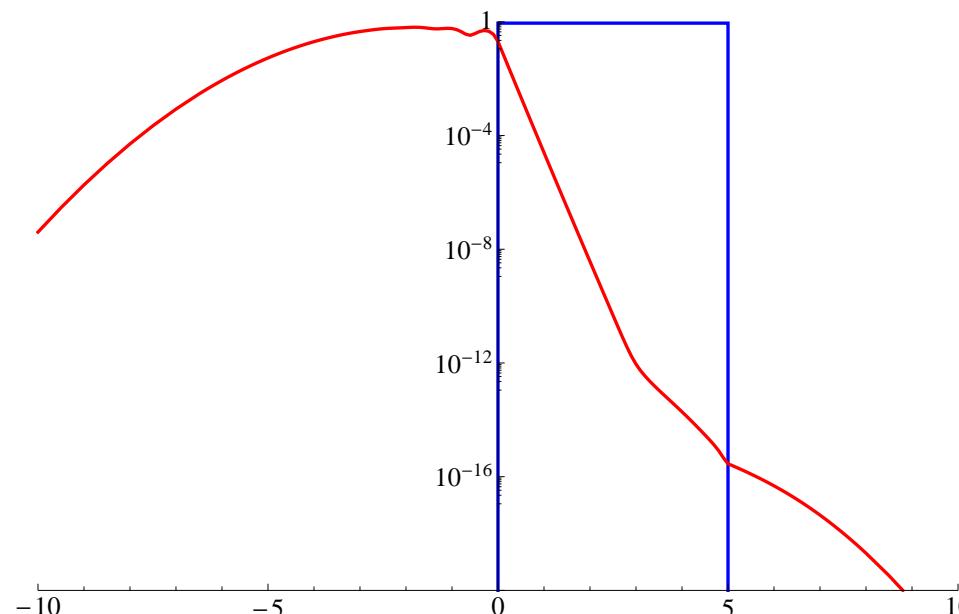
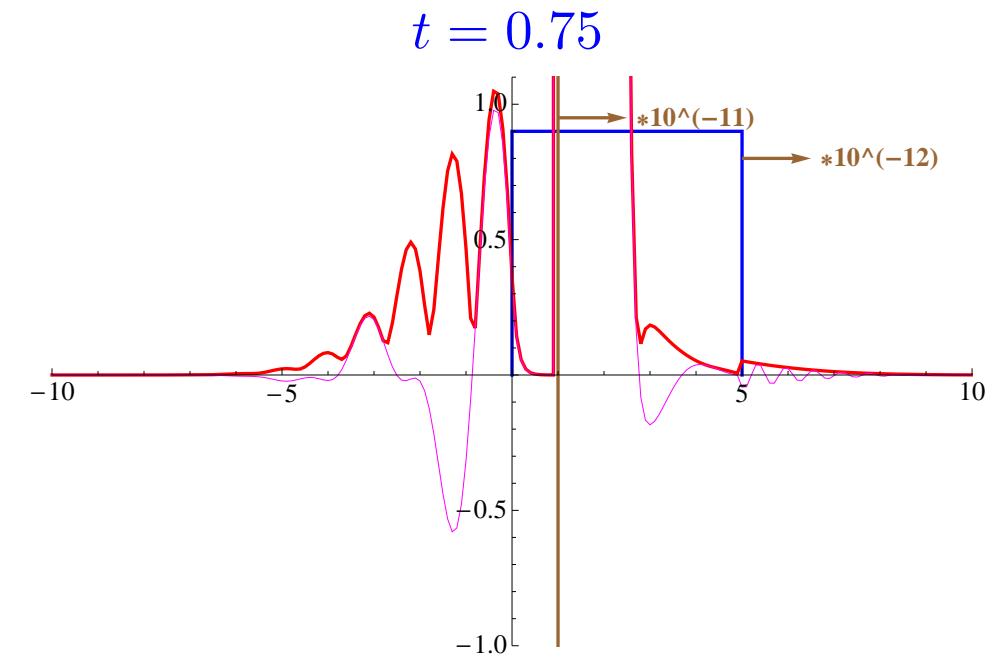
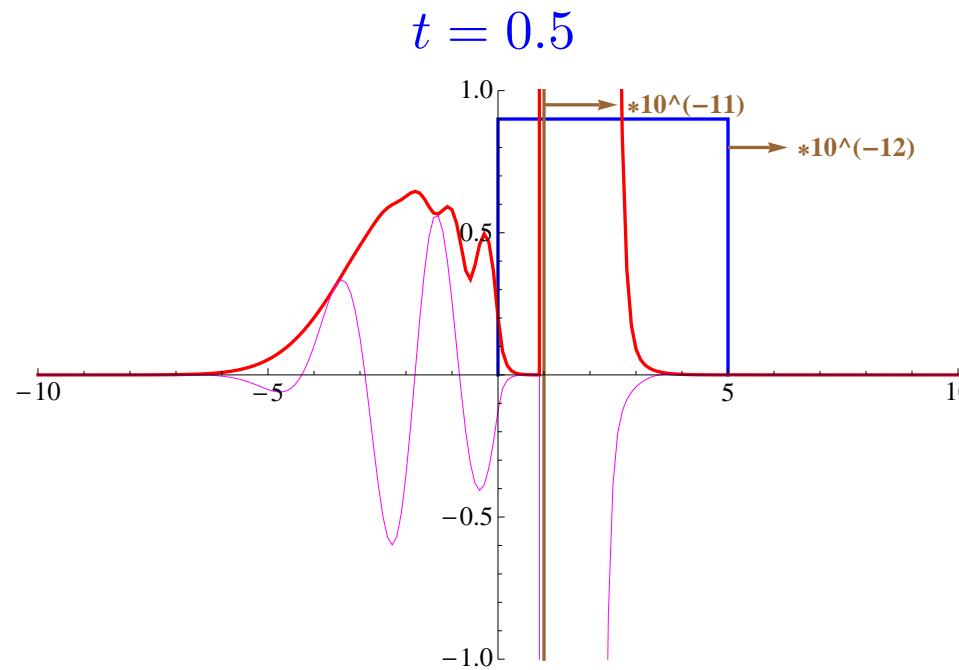
⇒ classically allowed transmission



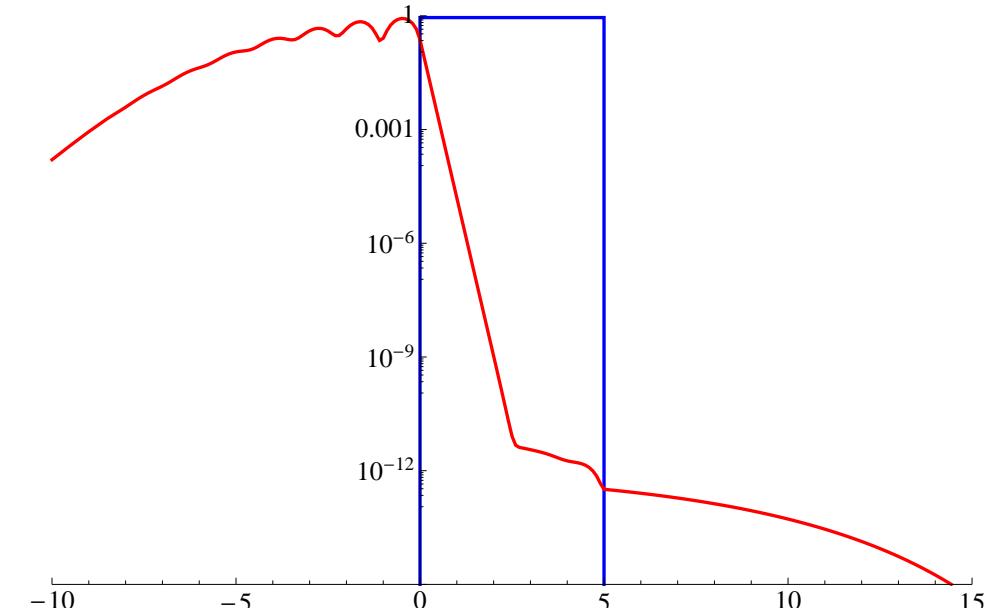
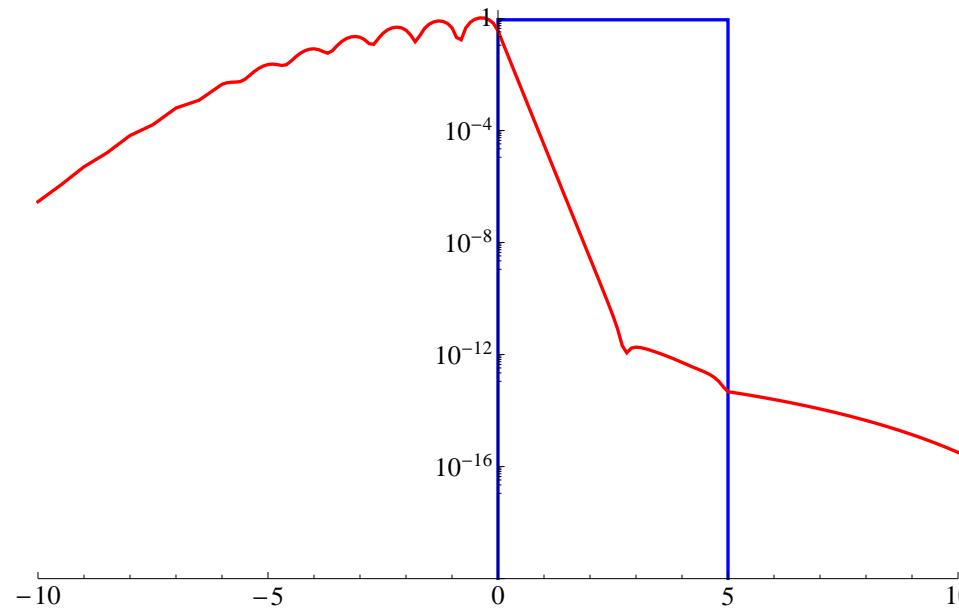
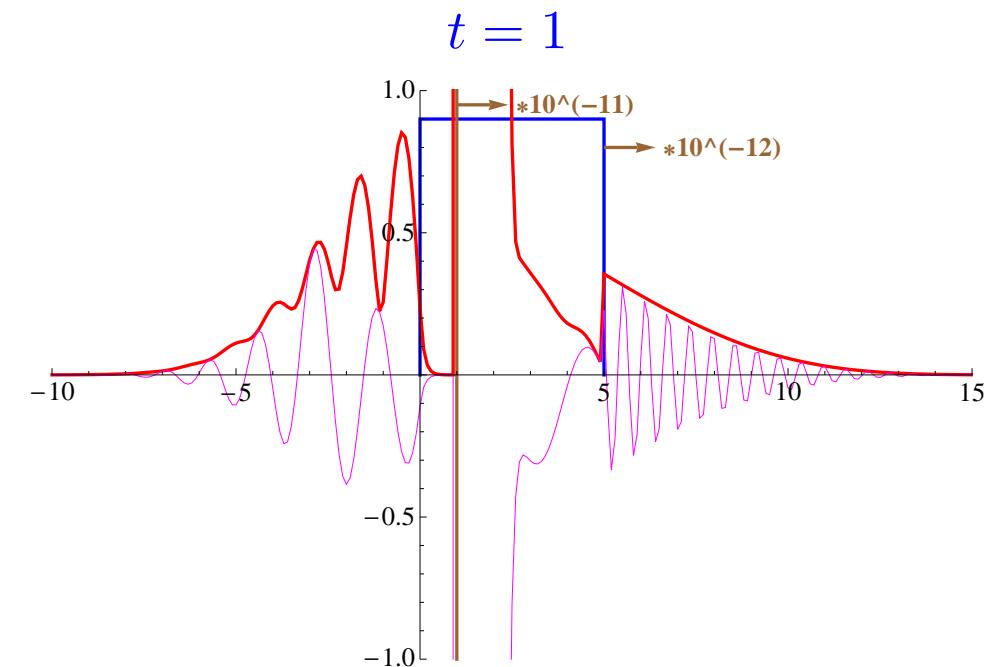
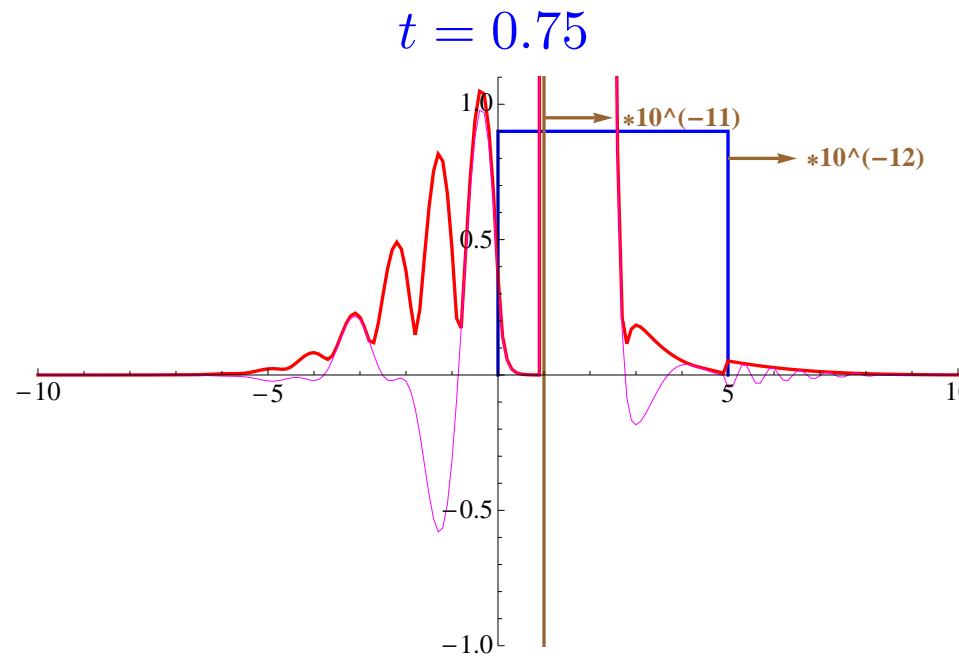
## B: “Slow” wave packet, thick barrier



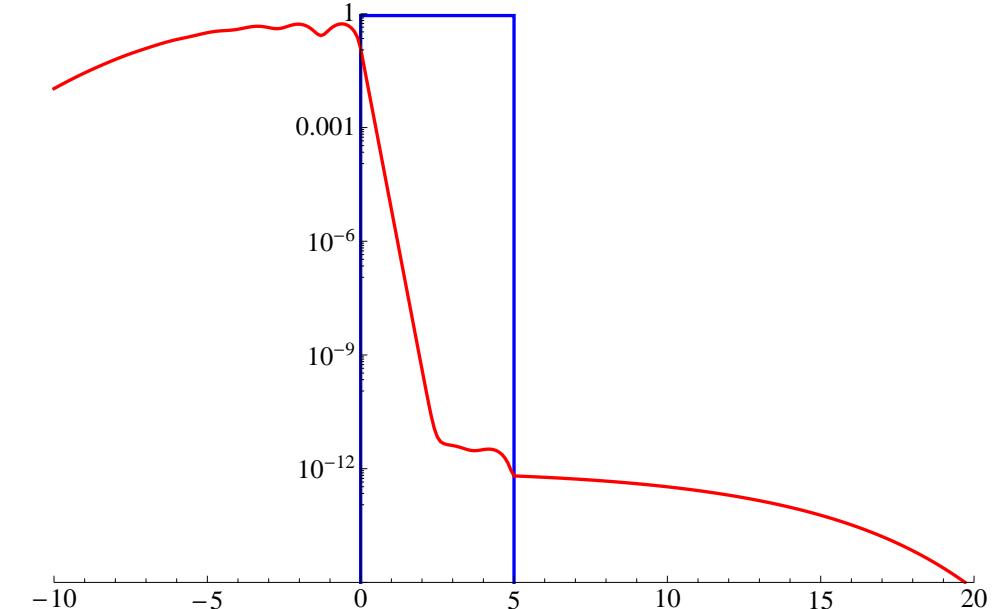
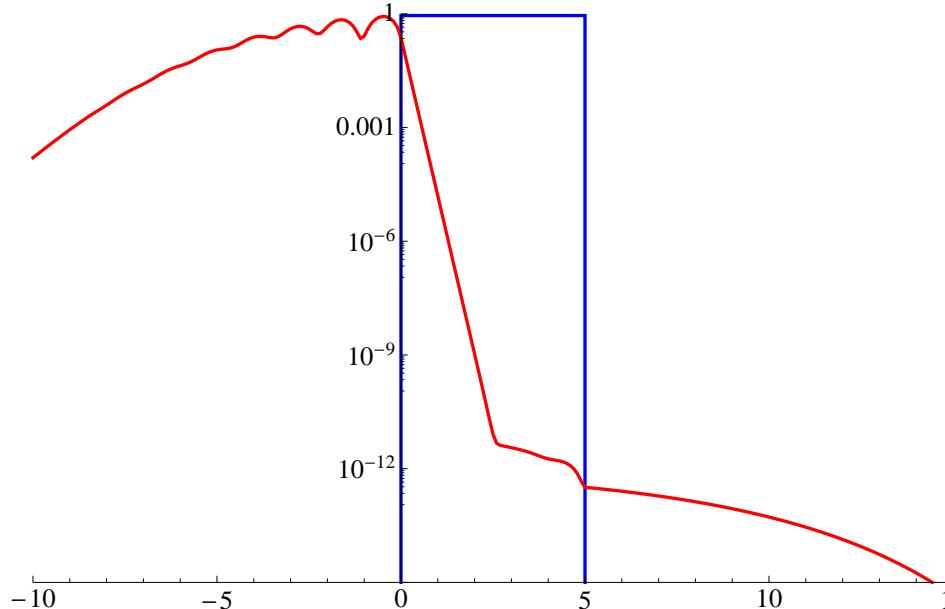
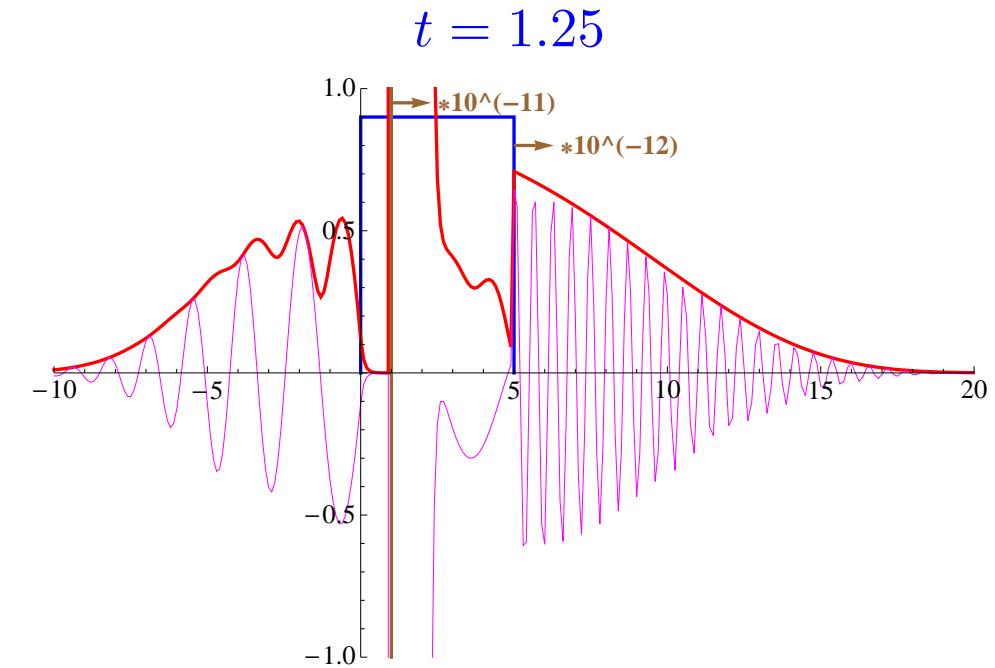
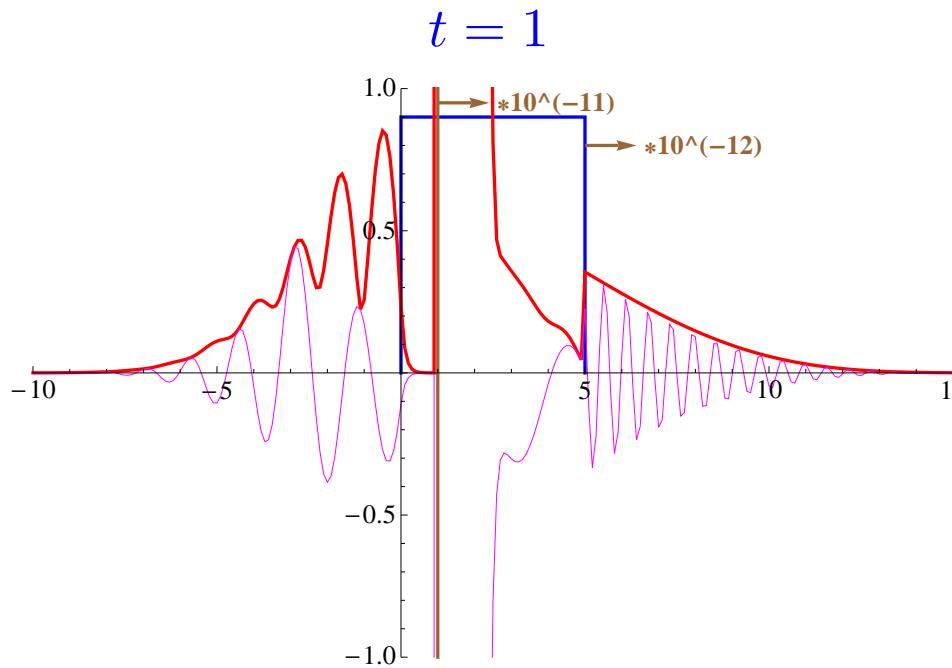
## B: “Slow” wave packet, thick barrier

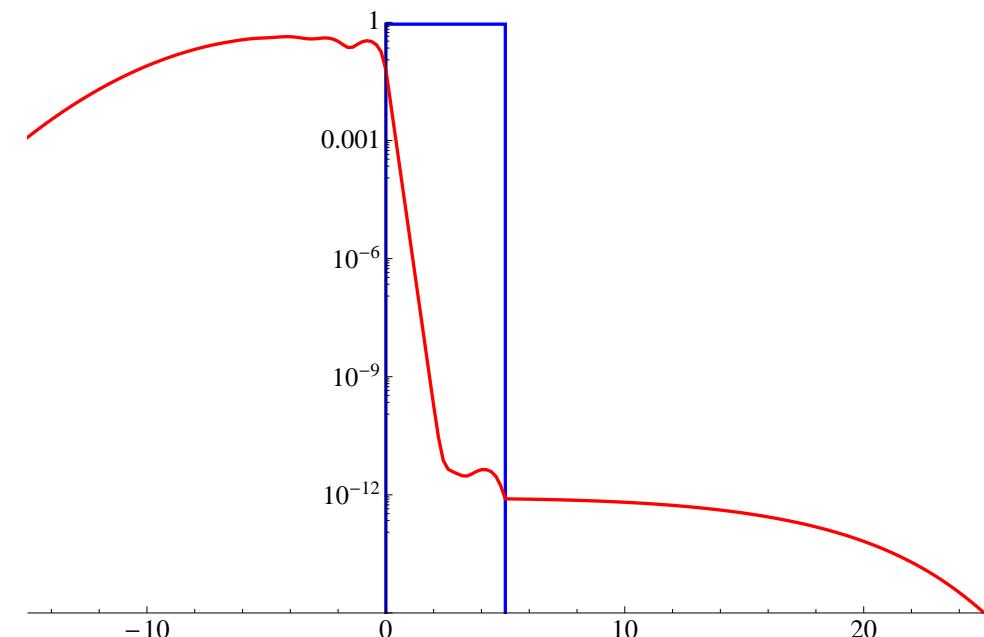
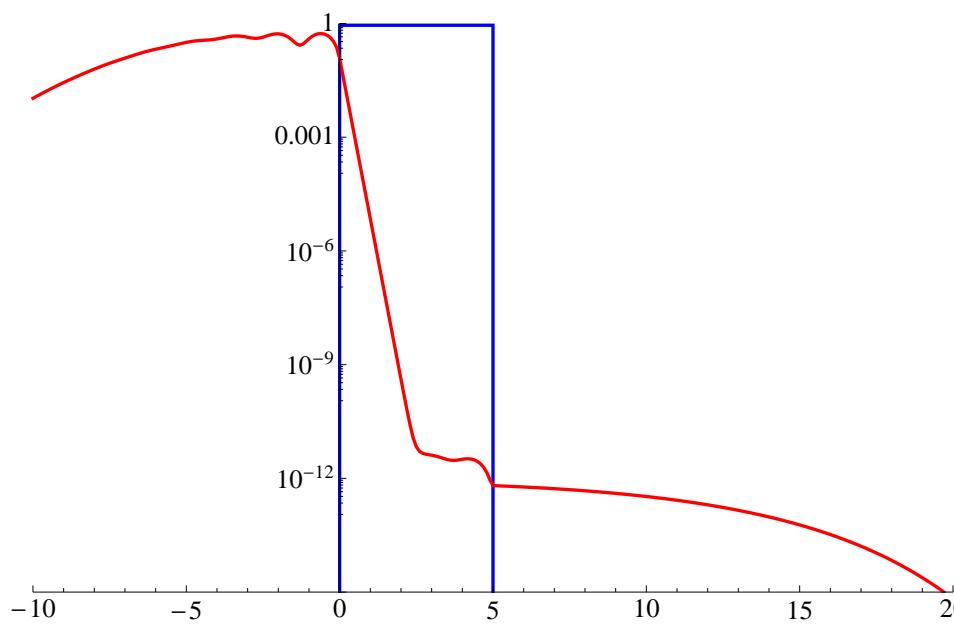
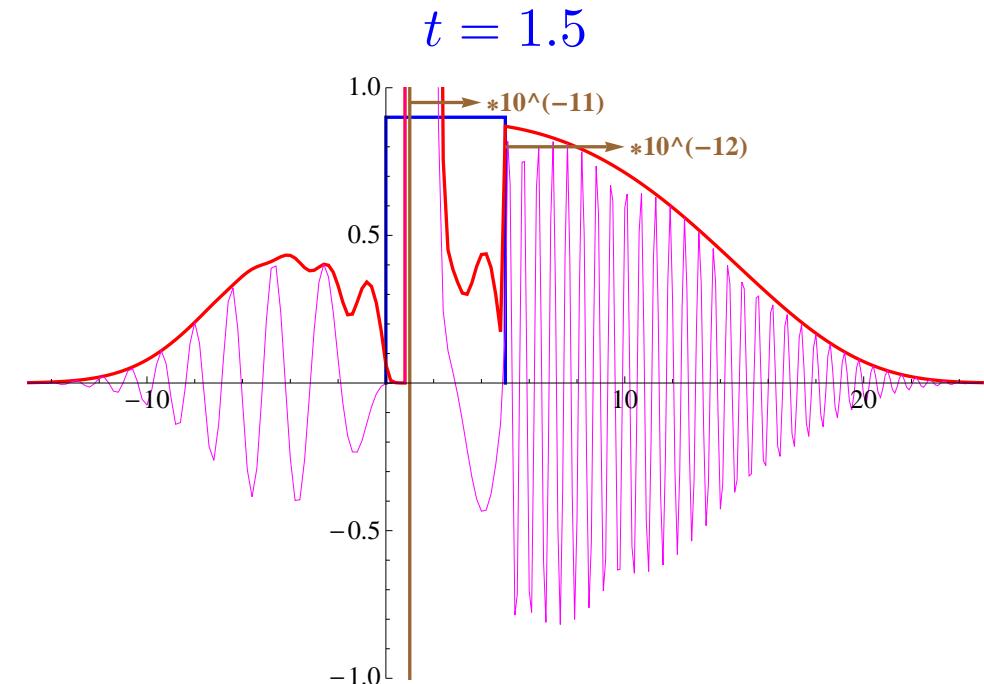
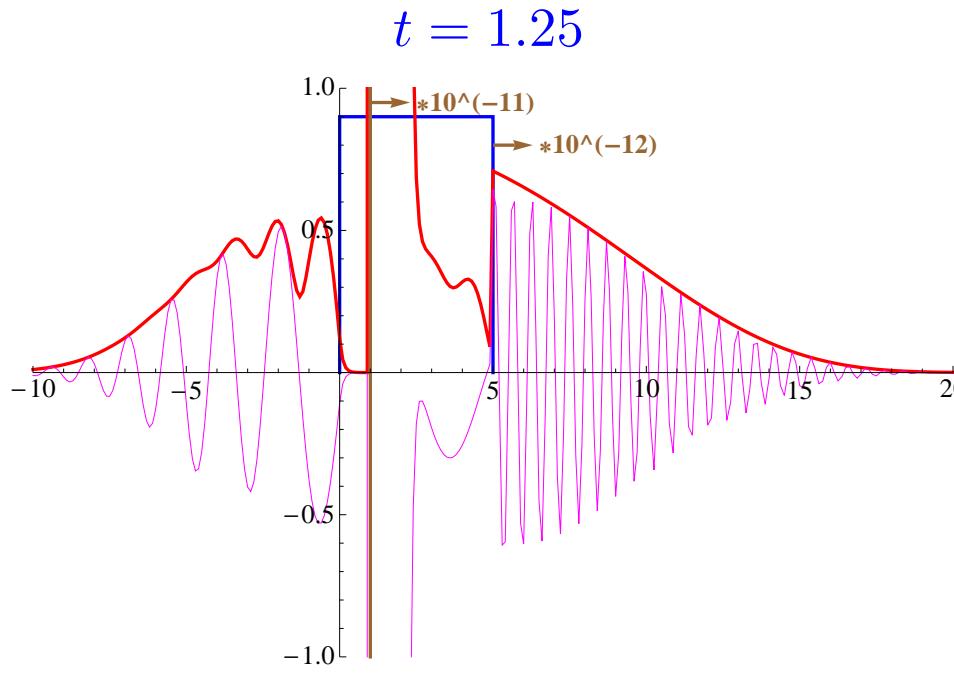


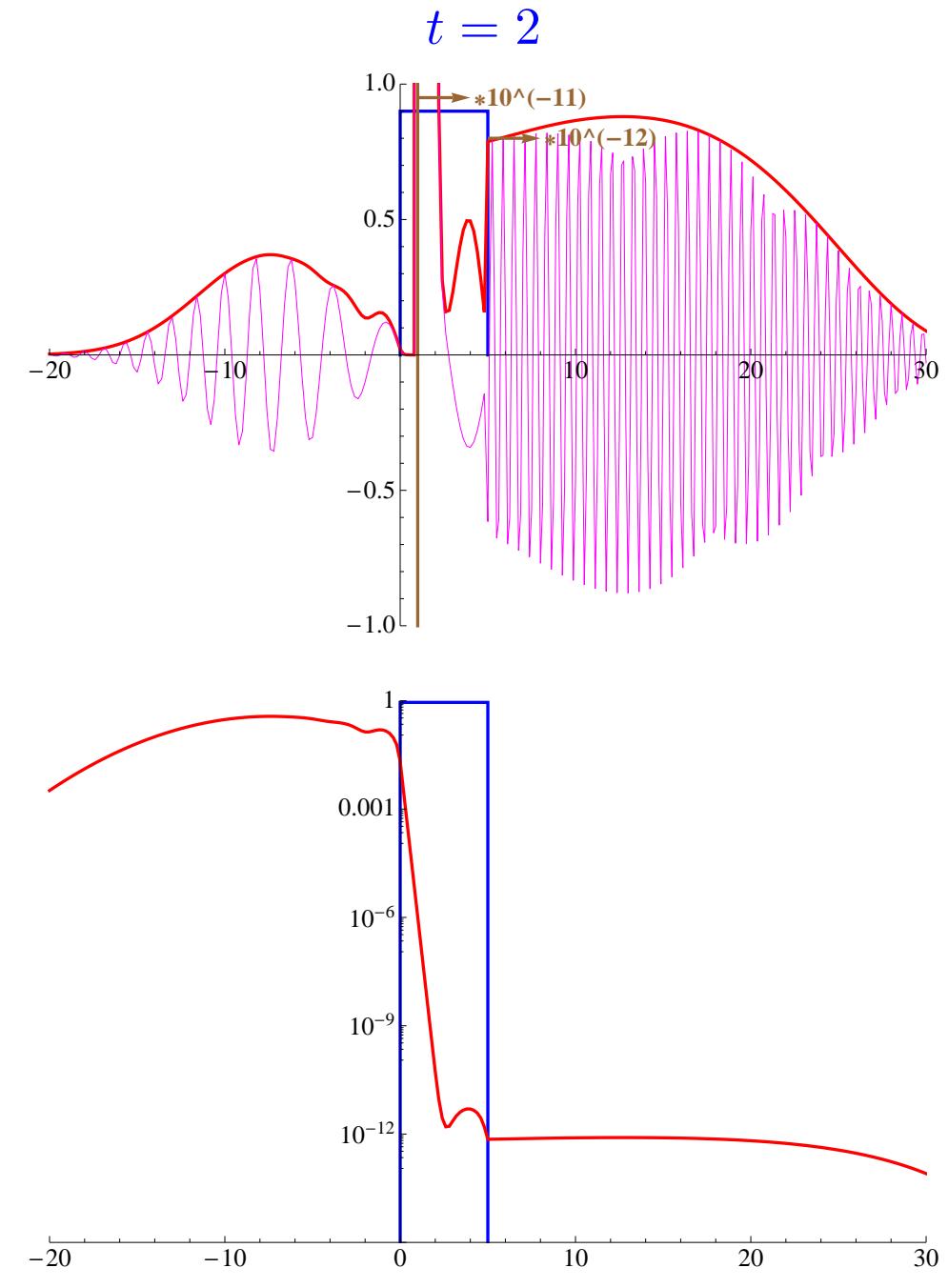
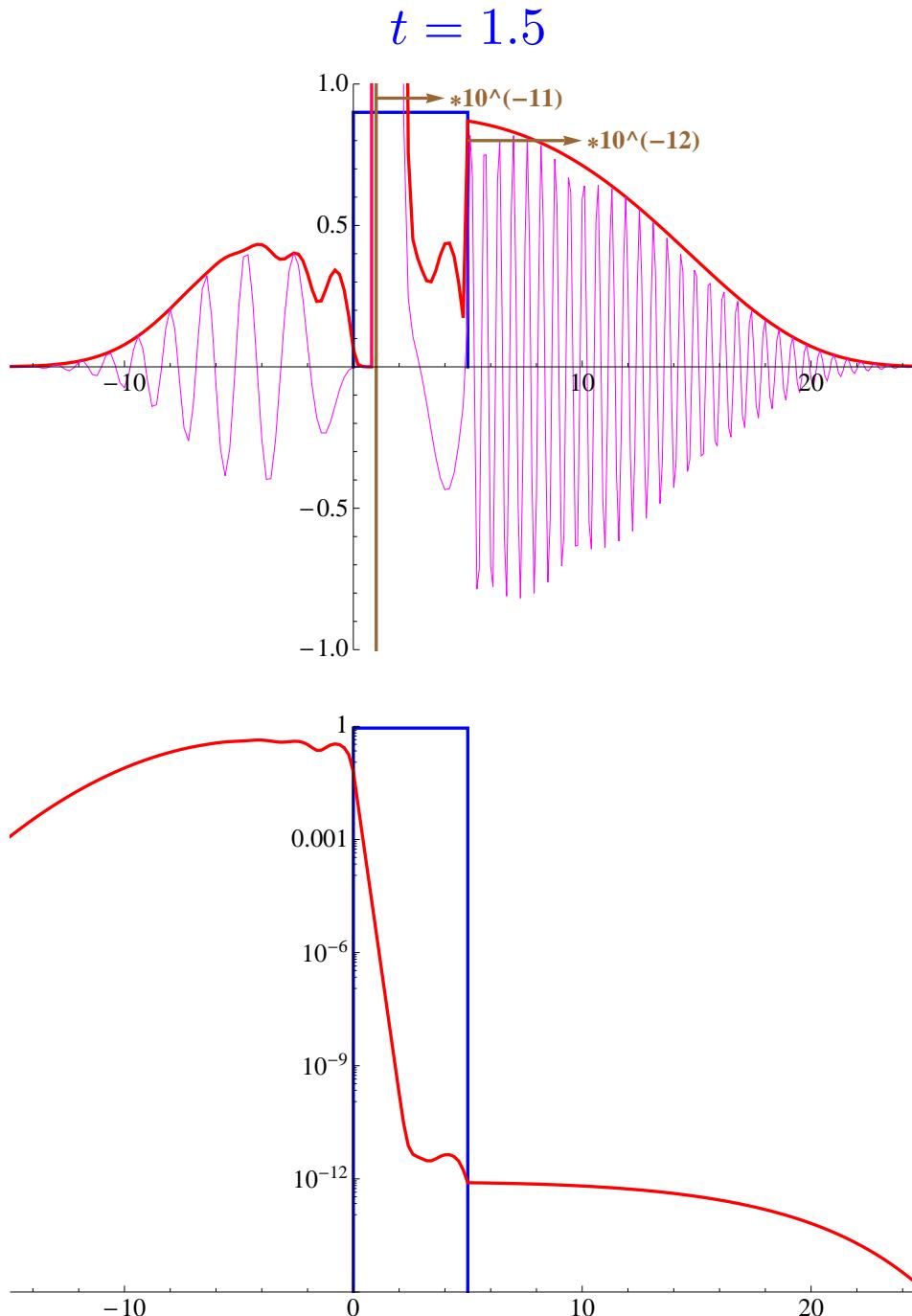
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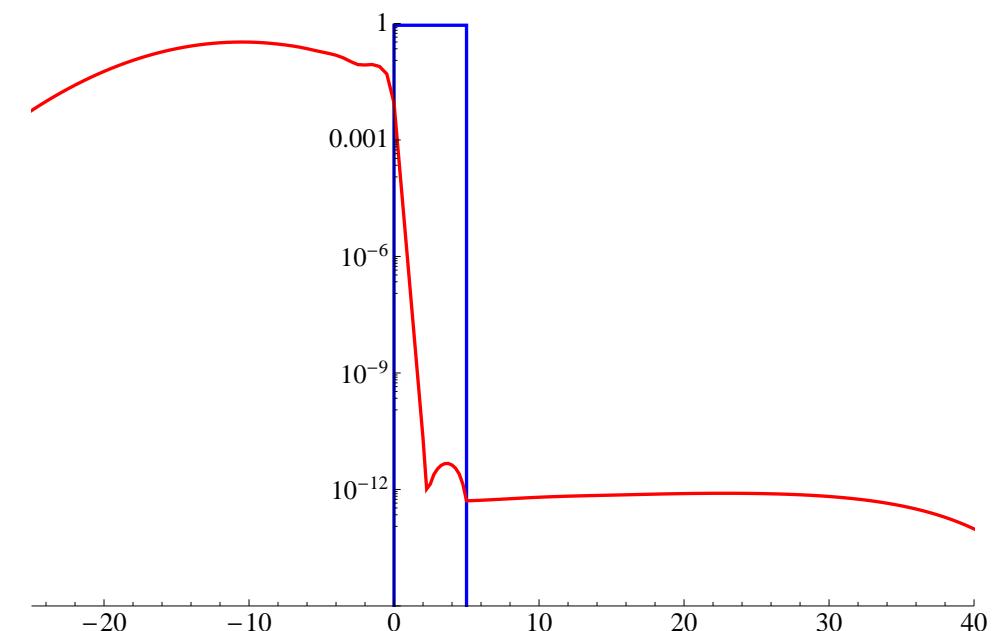
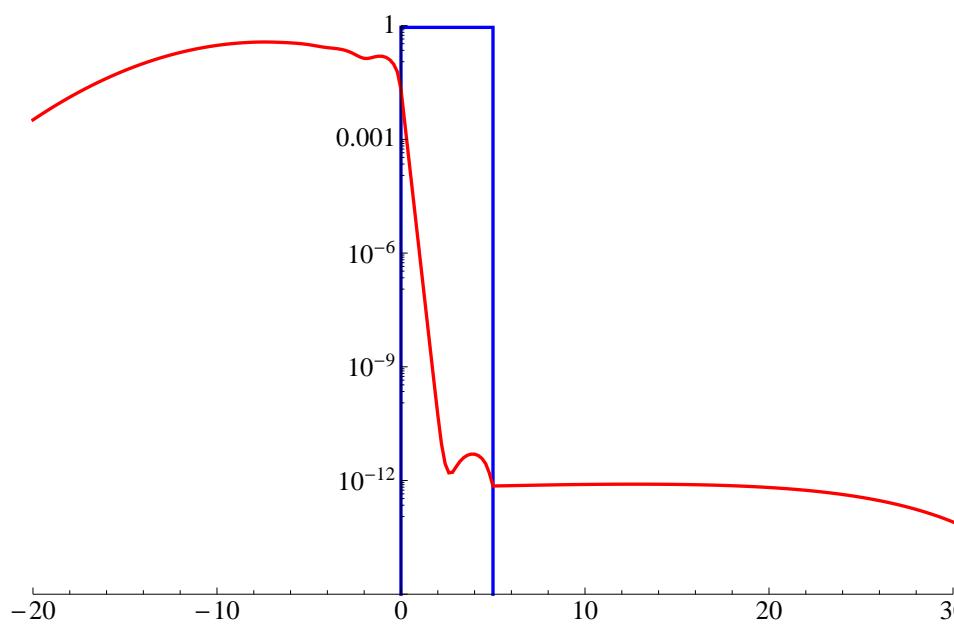
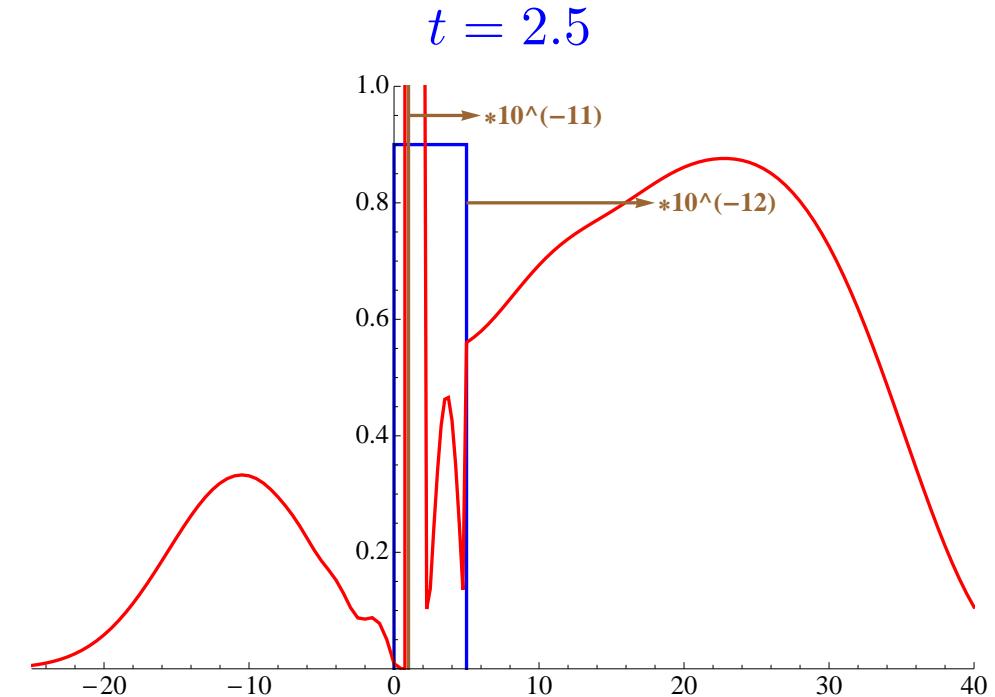
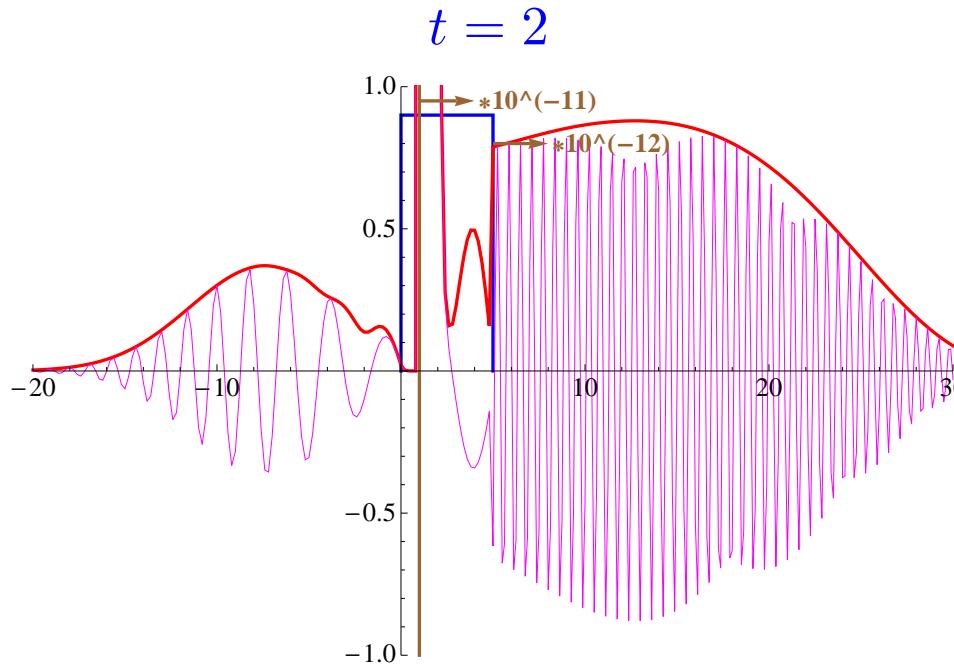


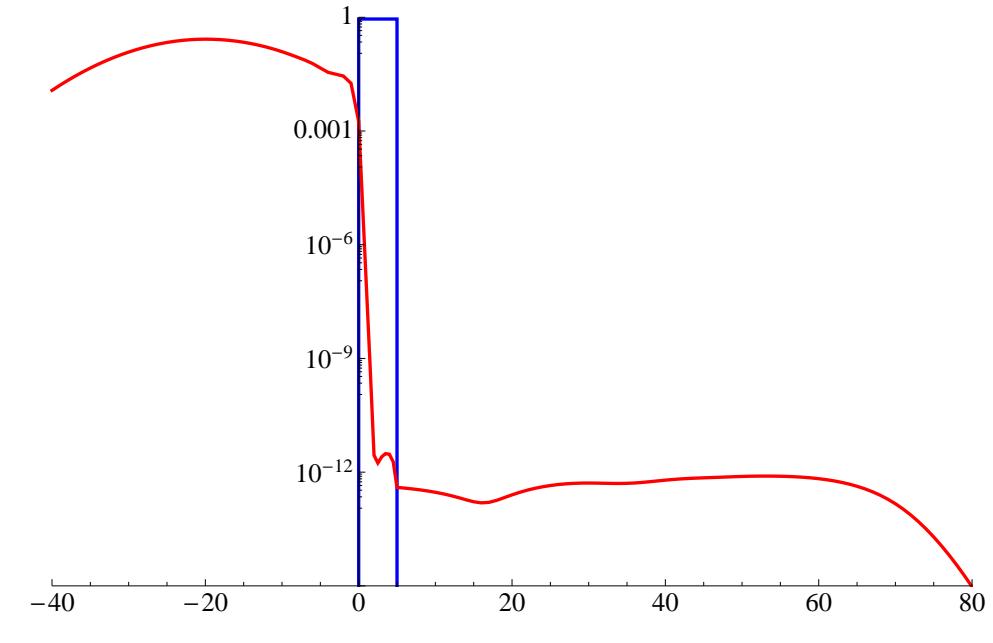
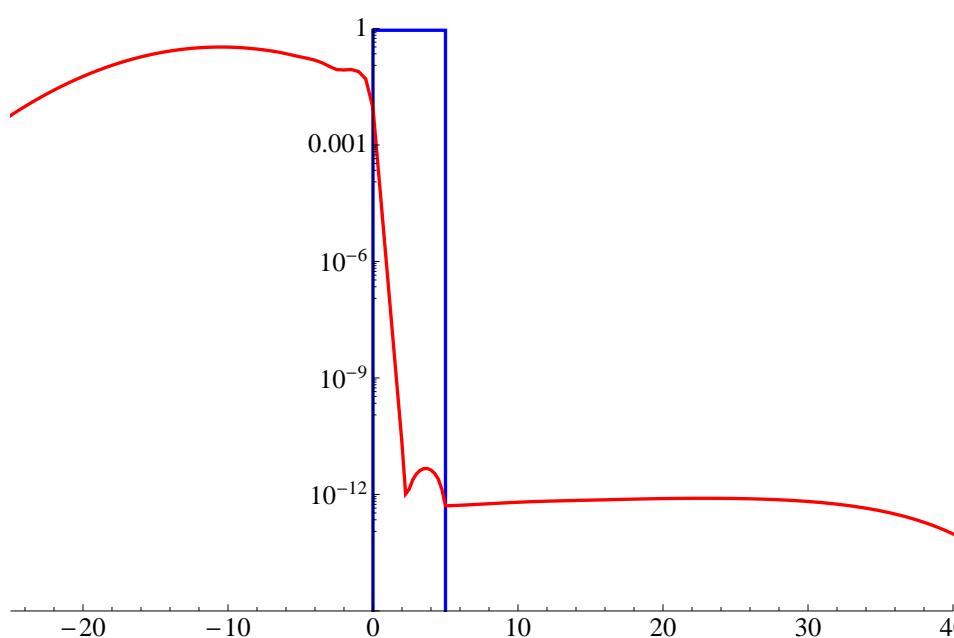
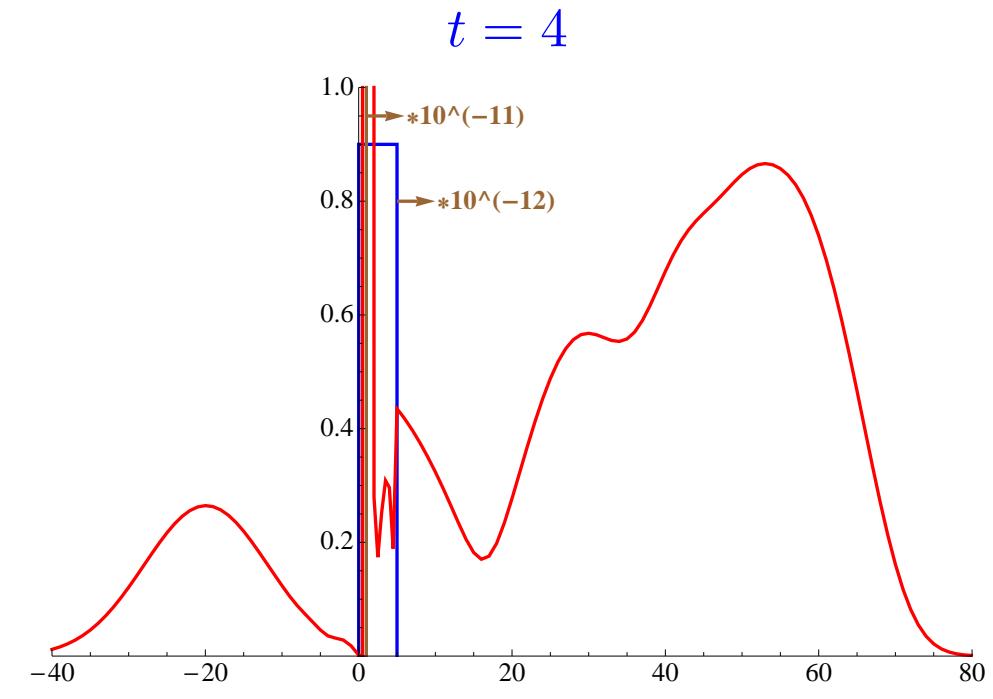
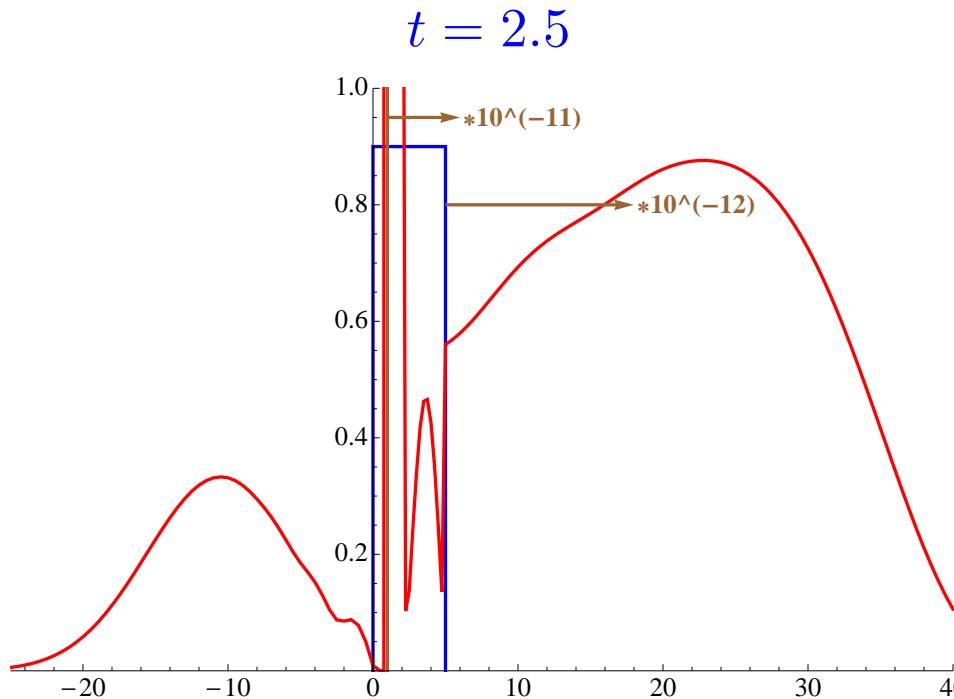
## B: “Slow” wave packet, thick barrier

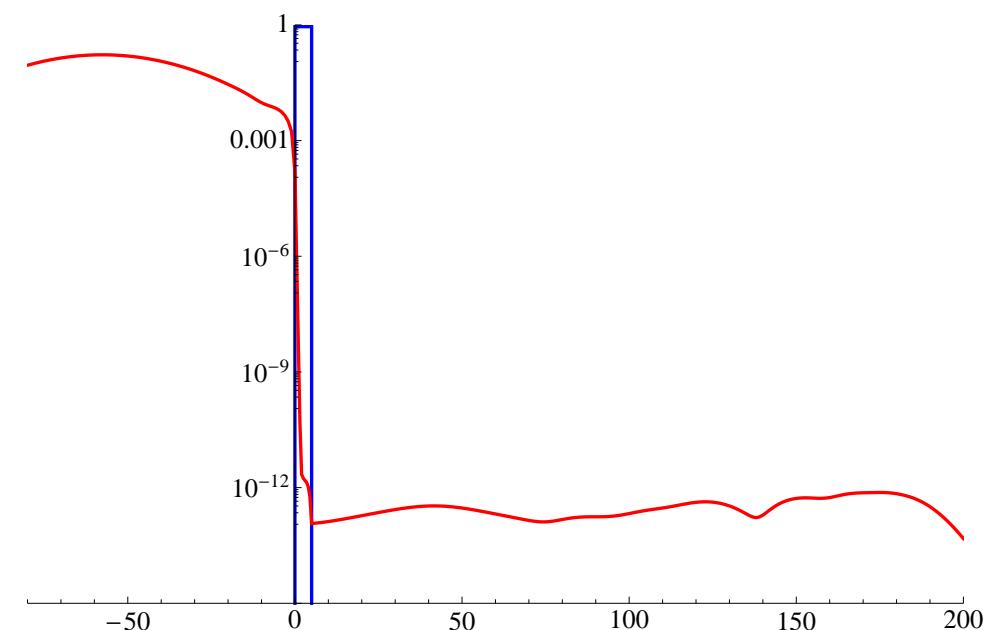
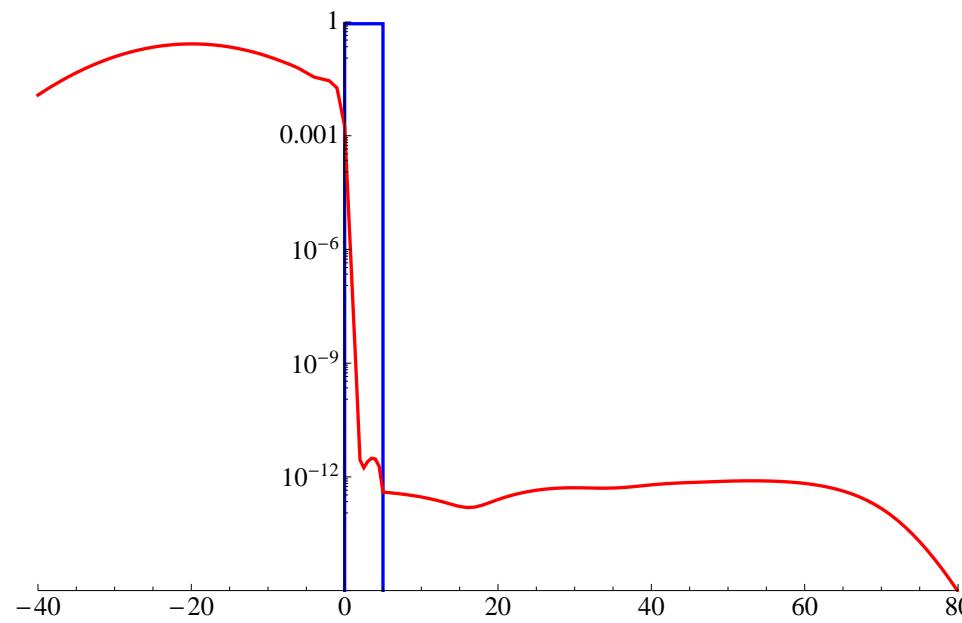
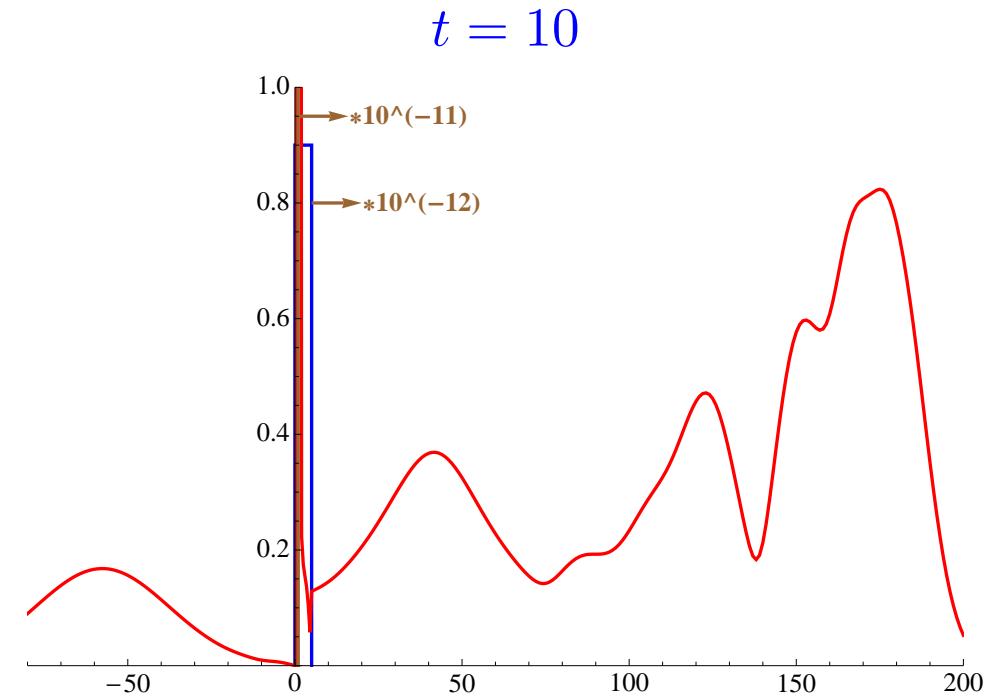
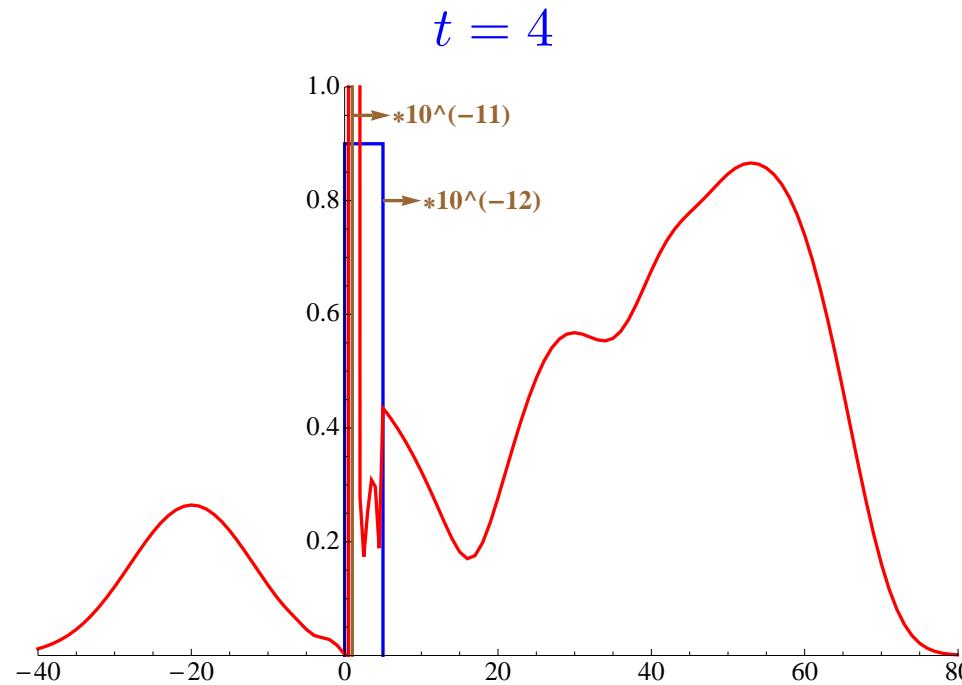


**B: “Slow” wave packet, thick barrier**

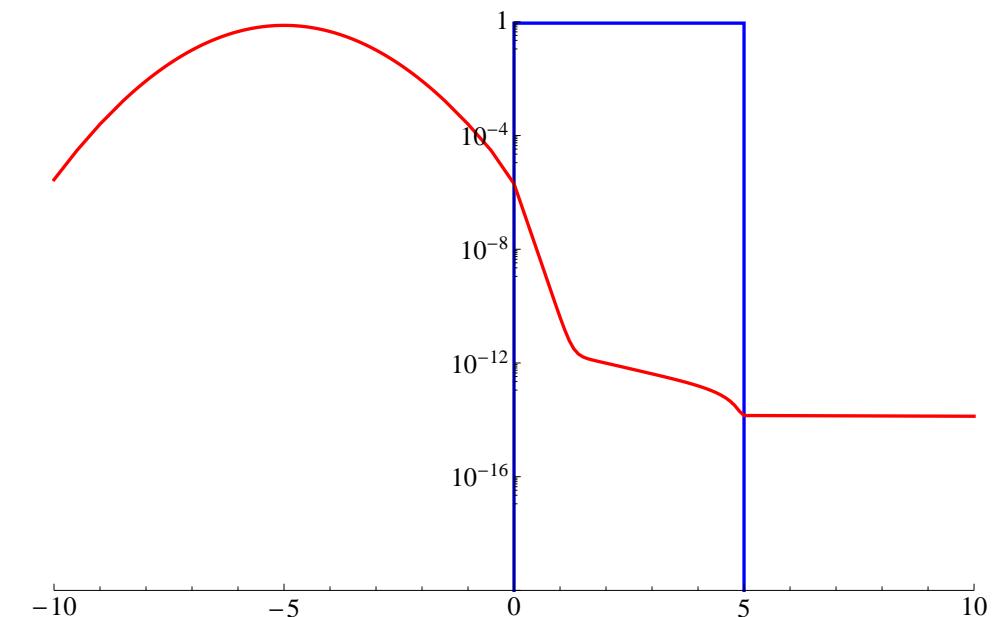
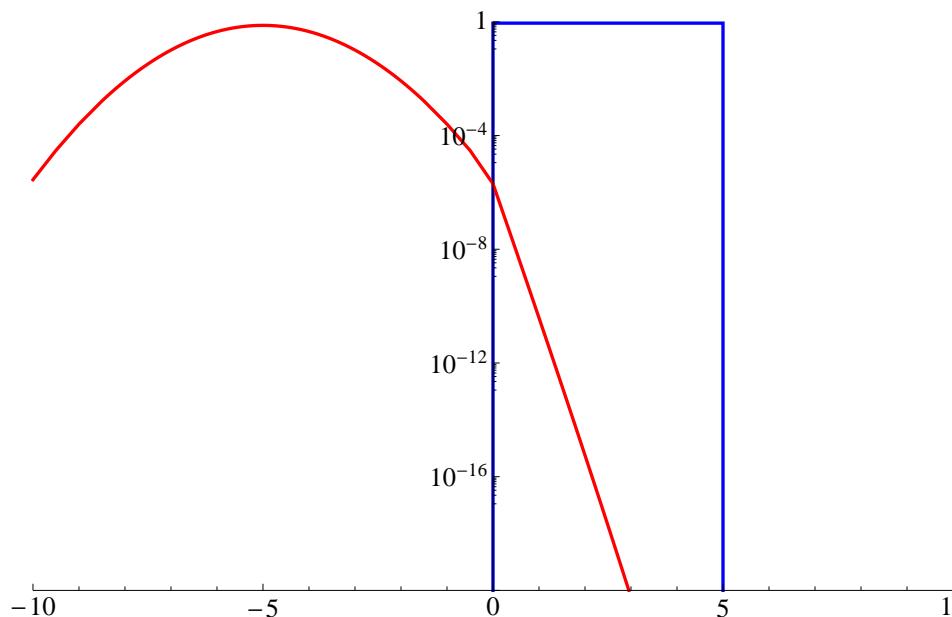
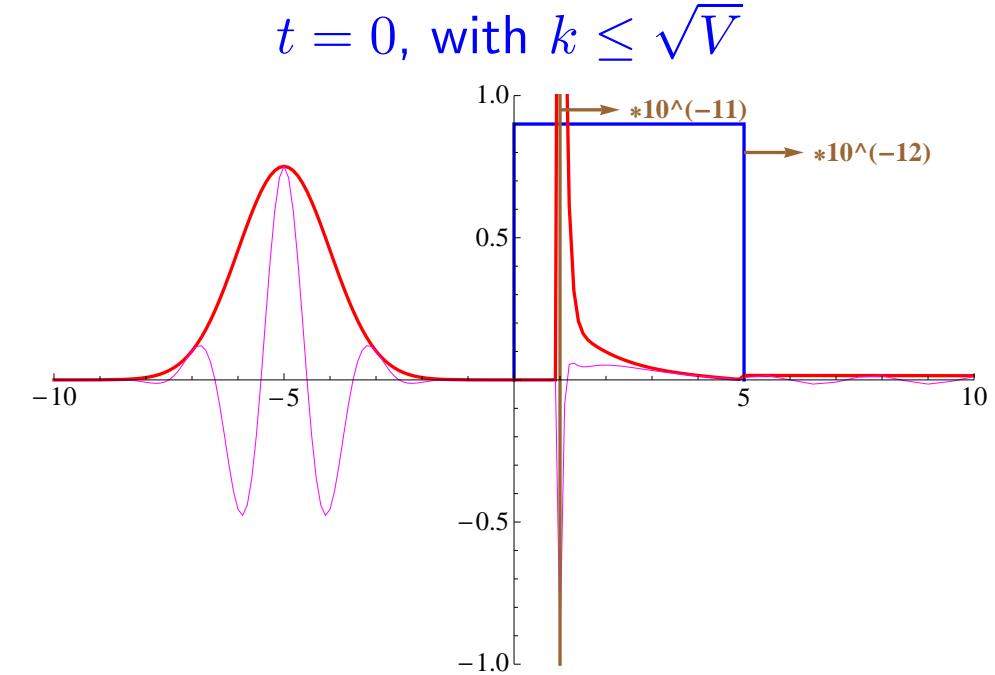
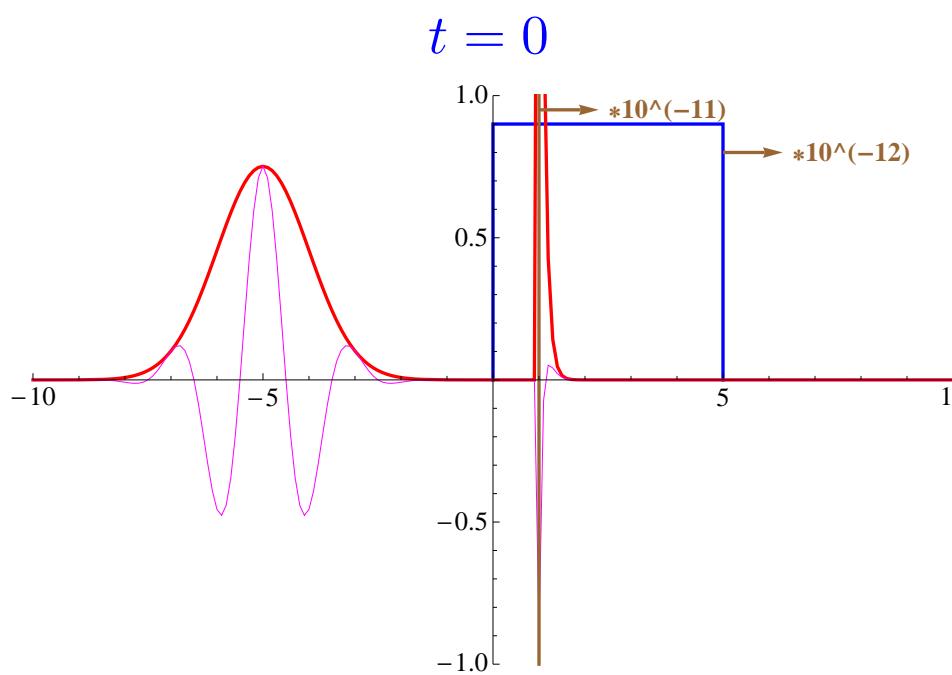
**B: "Slow" wave packet, thick barrier**

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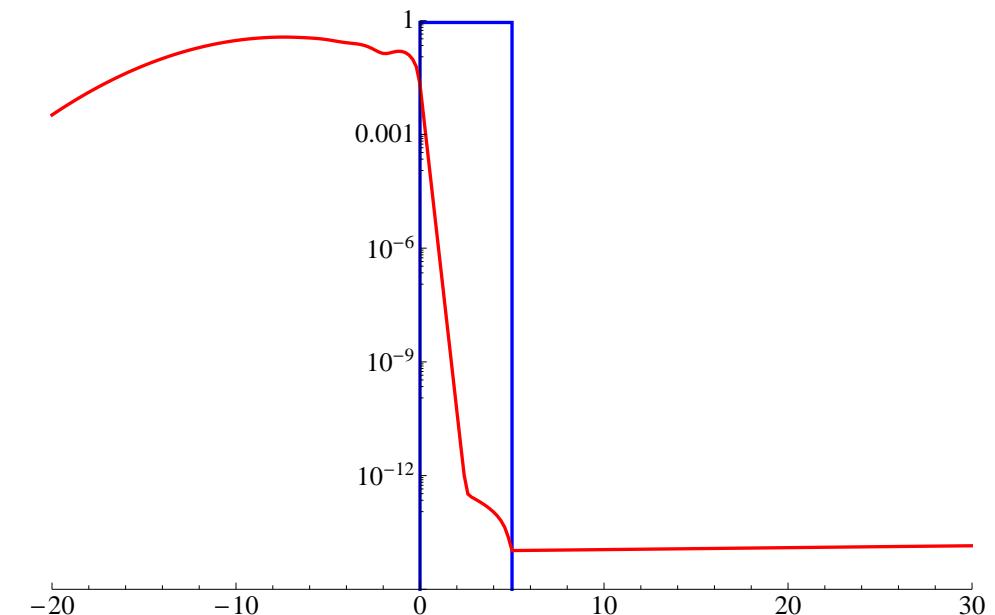
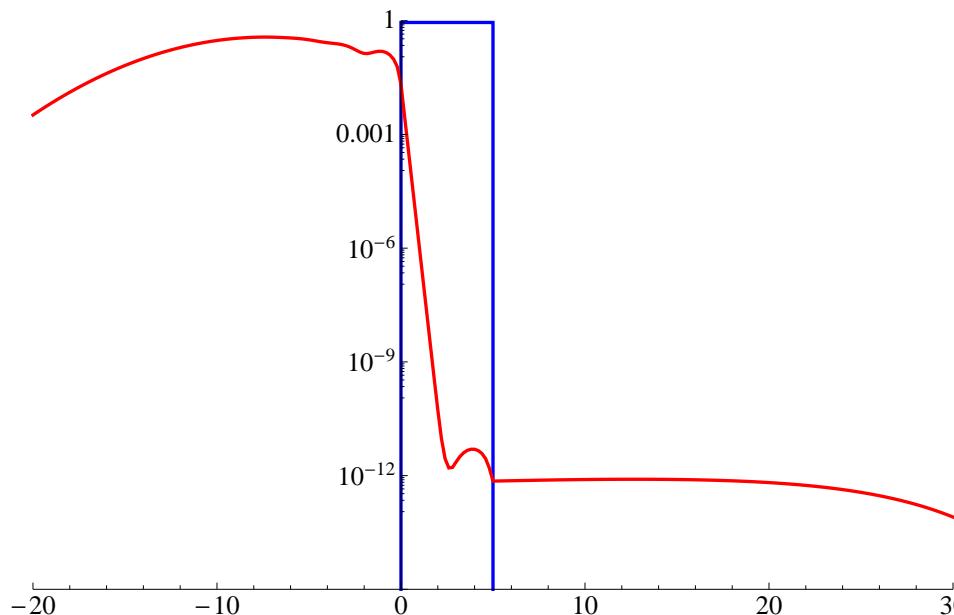
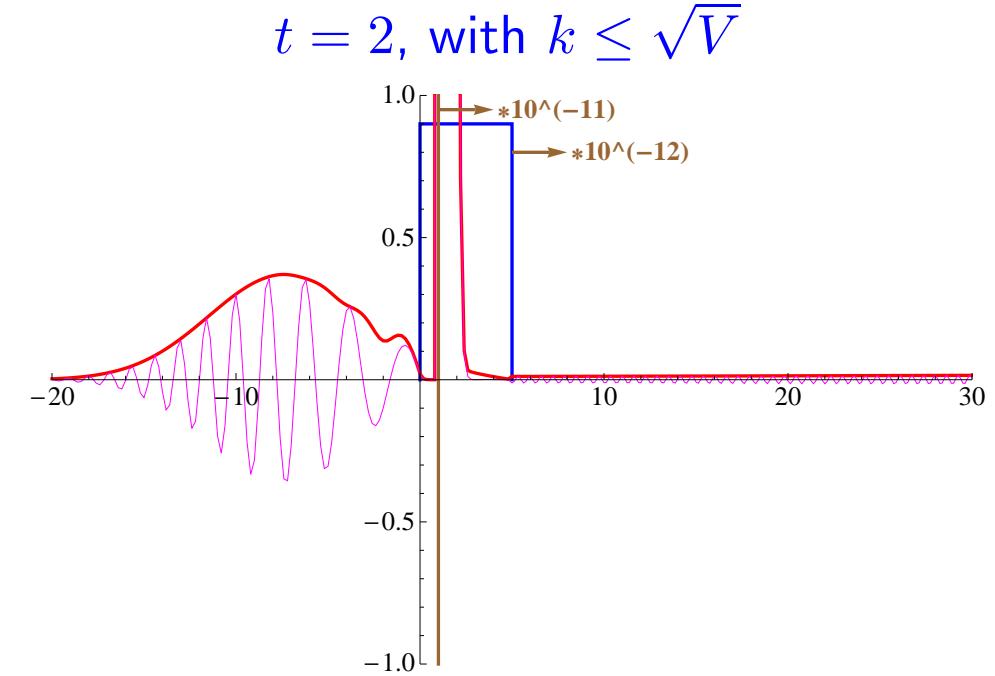
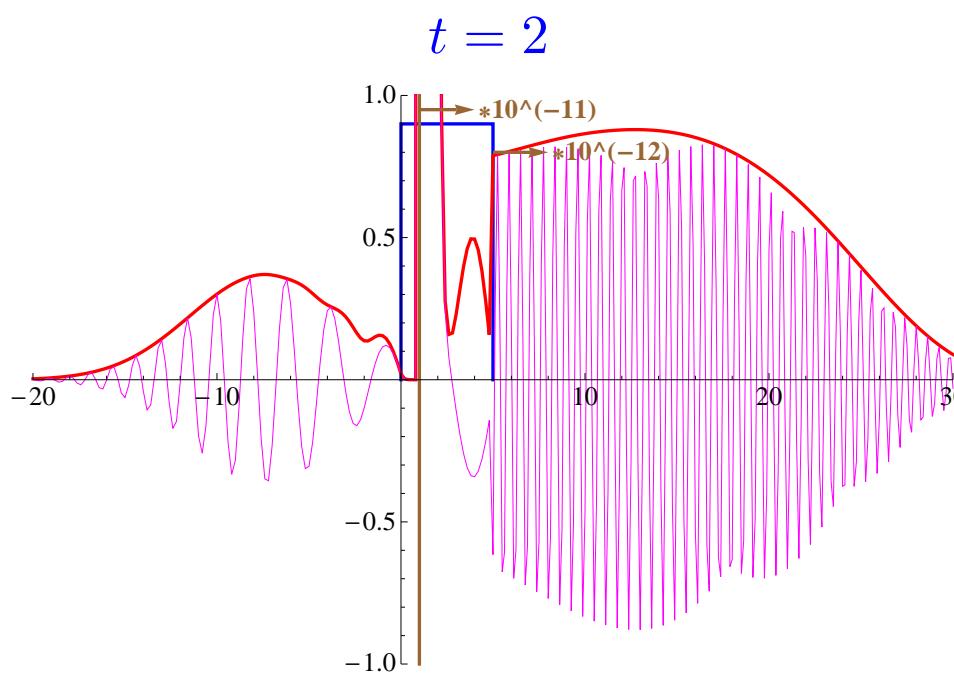
**B: “Slow” wave packet, thick barrier**

**B: “Slow” wave packet, thick barrier**

## B: “Slow” wave packet, thick barrier — omitting high momenta



## B: “Slow” wave packet, thick barrier — omitting high momenta (2)



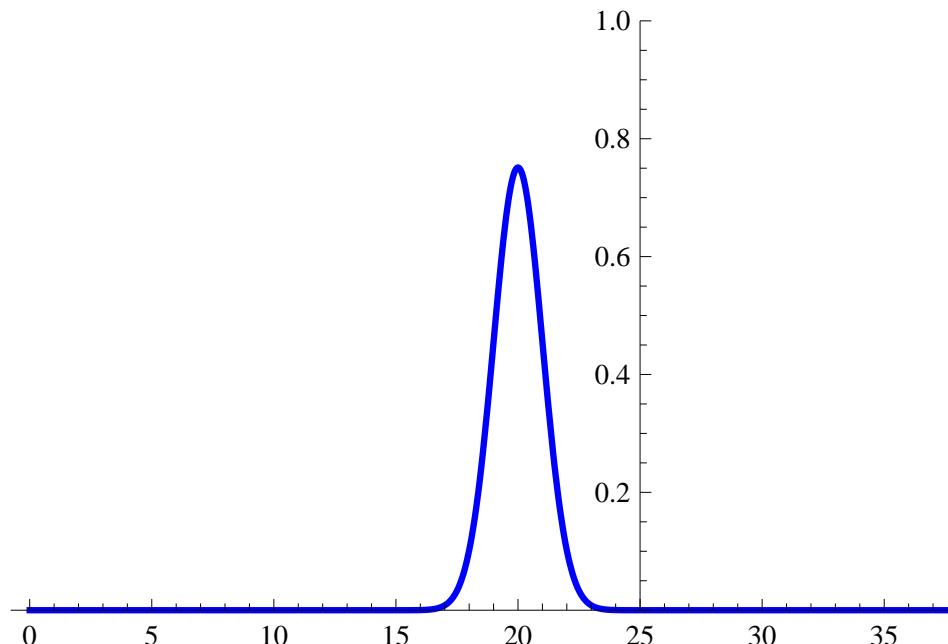
## B: “Slow” wave packet, thick barrier — Conclusion

- Thick barrier: strong exponential damping through the tunneling ( $\sim 10^{-20}$  at  $k_0$ )  
→ dominated by classically allowed transmission of high momenta ( $\sim 10^{-11}$ ).
- Incoming and reflected wave packet: like in case A – wave length and velocity of maximum correspond to  $k_0 = \pi$  and to a reflection at  $x = 0$ .
- In the tunnel: structure appears (ankle, local maximum)
- Transmitted wave packet: wave length and velocity of maximum  
 $\rightsquigarrow k \approx 10 = \sqrt{V} \gg k_0$
- Omitting high momenta  $k > \sqrt{V}$  in integration:  
no transmission from low momenta
- Packet maximum enters tunnel at  $t = 0.8$  and exits at  $t = 1.6$   
 $\Rightarrow \Delta t_{\text{tunnel}} \approx 0.8 > 0$  ( $v_{\text{tunnel}} = L/\Delta t_{\text{tunnel}} \approx 2\pi \rightsquigarrow k_{\text{tunnel}} \approx k_0$ )
- Transmitted wave packet develops tail of lower maxima.

## C: “Fast” wave packet, thick barrier

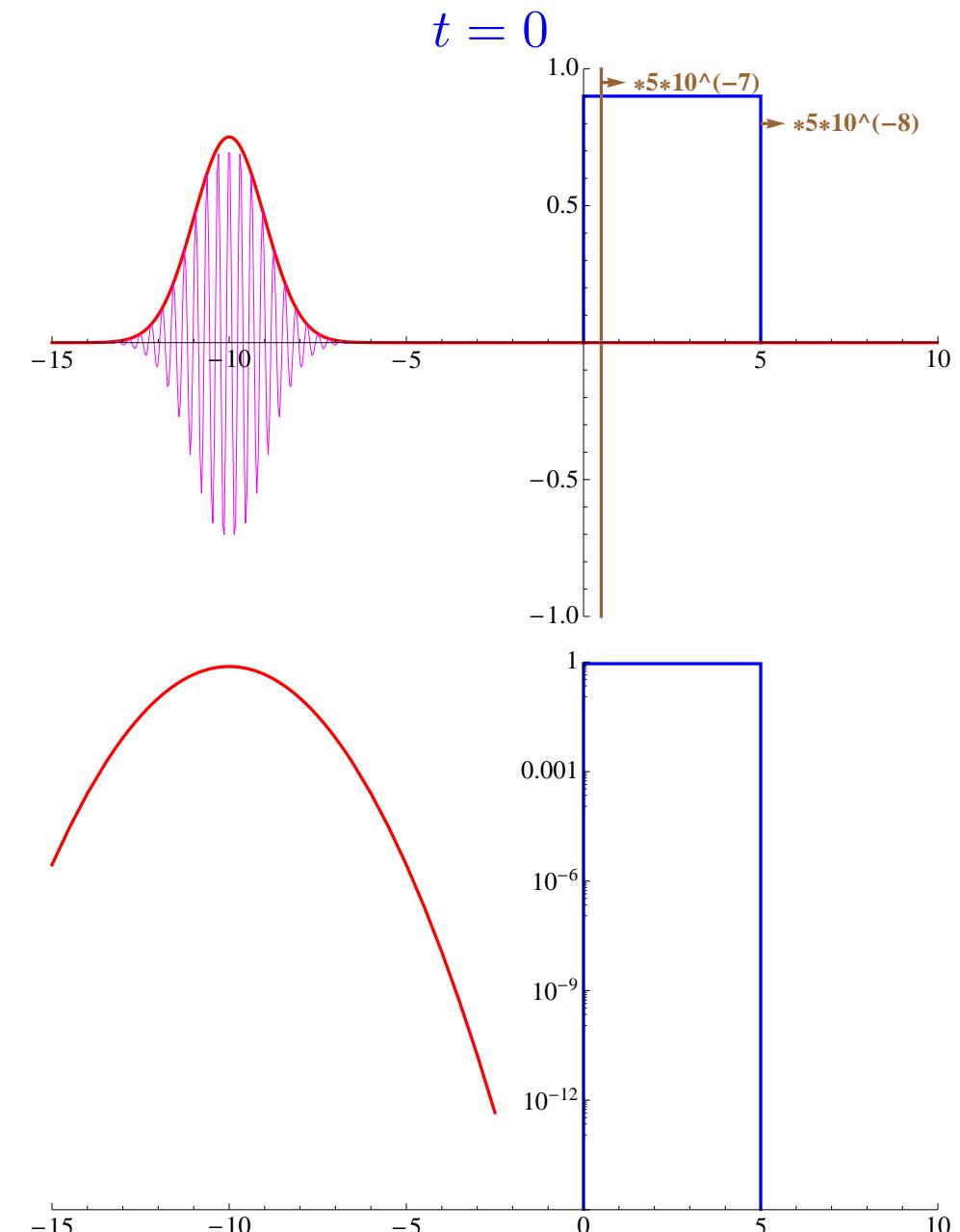
**Parameters:**  $k_0 = 20 \gg \sigma_k$ ,  
 $L = 5 > \sigma_x, \gg \pi/k_0$ ,  $\sigma_k = 1/\sqrt{2} \approx 0.7$ ,  
 $x_0 = -10$ ,  $2m = 1$ ,  $V = 25^2$

**Momentum distribution**  $|A(k)|$ :



**Right plots:** barrier,  $|\psi(t, x)|$ ,  $\text{Re } \psi(t, x)$ ;

top: to scale by  $5 \cdot 10^{-7}$  for  $0.5 < x < L$   
and by  $5 \cdot 10^{-8}$  for  $x > L$ ;

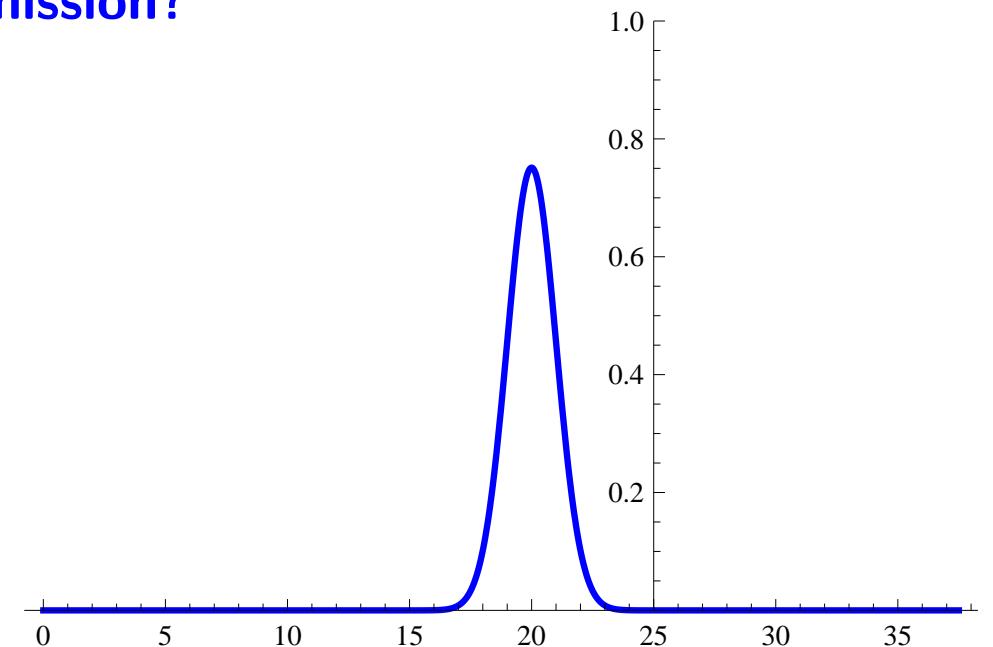


### C: “Fast” wave packet, thick barrier — Which momenta contribute to the transmission?

Momentum distribution of incoming wave

packet, i.e. part of  $\tilde{\psi}(k; t, x) \propto e^{ikx}$  at

$x < 0$ :  $|A(k)|$



Momentum distribution of transmitted

wave packet, i.e.  $\tilde{\psi}(k; t, L)$  at  $x > L$ :

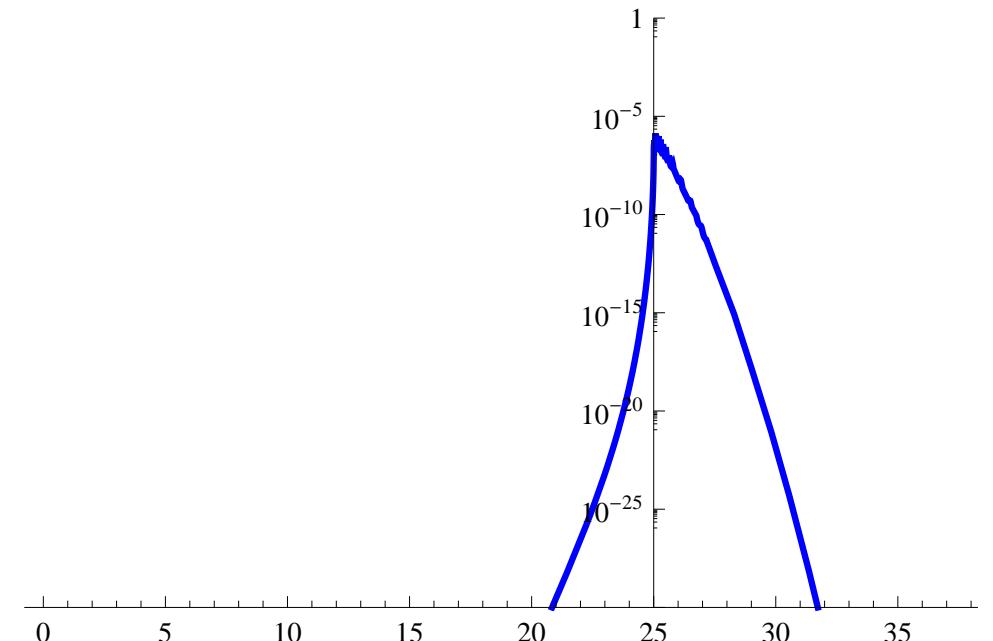
$|A(k) \tilde{\psi}(k; t, L)|$  (logarithmic scale)

⇒ like in case B:

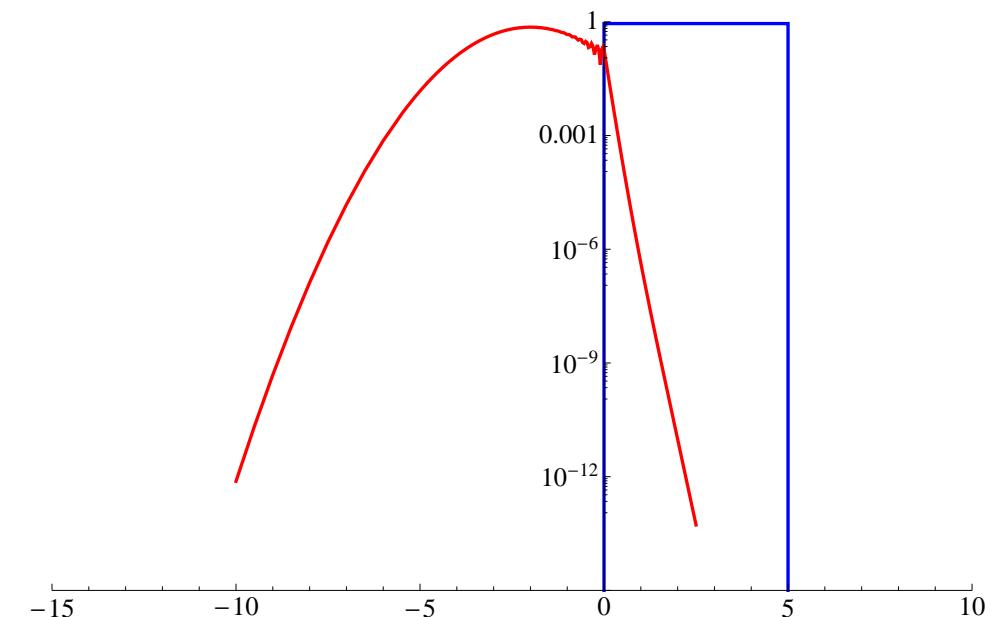
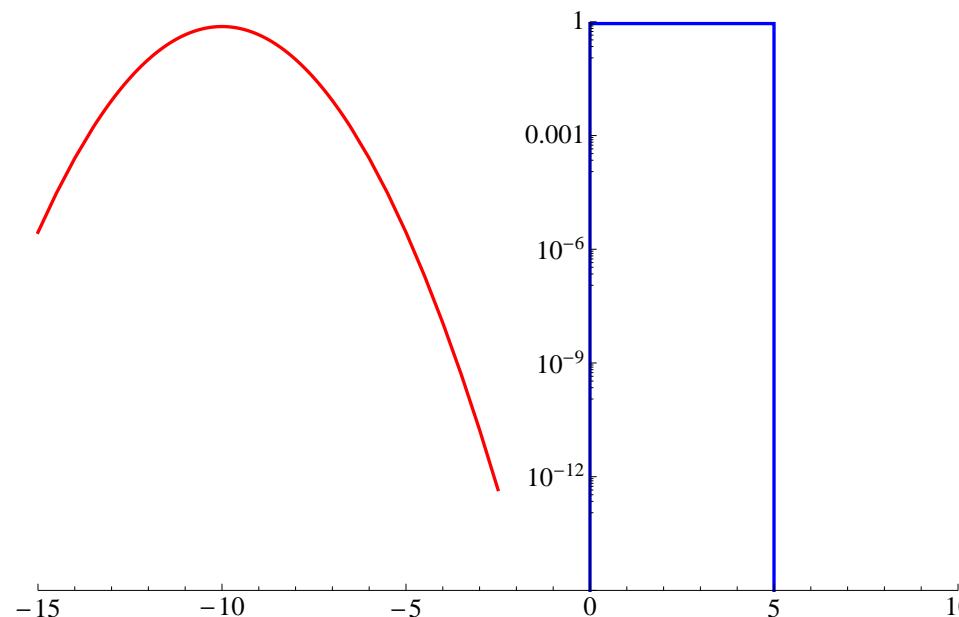
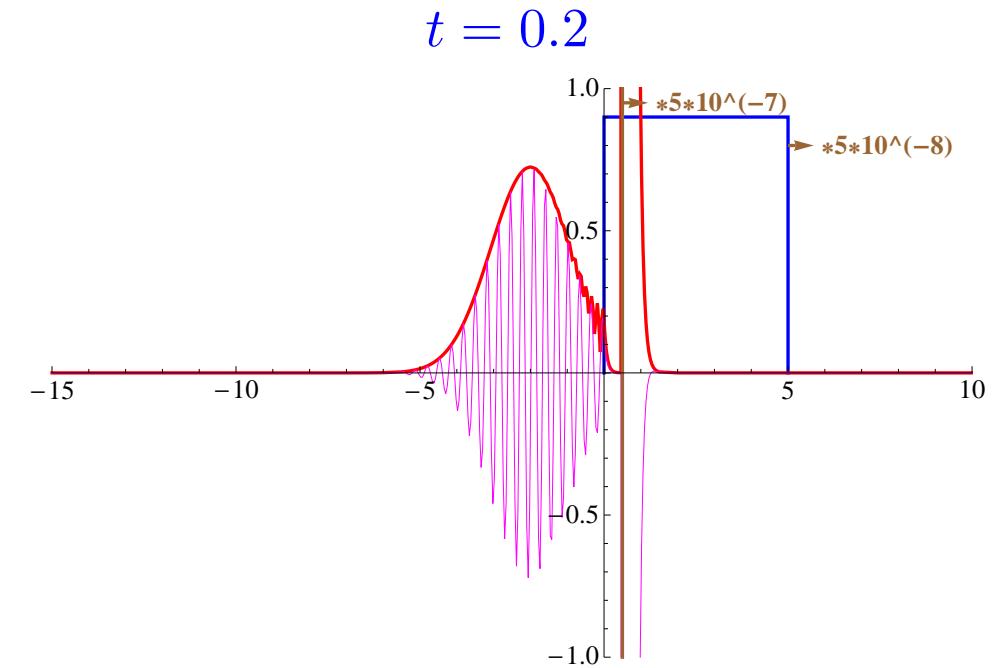
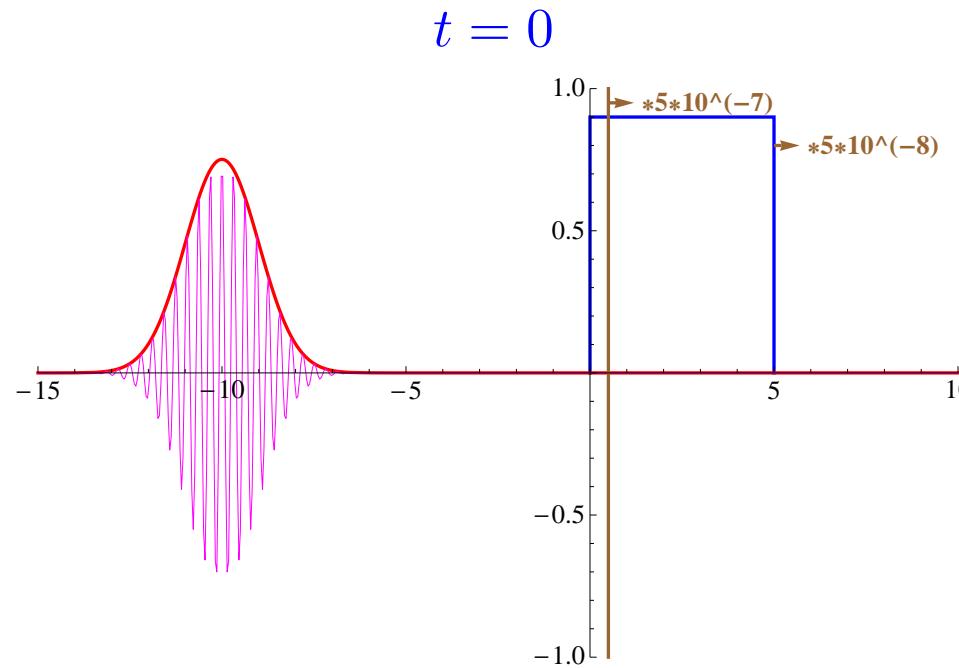
transmission dominated by

high momenta  $k \gtrsim \sqrt{V}$

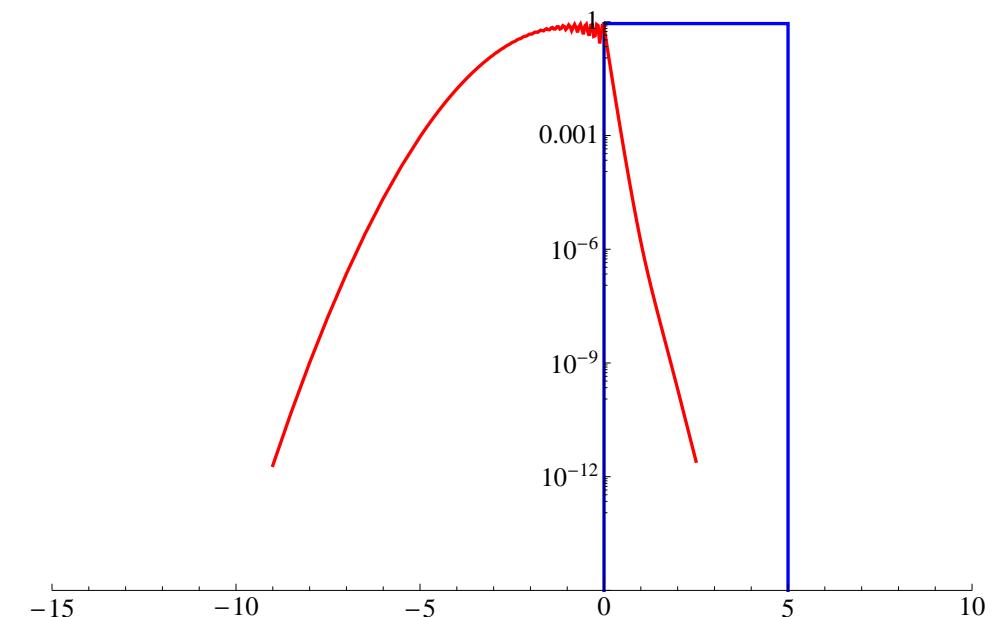
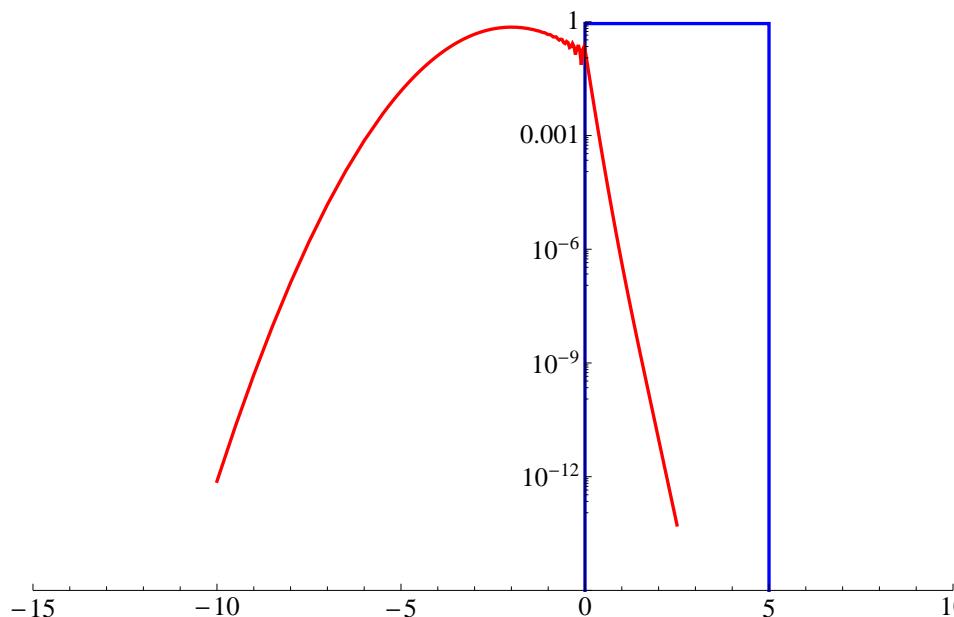
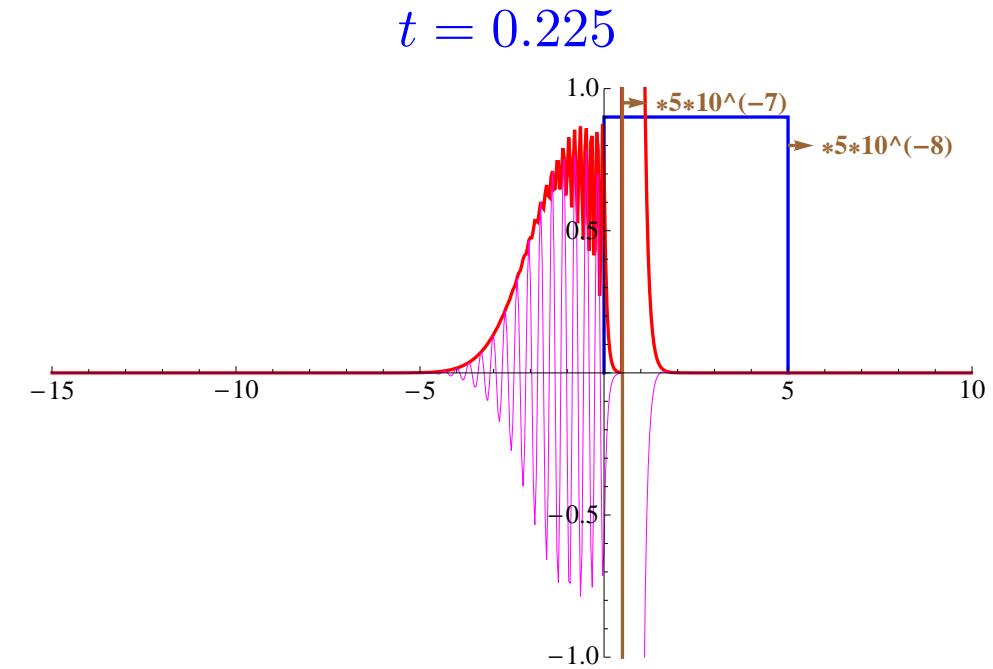
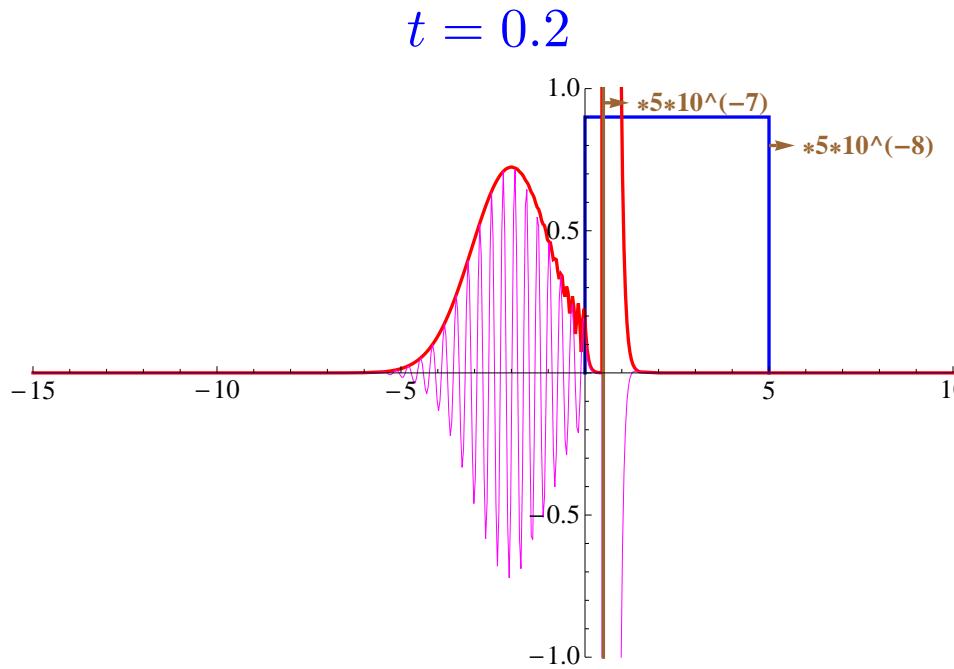
⇒ classically allowed transmission



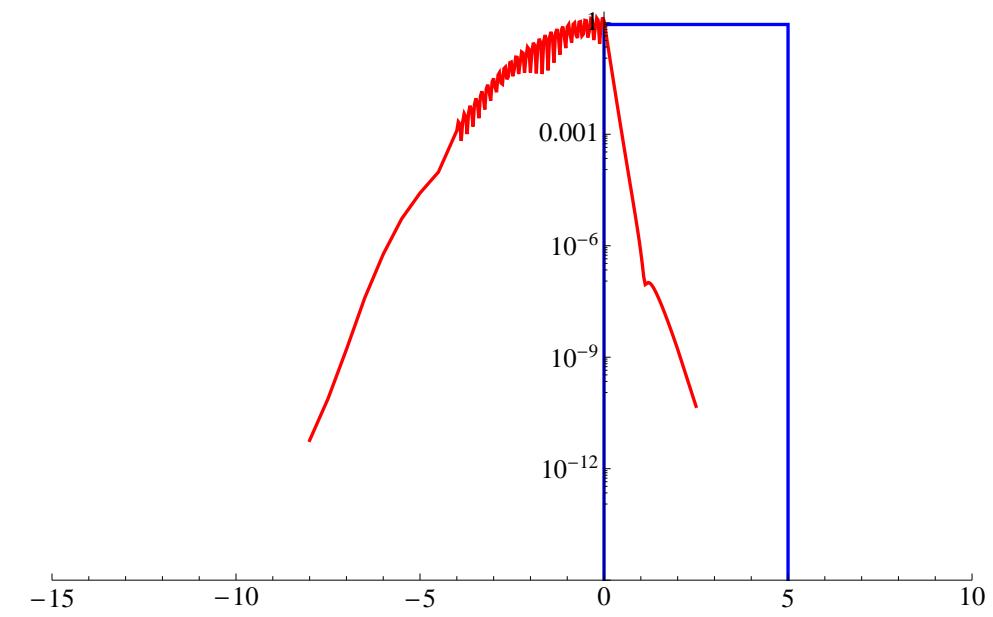
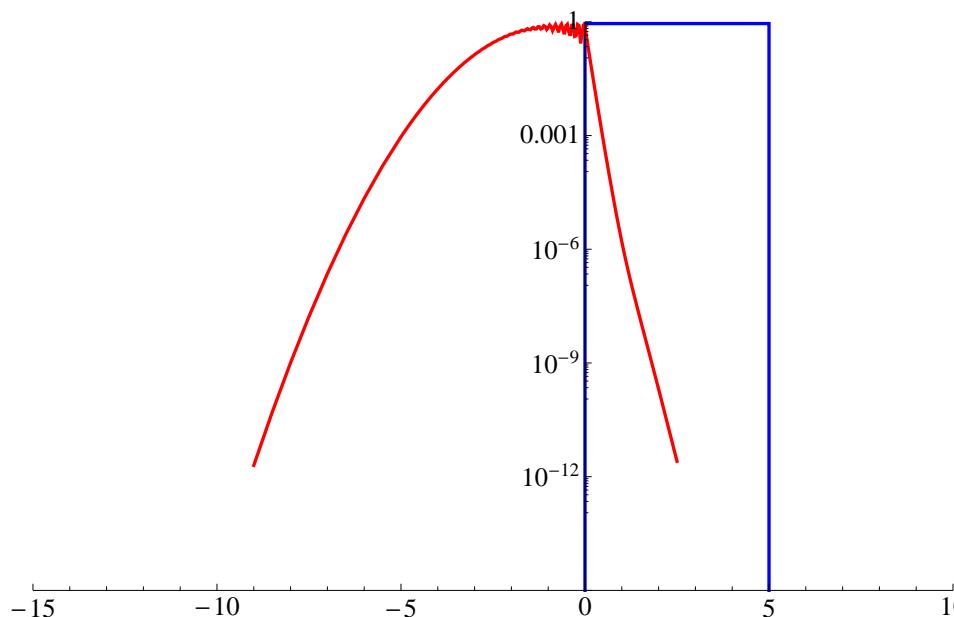
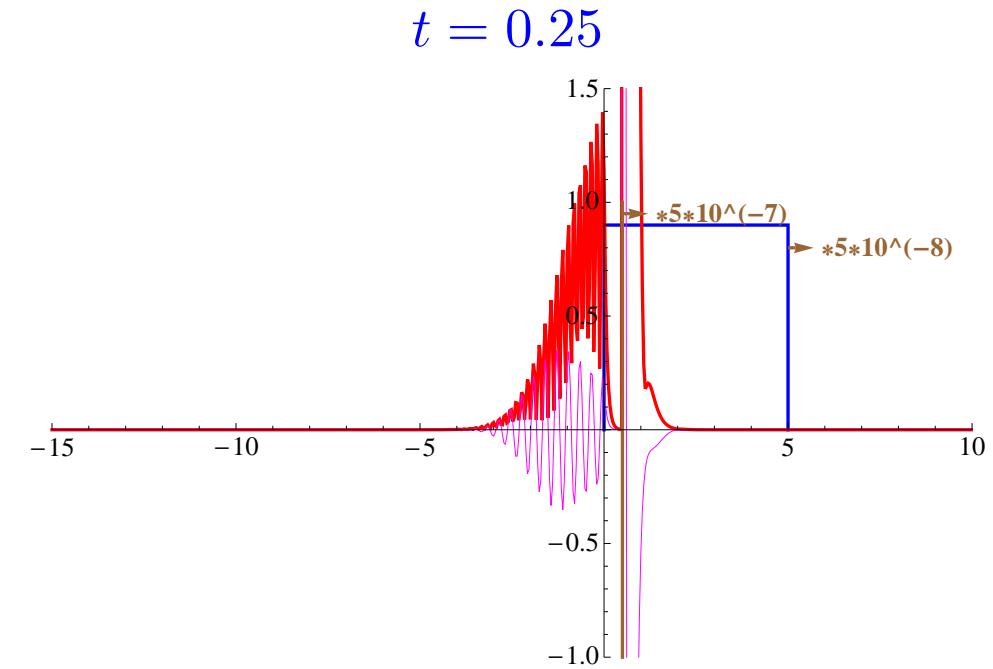
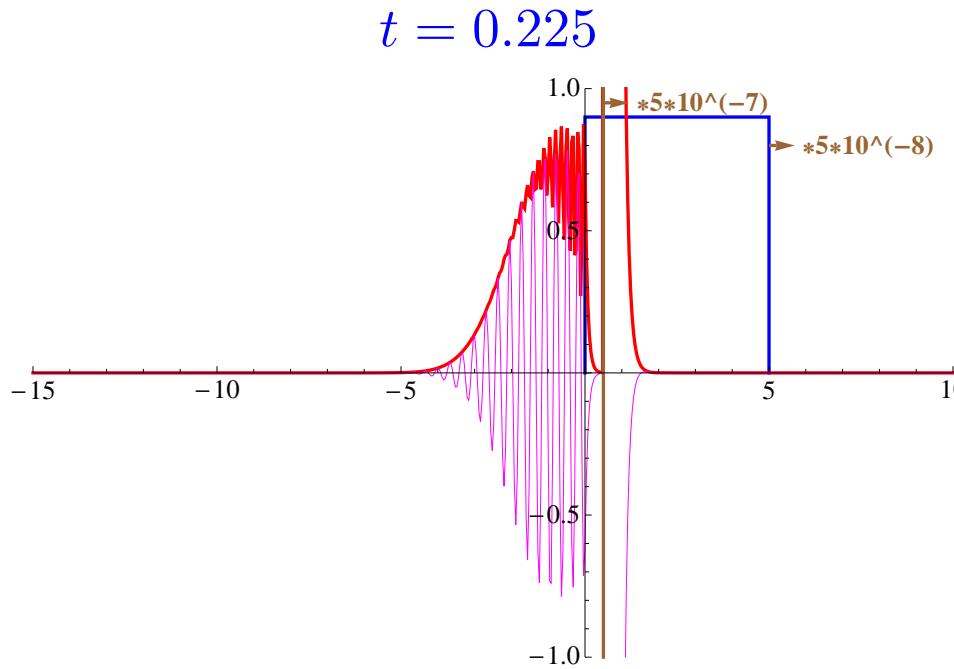
### C: “Fast” wave packet, thick barrier



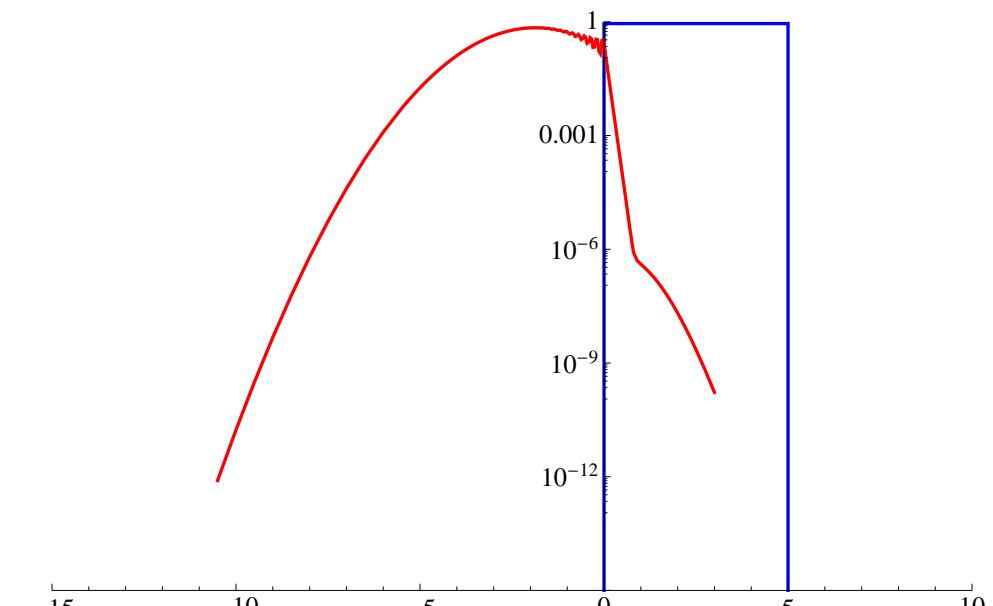
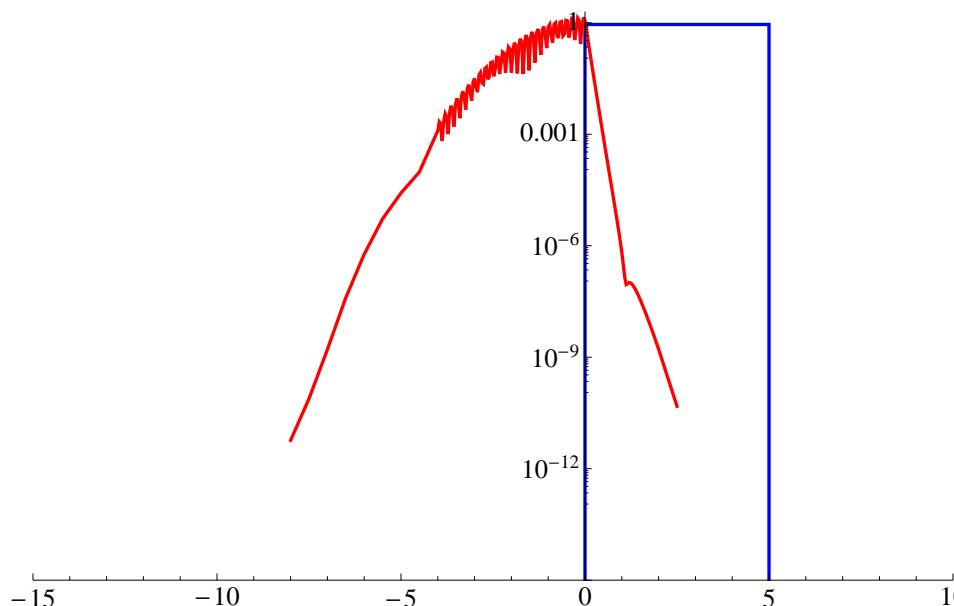
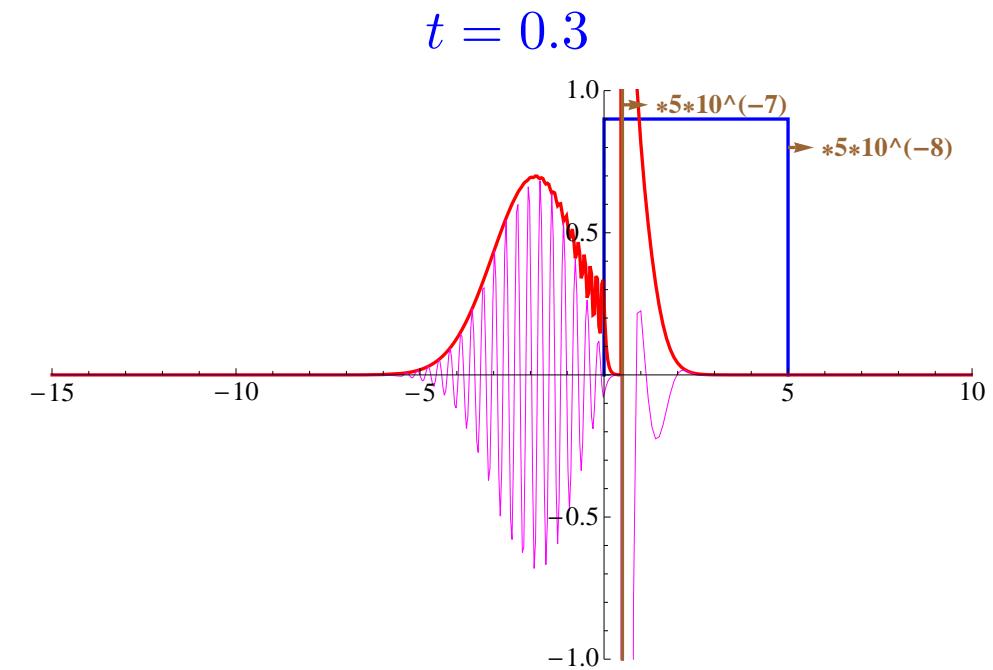
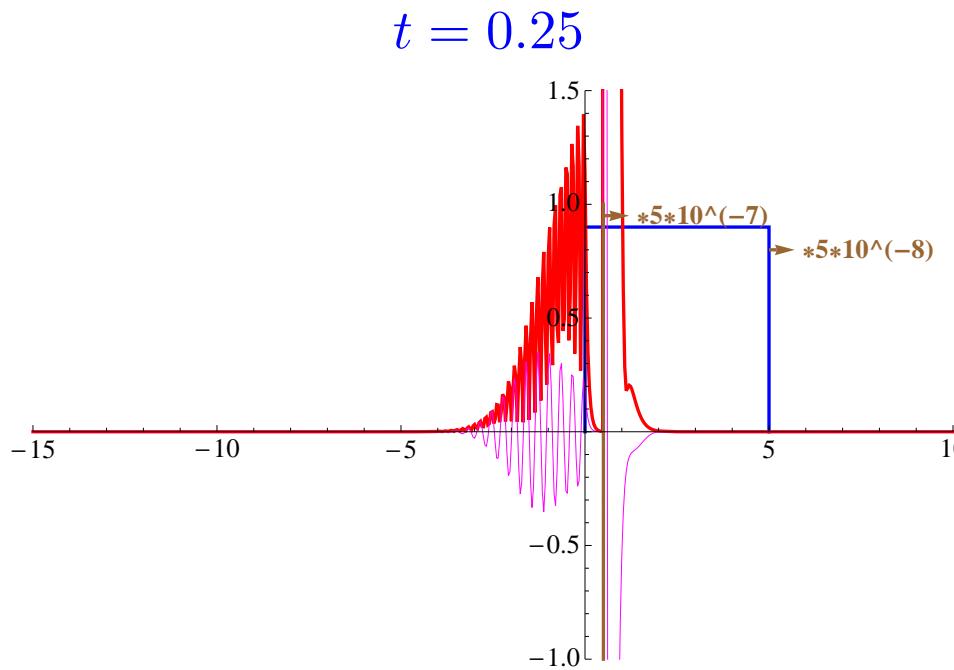
### C: “Fast” wave packet, thick barrier



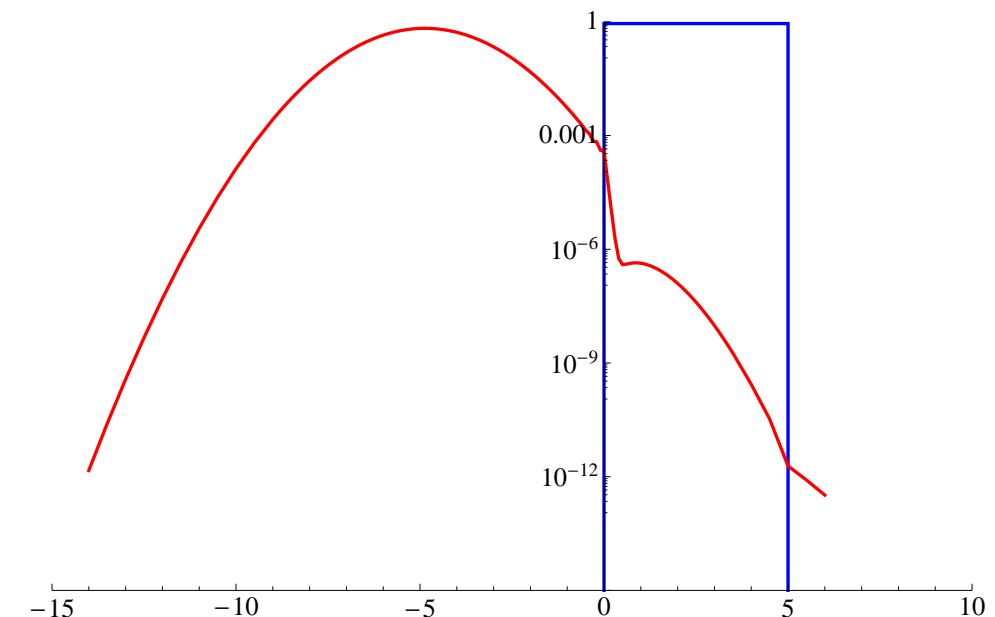
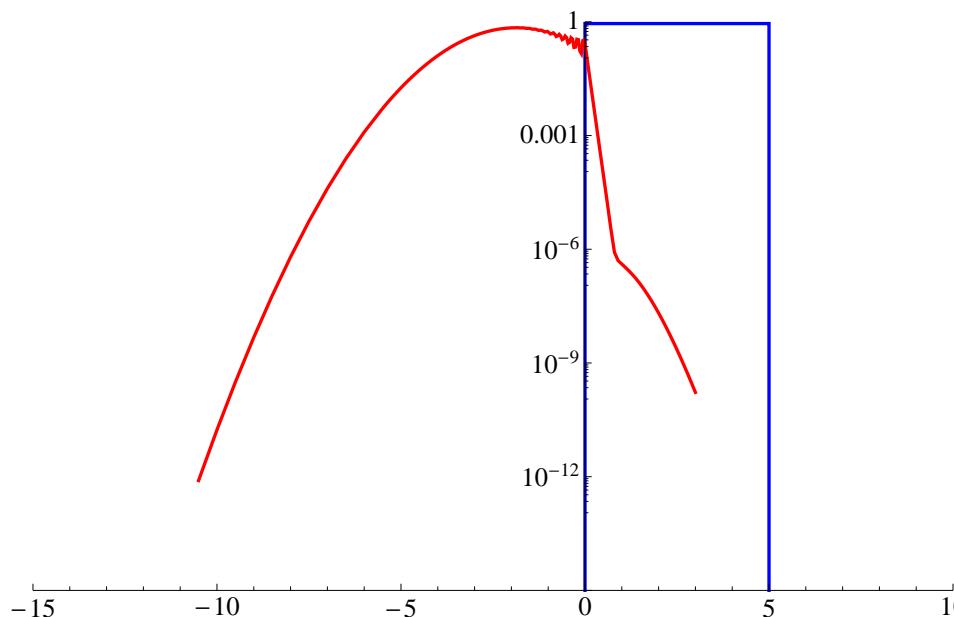
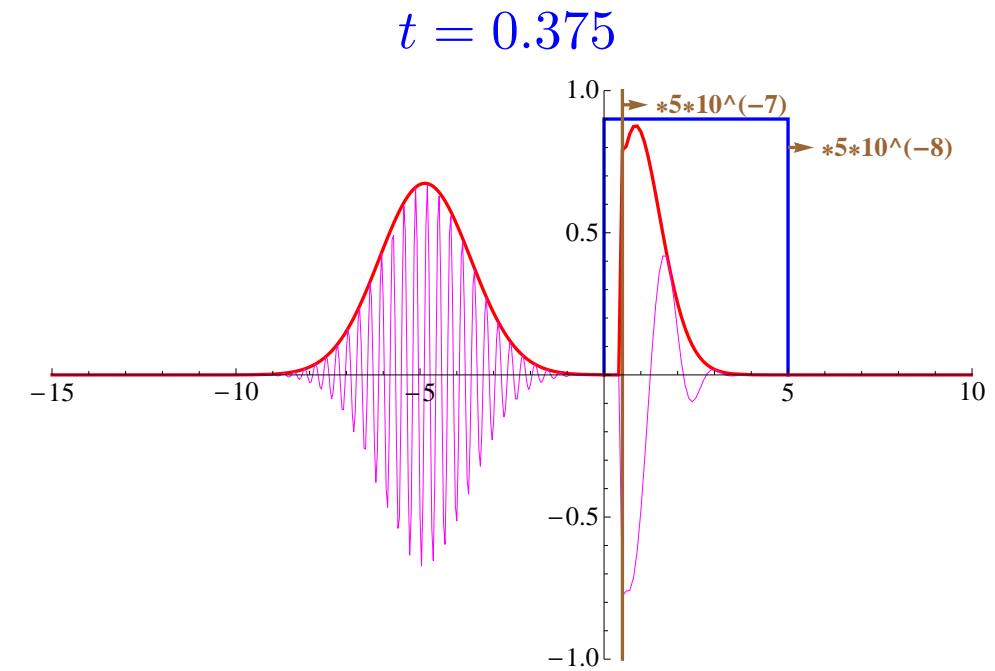
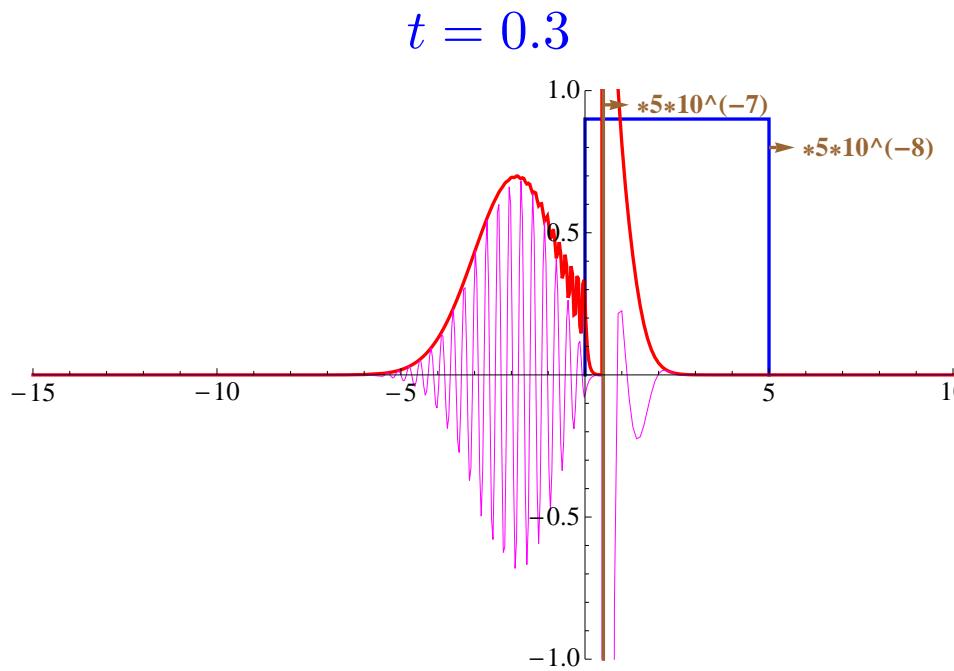
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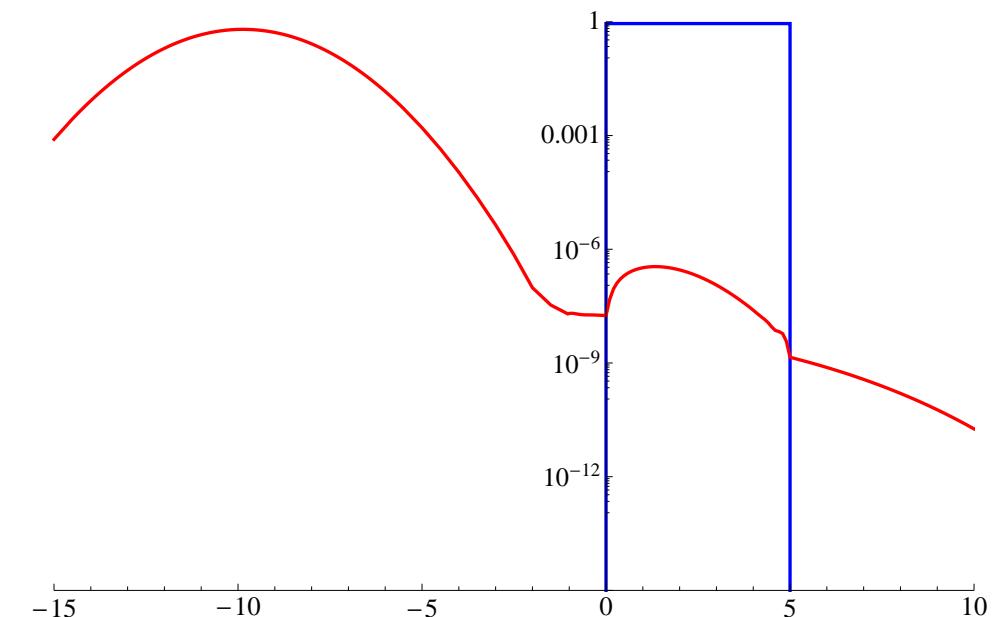
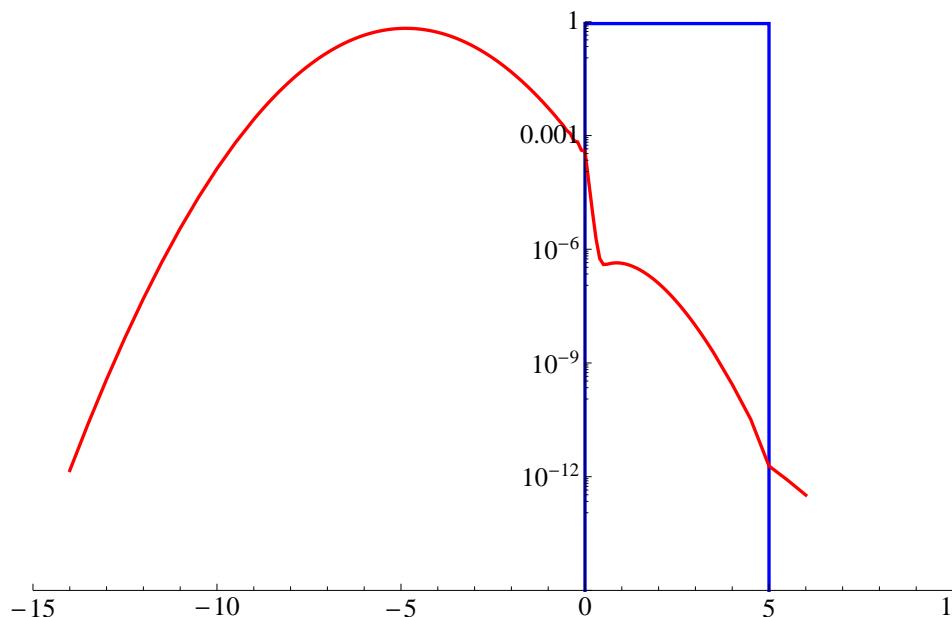
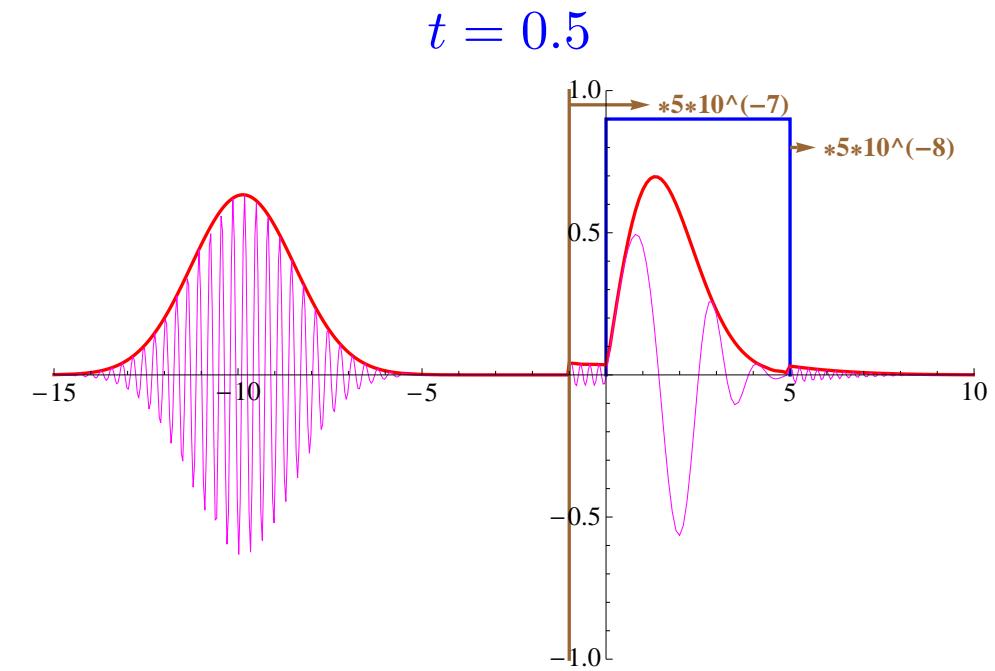
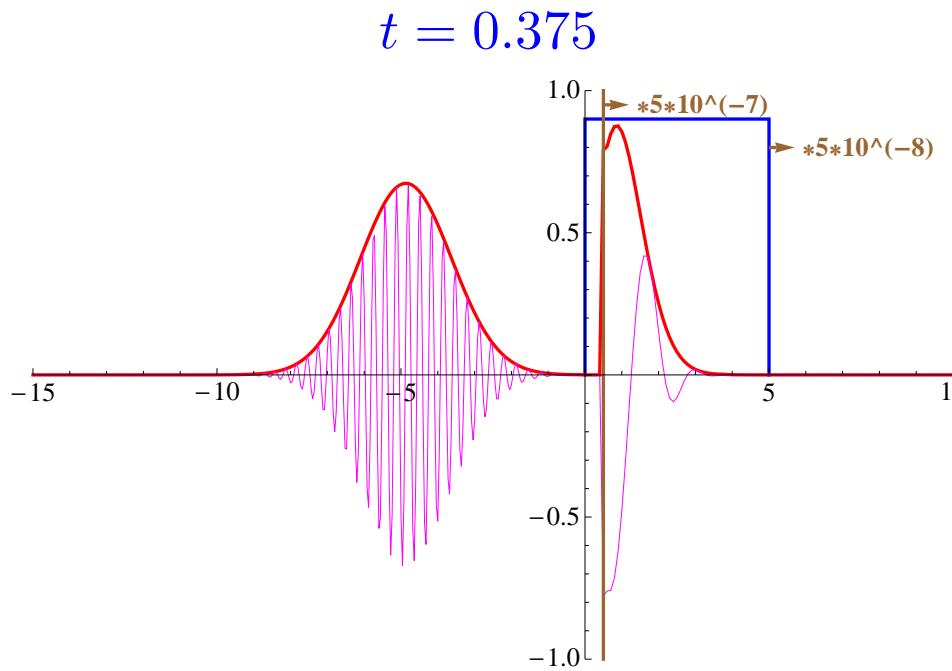
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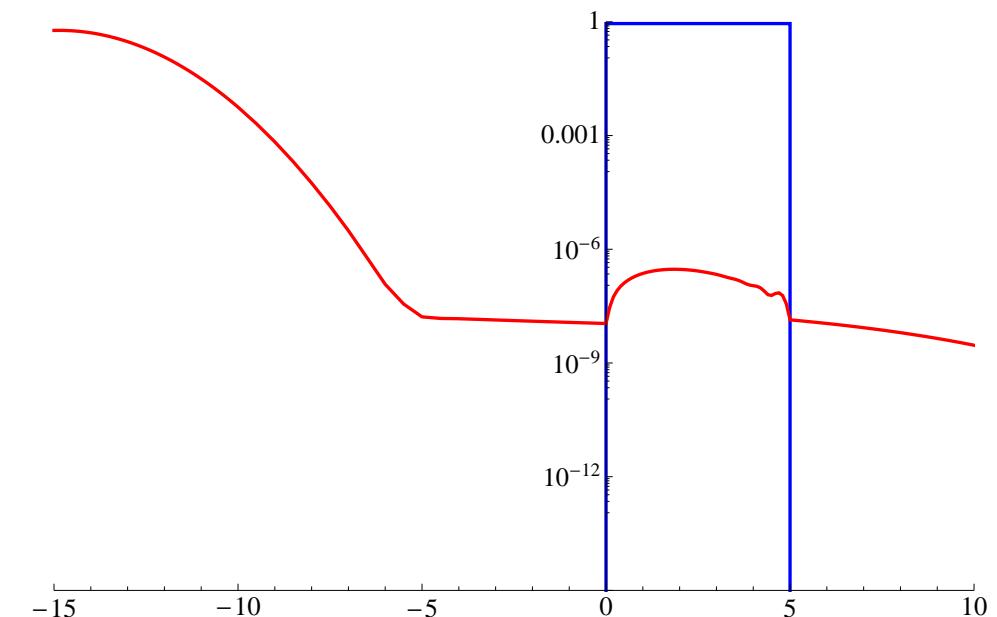
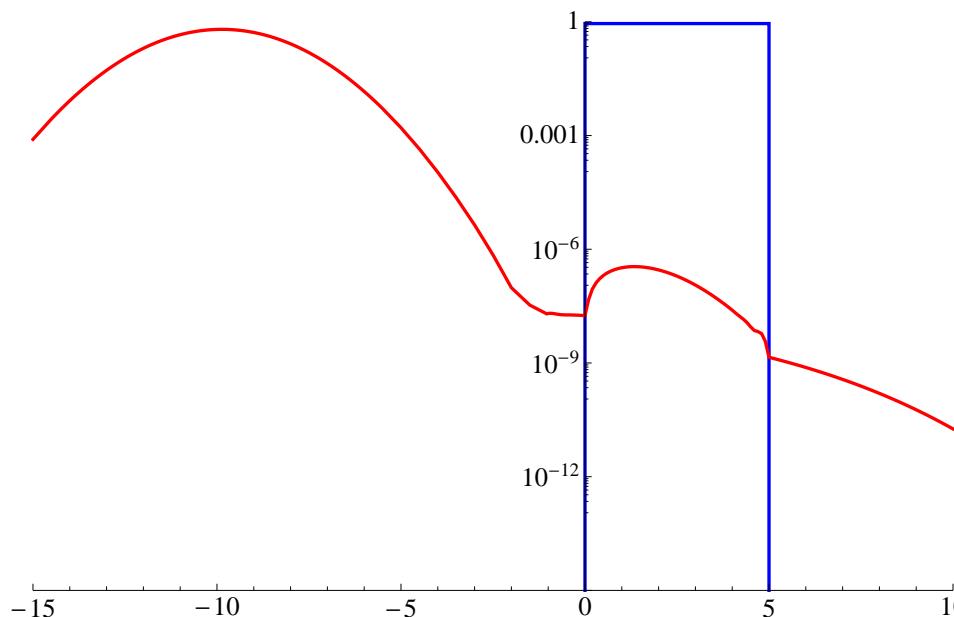
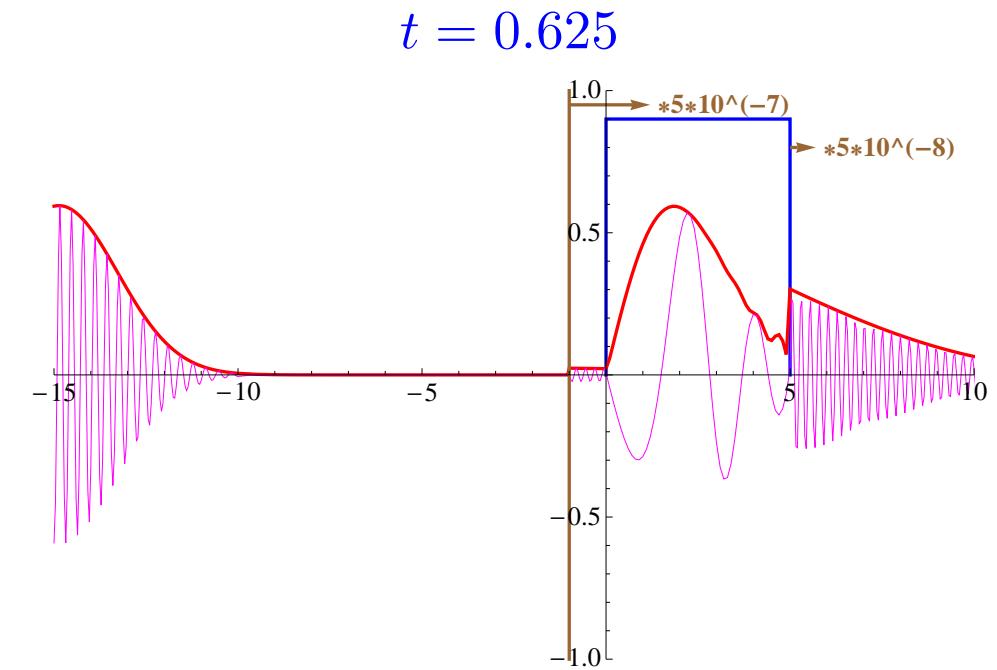
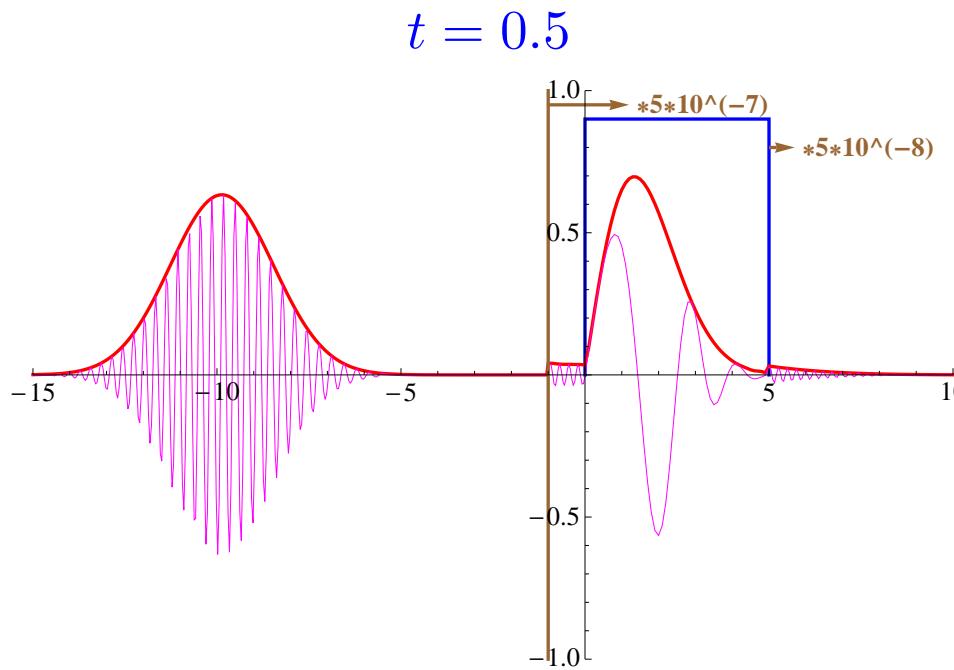
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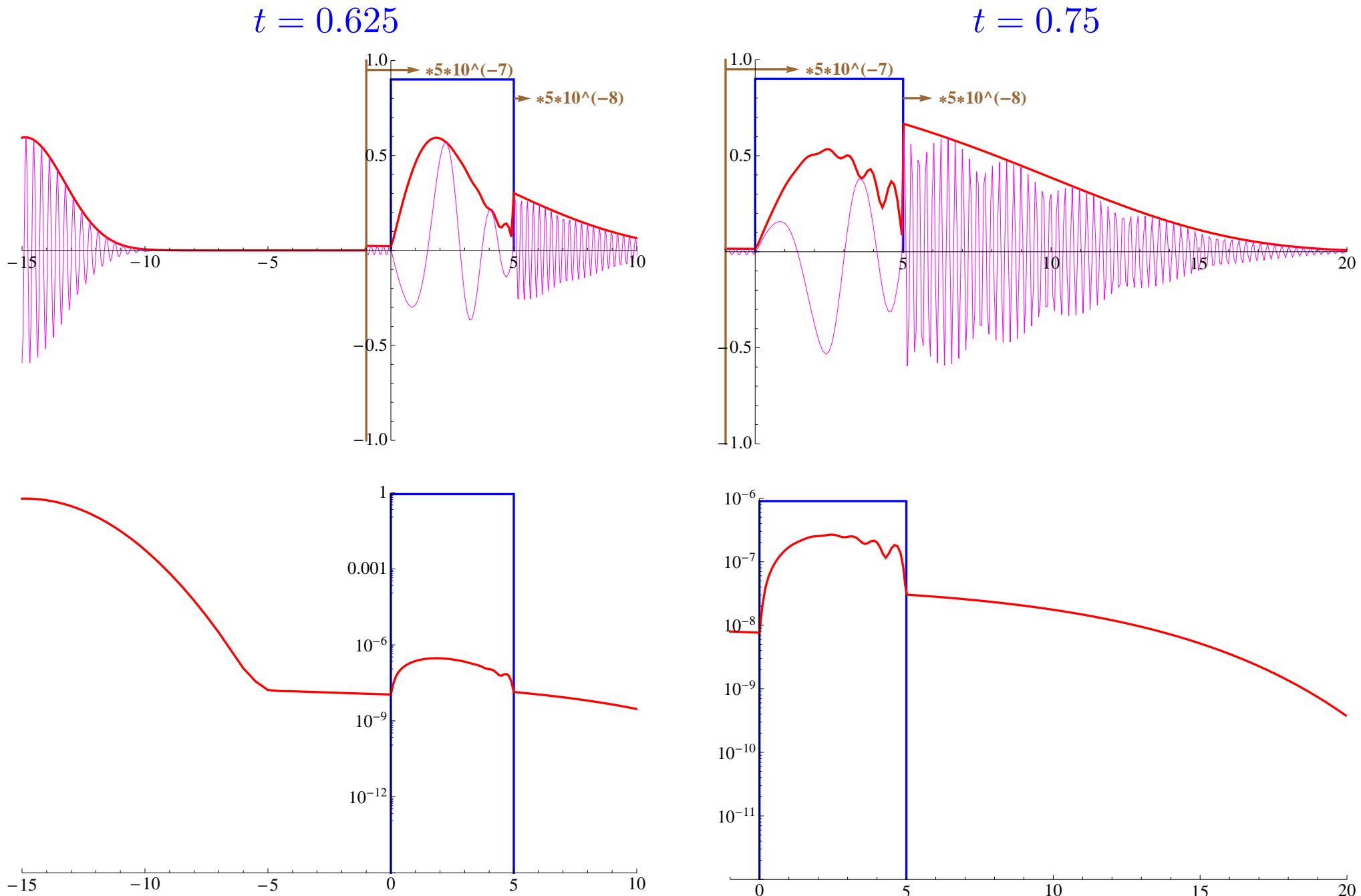
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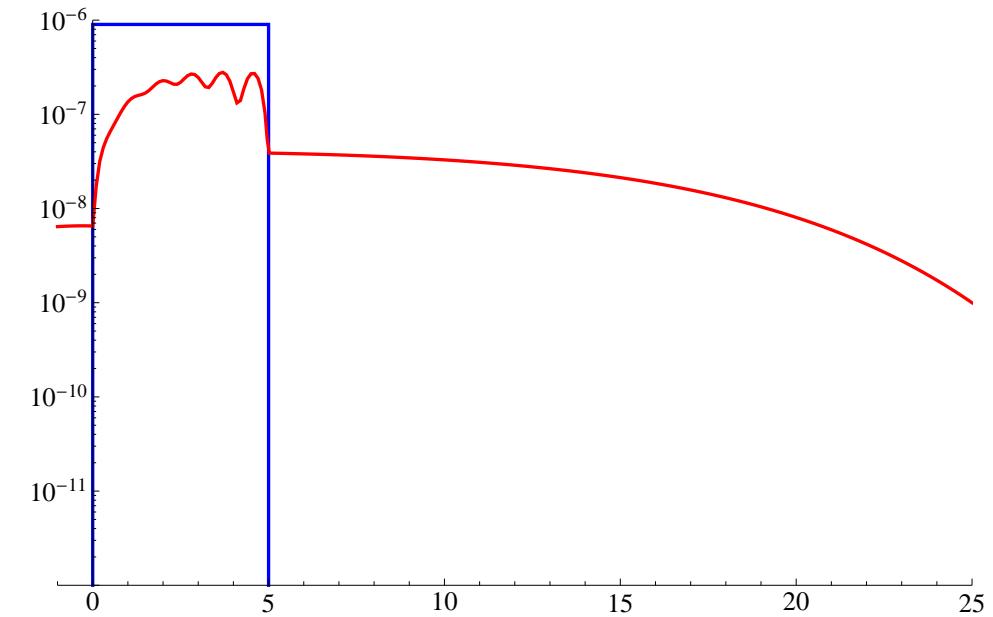
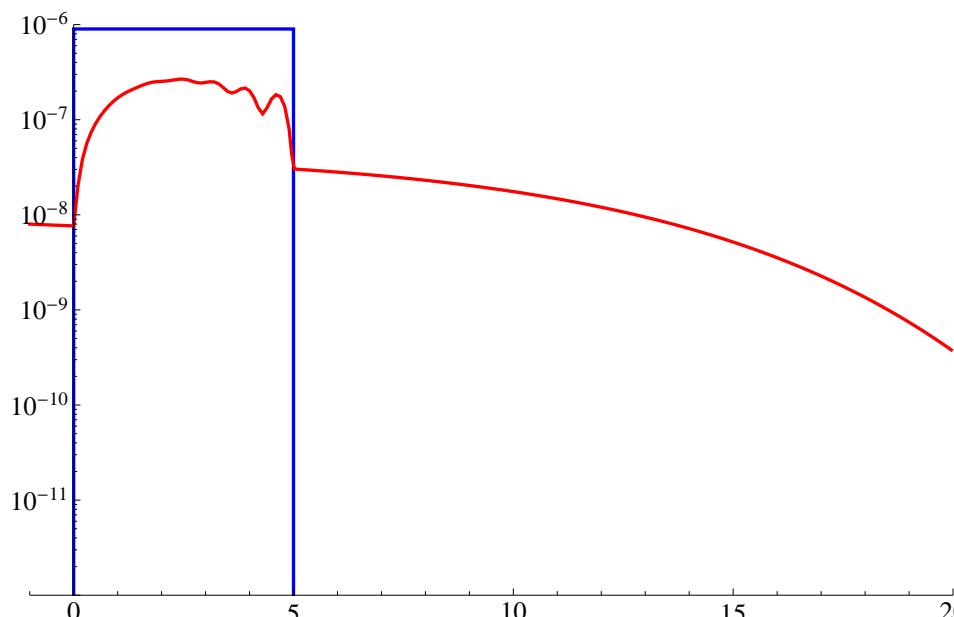
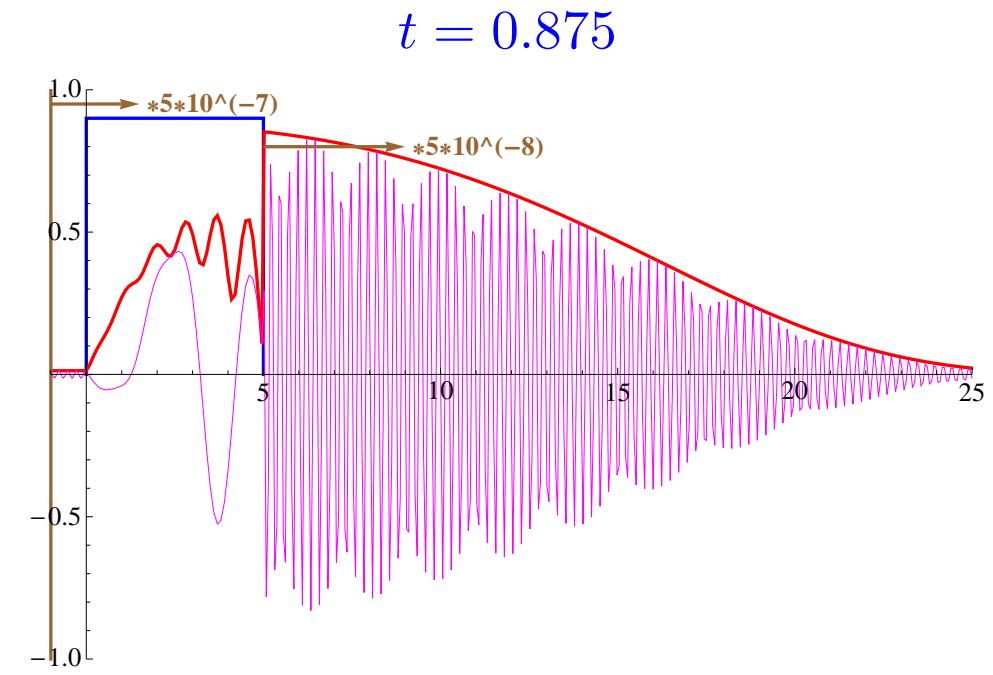
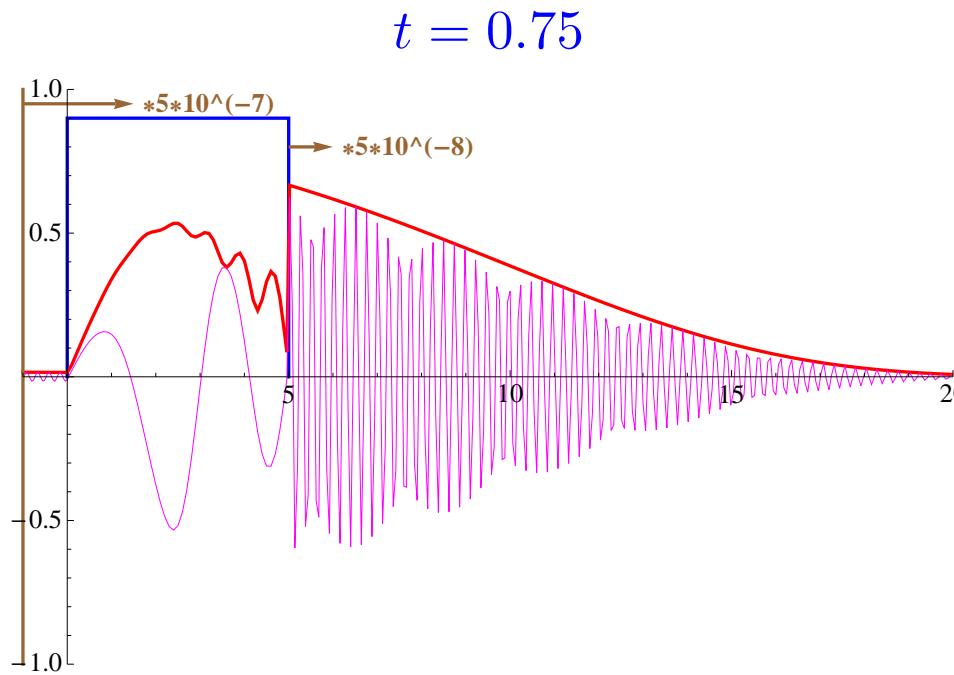
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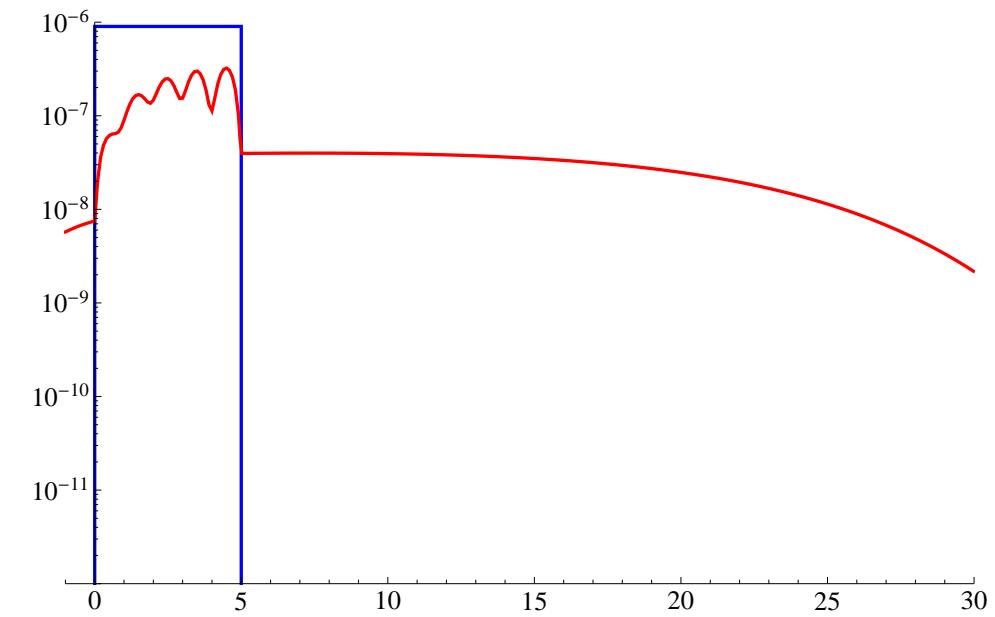
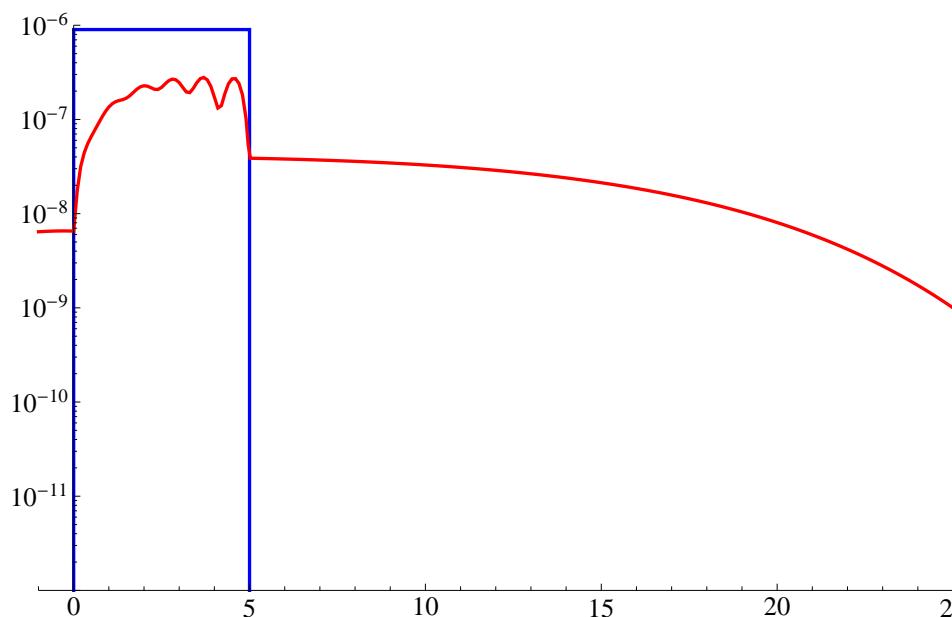
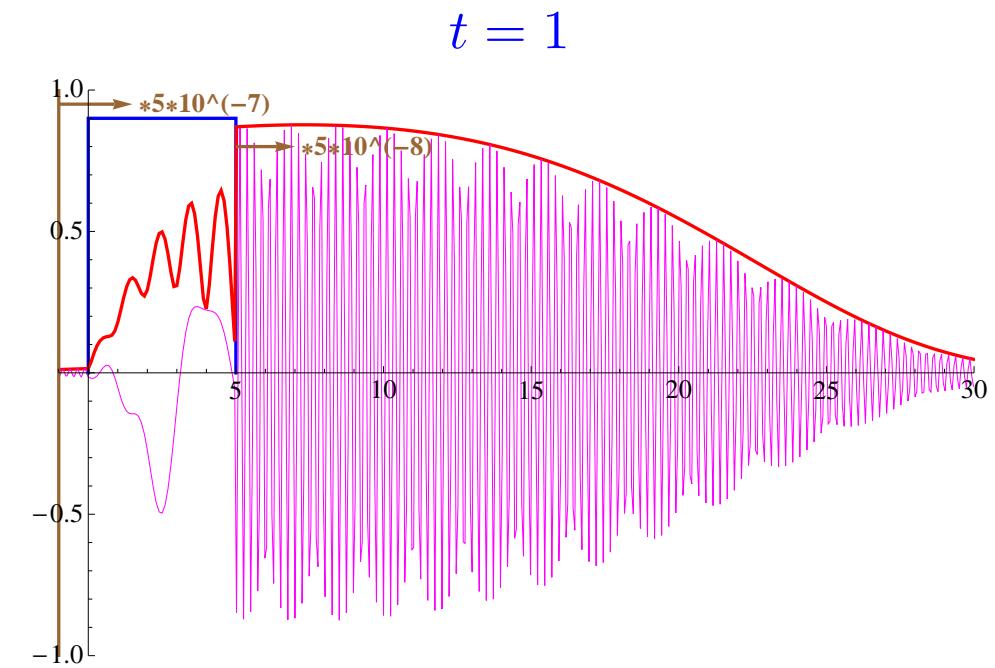
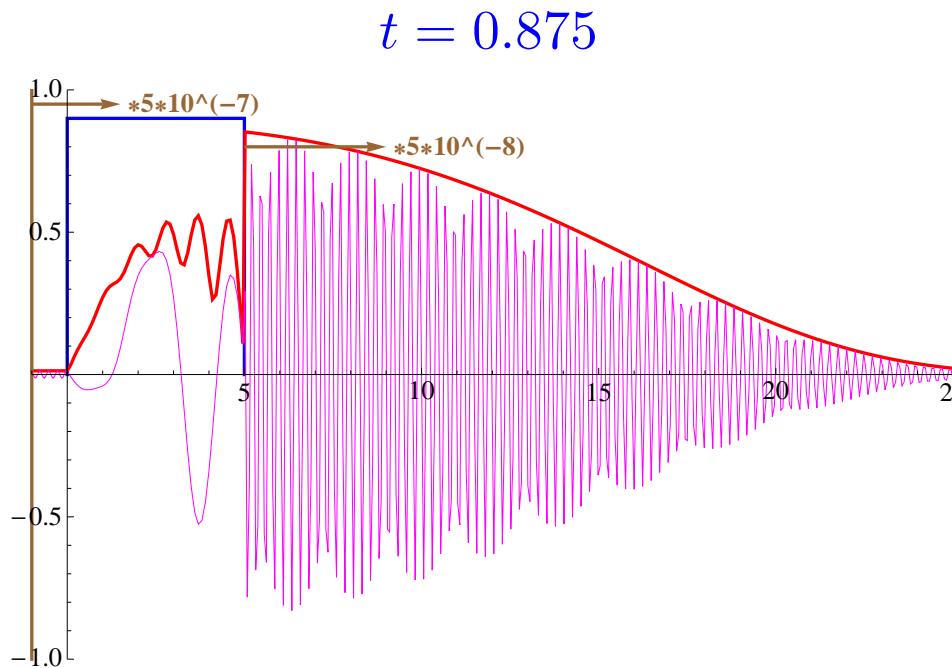
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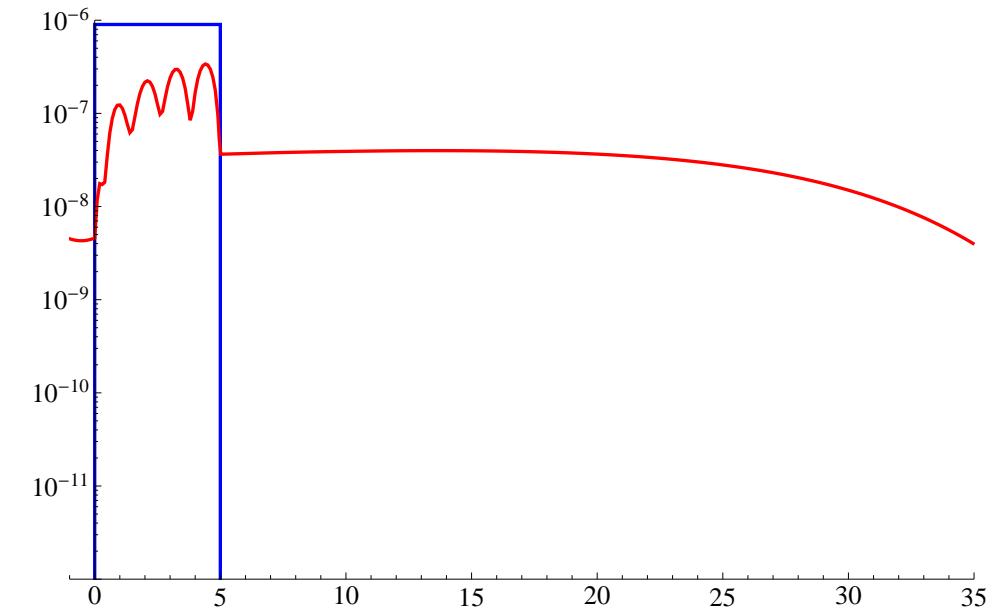
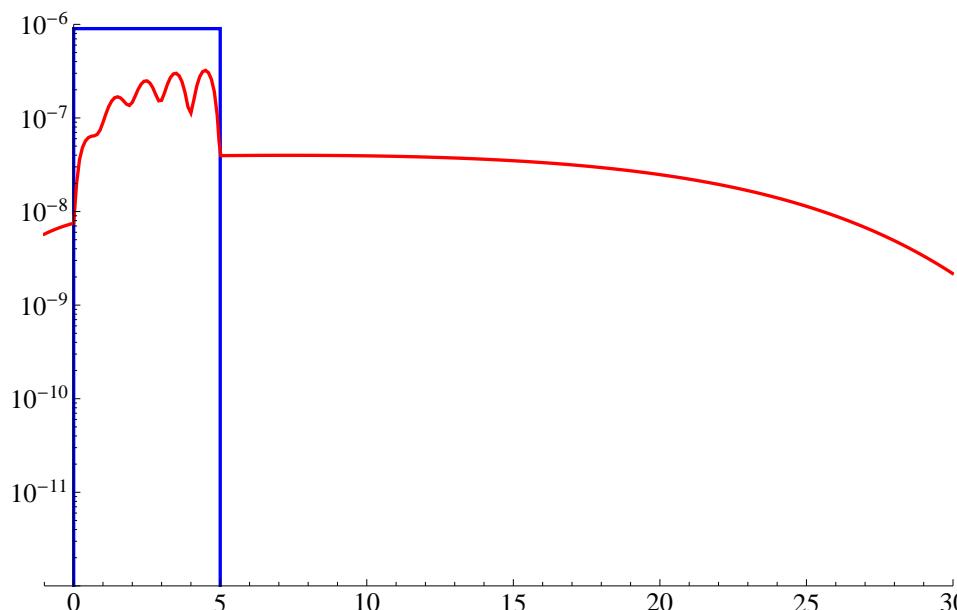
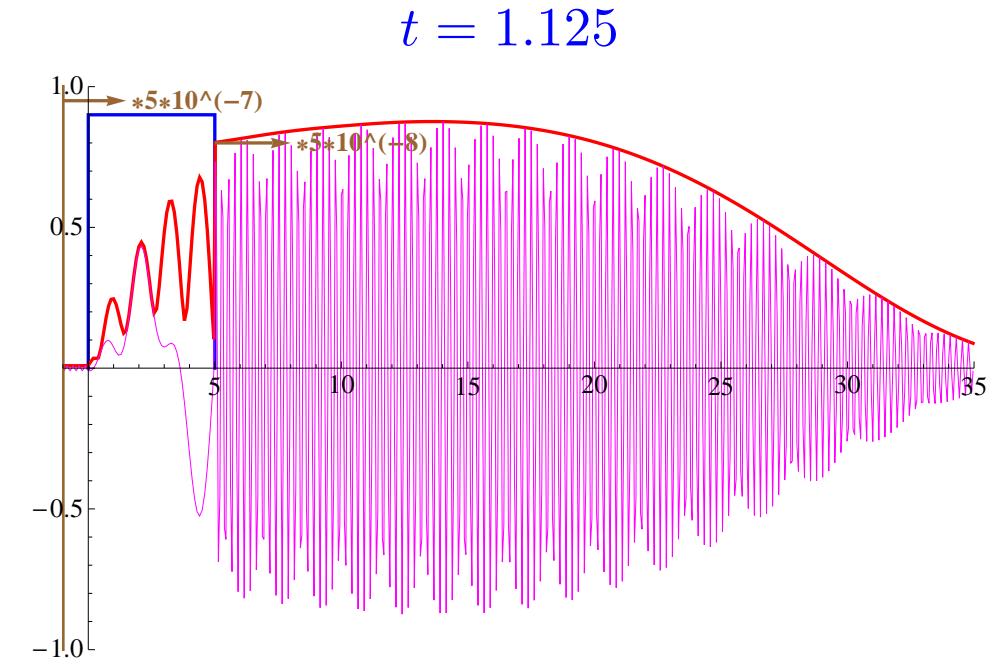
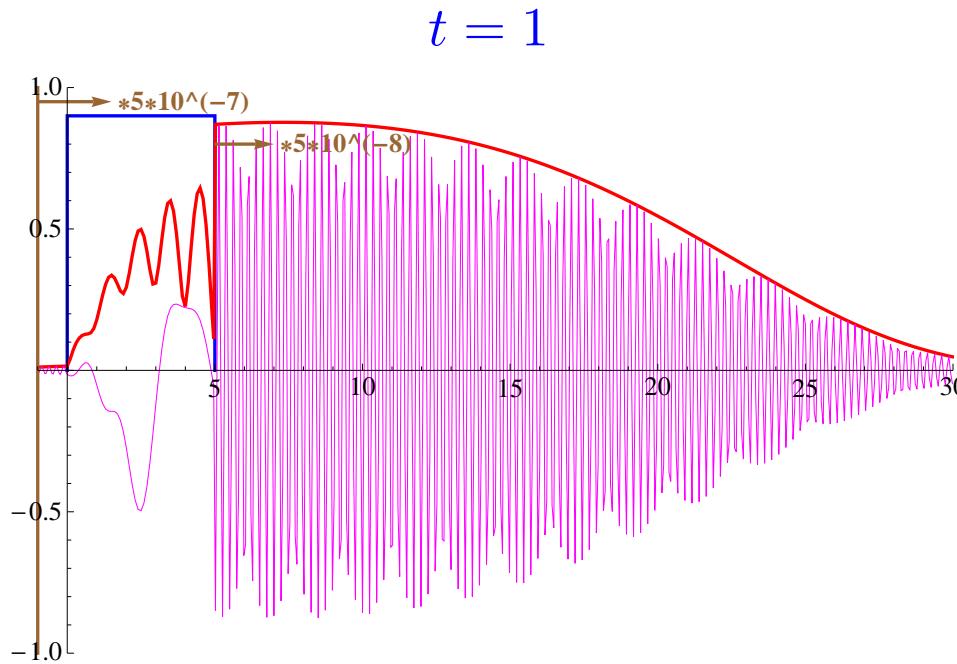
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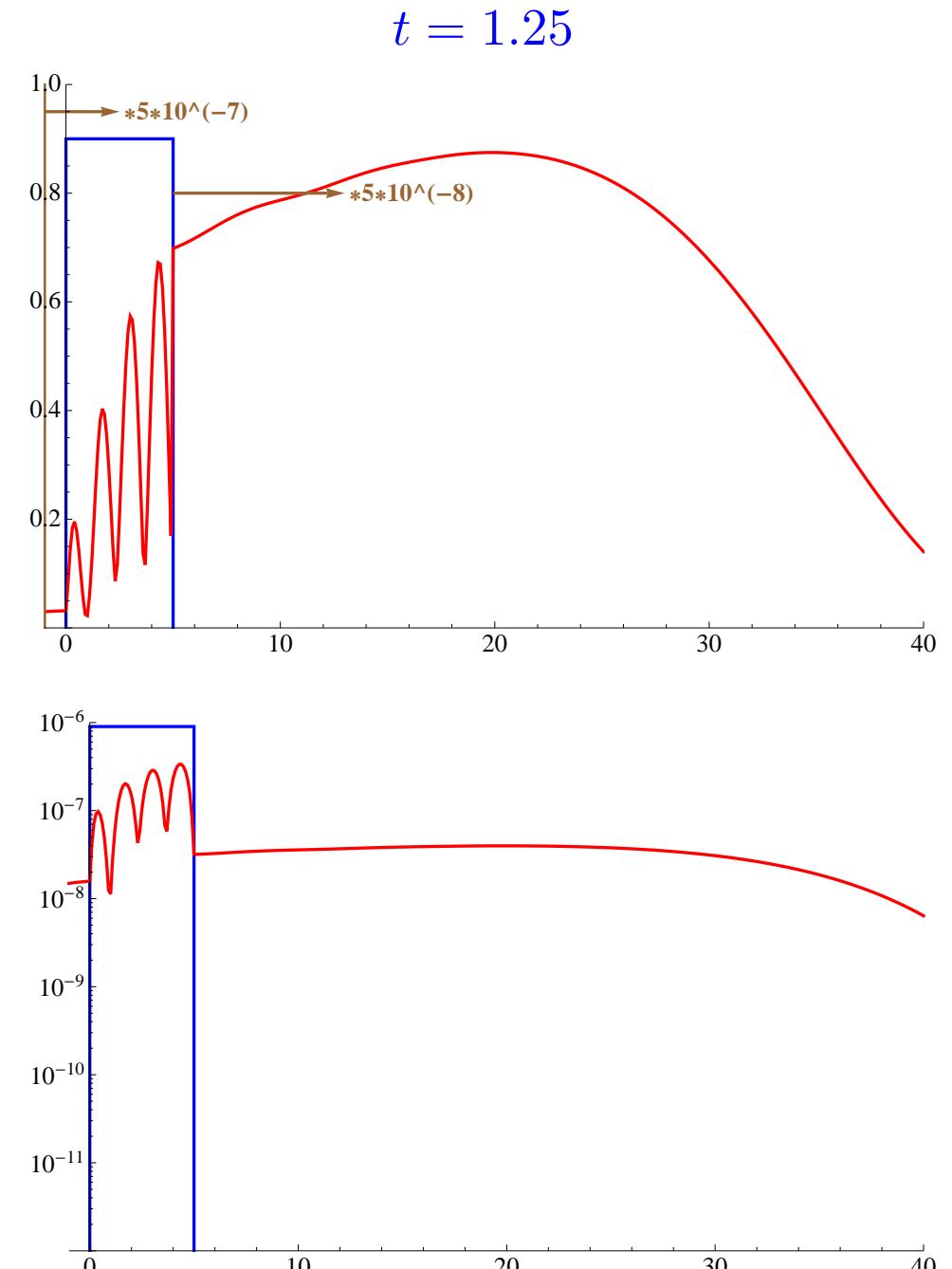
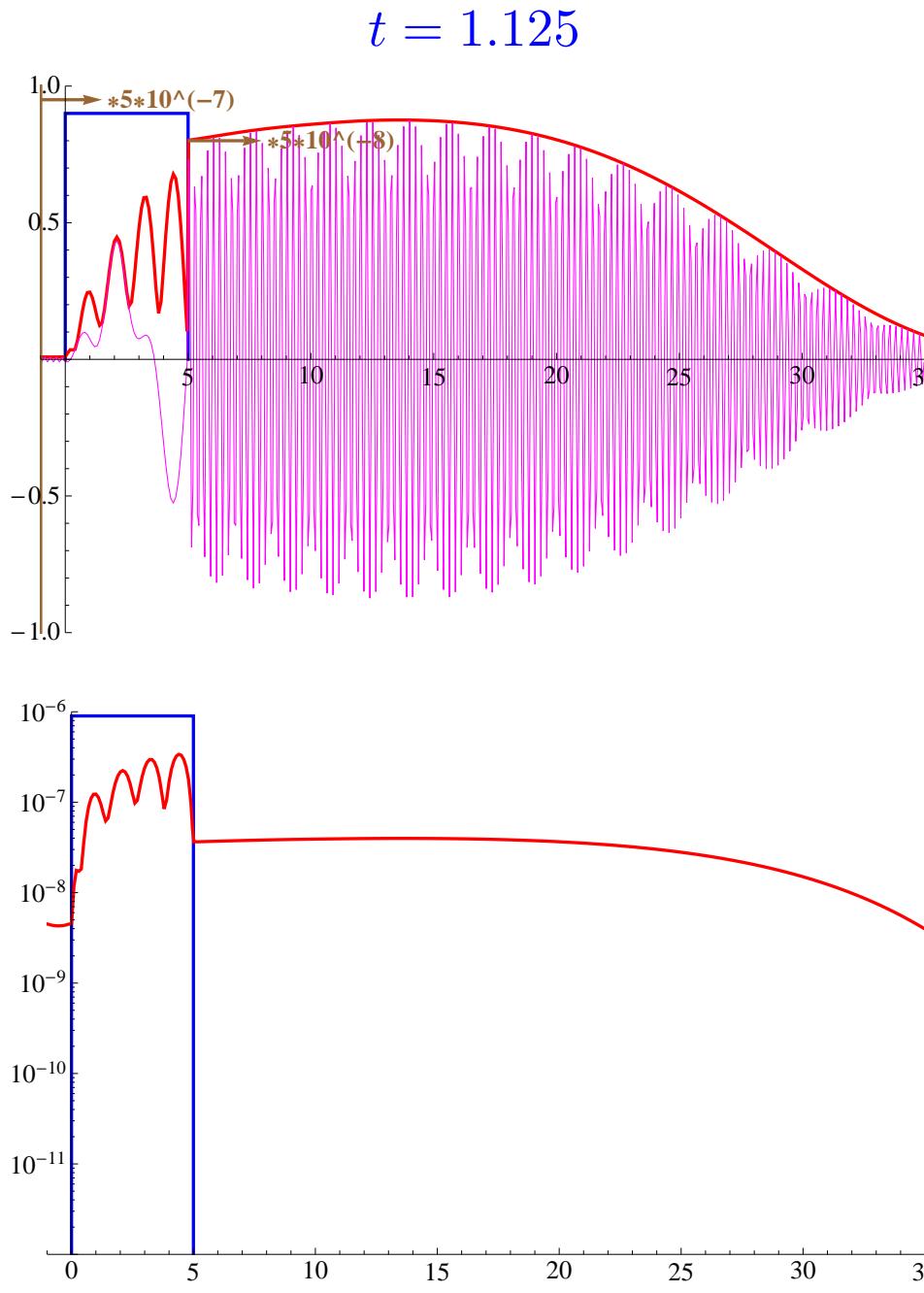
### C: “Fast” wave packet, thick barrier



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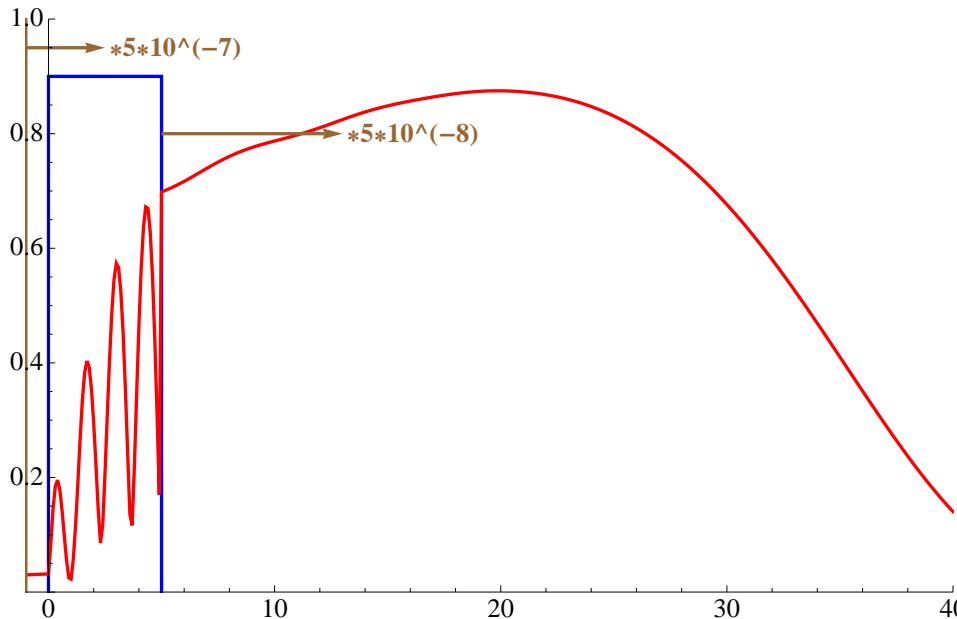


## C: “Fast” wave packet, thick barrier

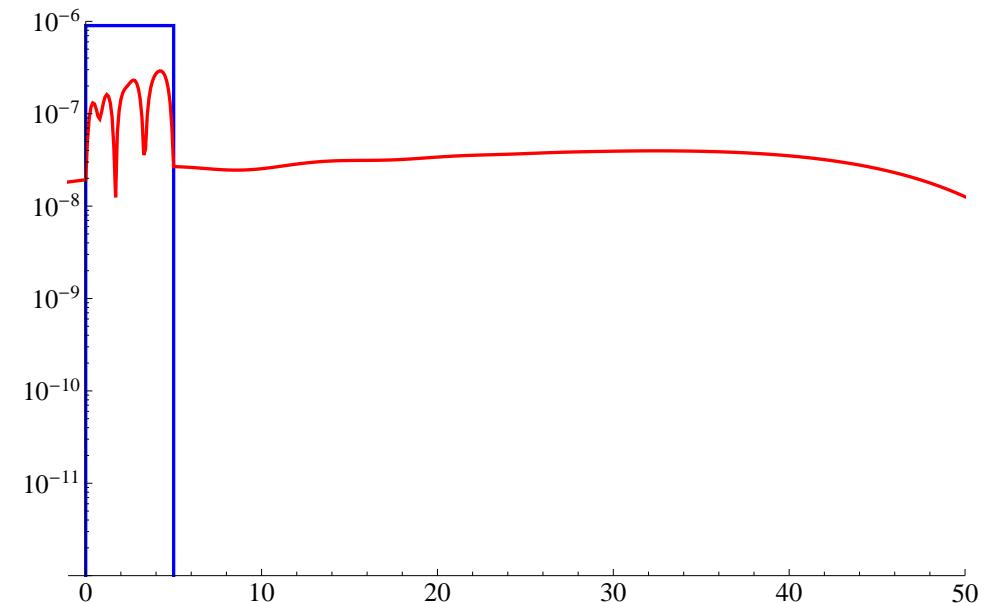
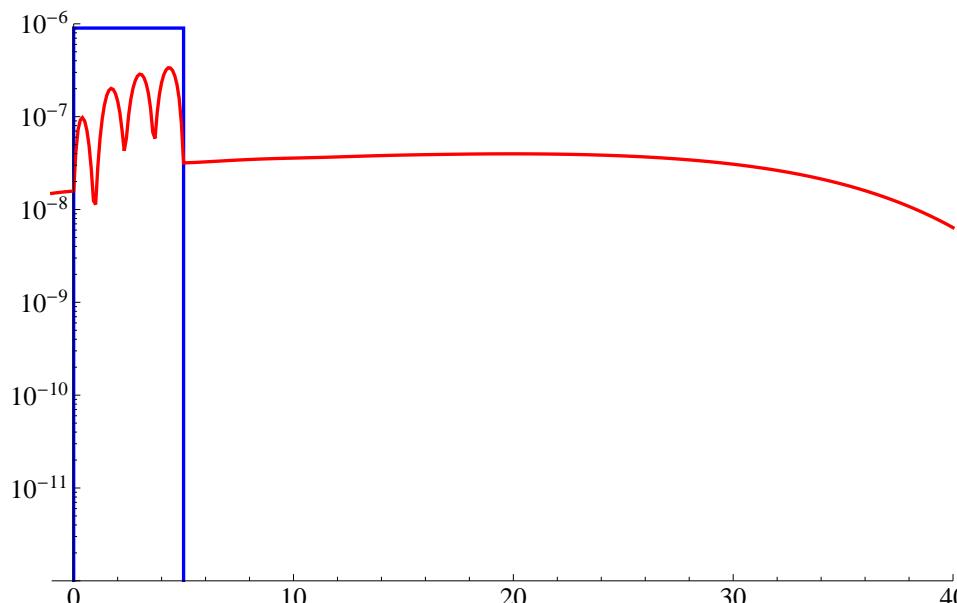
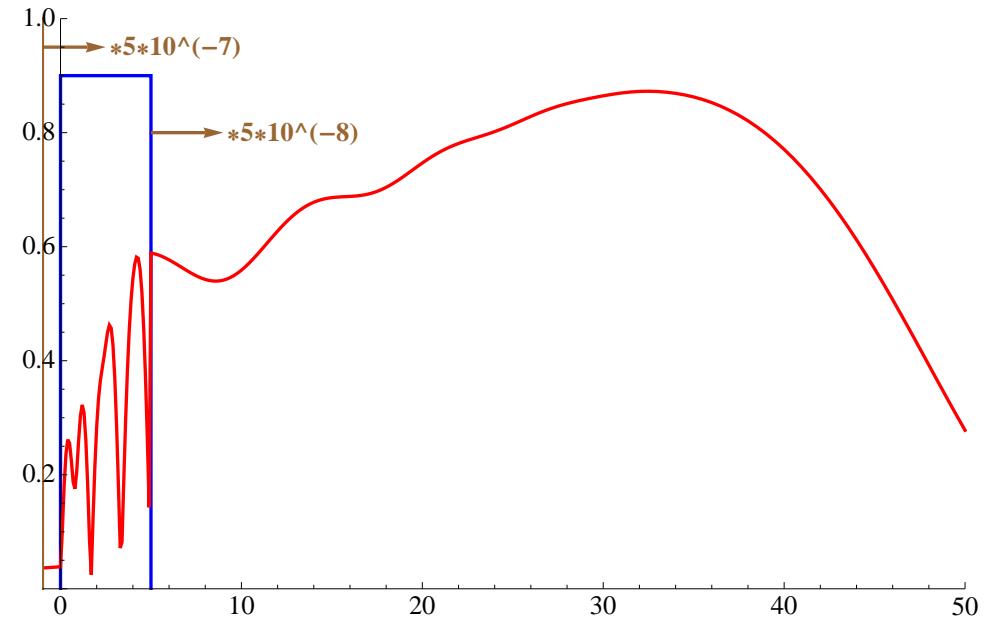


### C: “Fast” wave packet, thick barrier

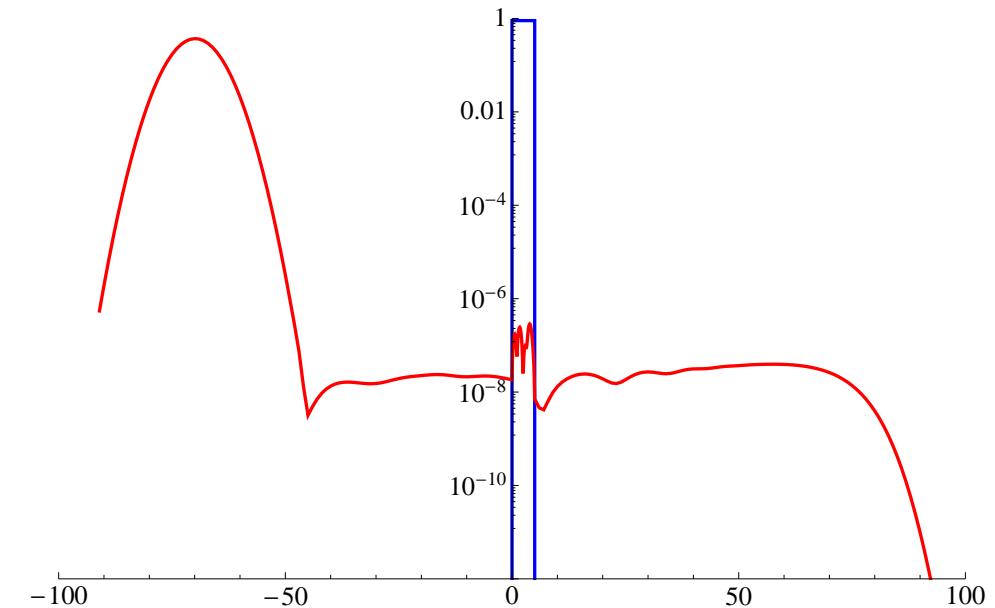
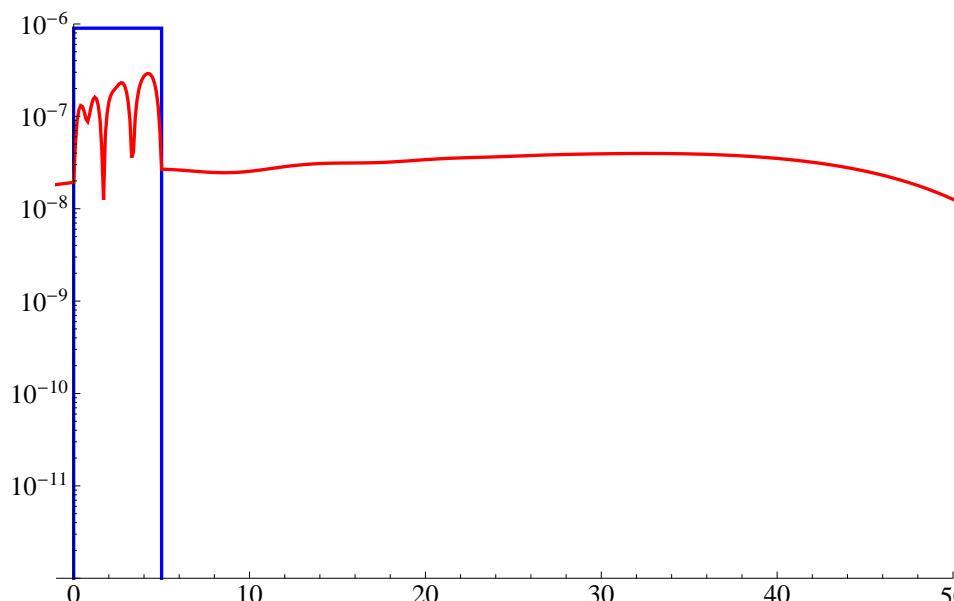
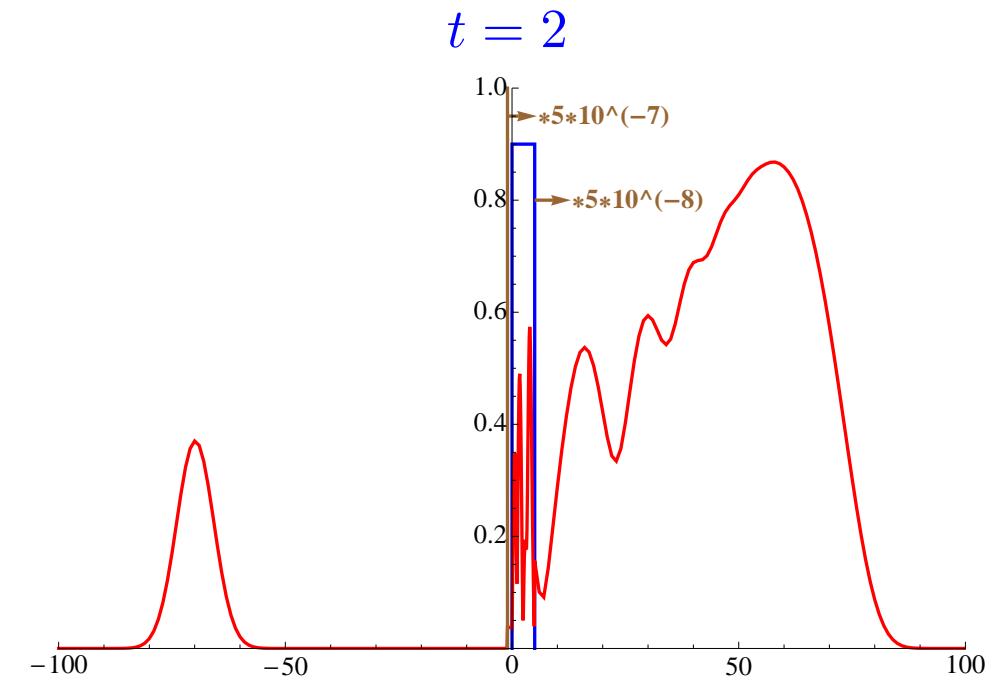
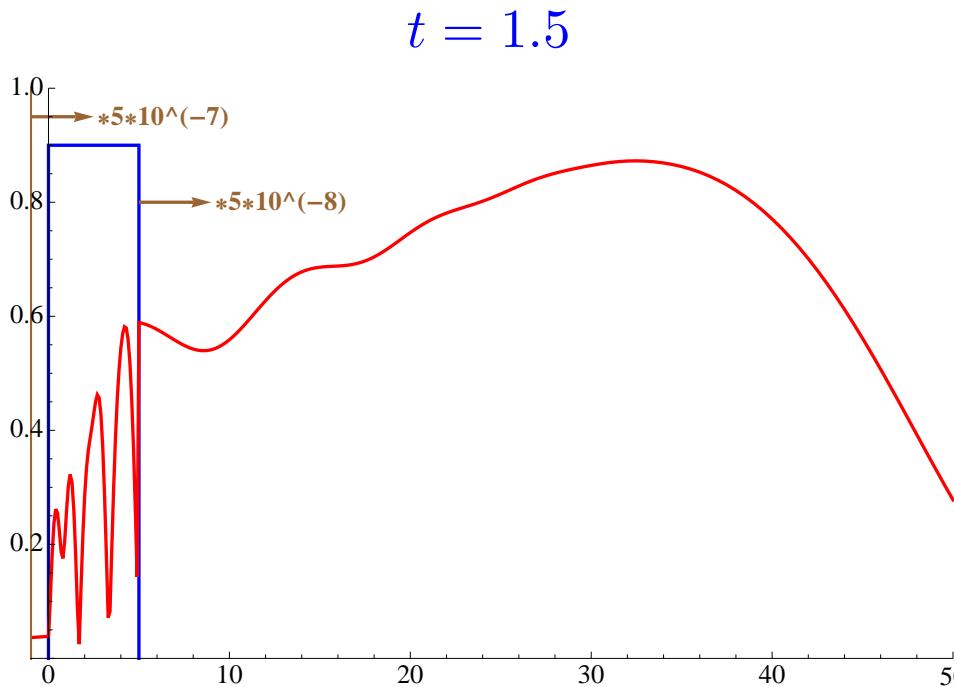
$t = 1.25$



$t = 1.5$



## C: “Fast” wave packet, thick barrier



## C: “Fast” wave packet, thick barrier — Conclusion

- “Fast” wave packet:  $\sigma_k/k_0$  smaller, (almost) no low momenta  
    → wave packet stays more compact ⇒ better time resolution
- Thick barrier: like in case B – strong exponential damping ( $\sim 10^{-30}$  at  $k_0$ )  
    → dominated by classically allowed transmission of high momenta ( $\sim 10^{-7}$ ).
- Incoming and reflected wave packet: like in cases A&B – wave length and velocity of maximum correspond to  $k_0 = 20$  and to a reflection at  $x = 0$ .
- In the tunnel: a wave packet with moving maximum and large wave length appears, extends, is reflected from end of tunnel, develops fringes, becomes a standing wave
- Transmitted wave packet: wave length and velocity of maximum  
 $\rightsquigarrow k \approx 25 = \sqrt{V} > k_0$
- Packet maximum enters tunnel at  $t = 0.25$  and exits at  $t = 0.95$   
 $\Rightarrow \Delta t_{\text{tunnel}} \approx 0.7 > 0$  ( $v_{\text{tunnel}} = L/\Delta t_{\text{tunnel}} \approx 7.1 \rightsquigarrow k_{\text{tunnel}} \approx 3.6 \ll k_0$ )
- Like in case B: transmitted wave packet develops tail of lower maxima.